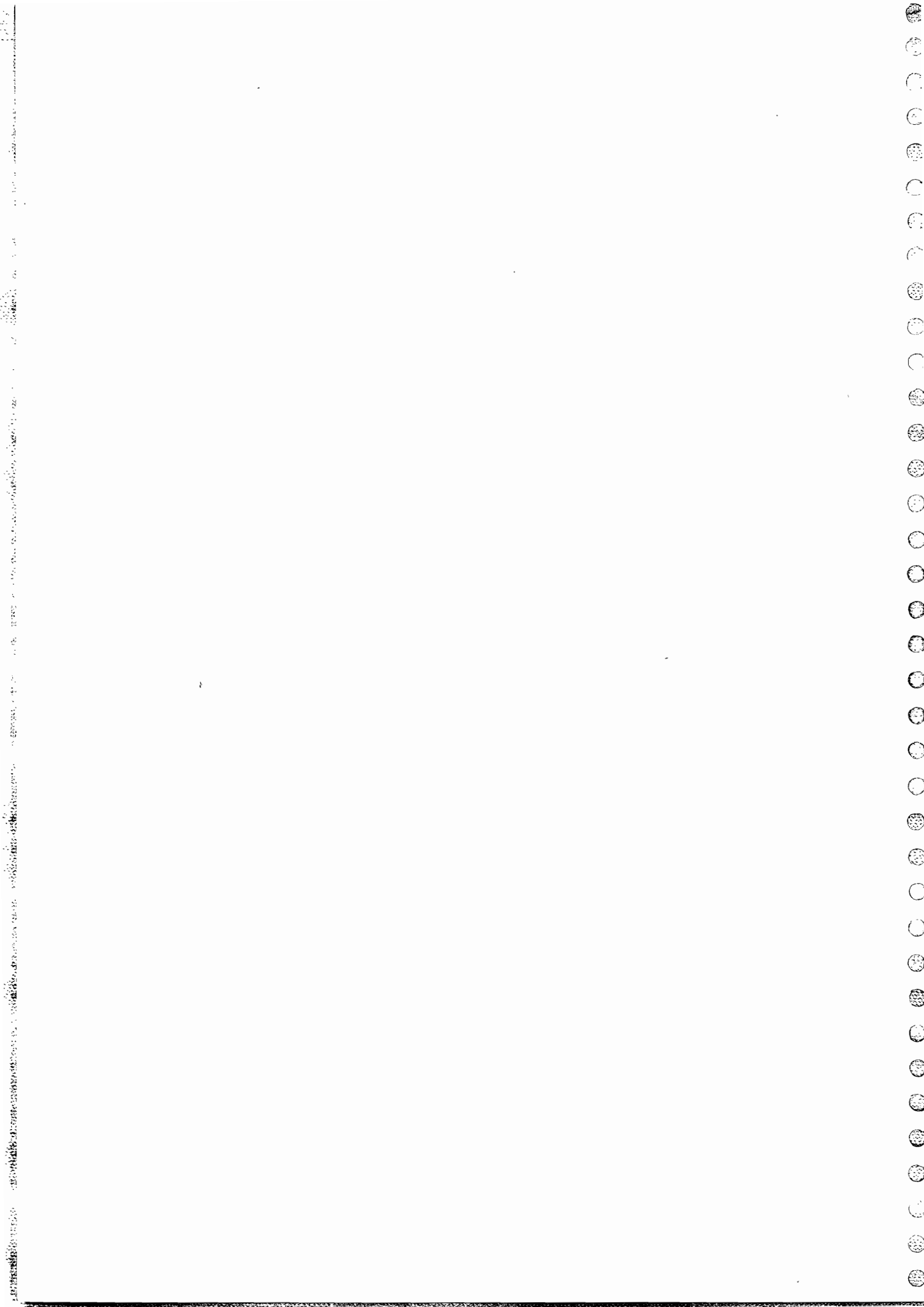


RISHAV GUPTA

PAPER I - Tut sheets

9873318559



Q1)

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Adiabatic Process is defines as the
process for which no heat is exchanged
by the system. For such a process,

$$P V^{\gamma} = \text{const} = \alpha \quad \text{--- (1)}$$

where P : Pressure of system

V : Volume of system

γ : ratio of specific heat at
const. Pressure and volume specific heat
ratio

$$\text{Work Done in Adiabatic Process} = \frac{(P_1 V_1 - P_2 V_2)}{\gamma - 1} \quad \text{--- (2)}$$

As given

$$V_2 = \frac{1}{2} V_1 \quad \text{--- (3)}$$

$\Rightarrow$ ~~from~~ ideal gas equation

$$\begin{aligned} \text{Volume of gas } V_1 &= \left(\frac{10^{-3}}{1.293} \right) \text{ m}^3 \\ &= 7.73 \times 10^{-4} \text{ m}^3 \end{aligned}$$

$$\text{Pressure } P_1 = 10^5 \text{ Pascal}$$

From (1) and (3)

$$P_1 V_1^{\gamma} = P_2 \left(\frac{V_1}{2} \right)^{\gamma}$$

$$\Rightarrow P_2 = P_1 \cdot 2^{\gamma} = \underline{\underline{2.63 \times 10^5 \text{ Pa}}}$$

From (2),

$$\begin{aligned} \text{Work Done} &= \left[10^5 \times 7.73 \times 10^{-4} - \frac{2.63 \times 10^5 \times 7.73 \times 10^{-4}}{2} \right] \times 0.4 \\ &= \underline{\underline{-60.87 \text{ J}}} \end{aligned}$$

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We can note that since the gas is compressed, Work done by the system is negative.

$$\text{External Work Done} = -dW = \underline{\underline{60.87J}}$$

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Q2) For an adiabatic process, heat exchanged is zero

$$\Rightarrow dQ = dU + PdV = 0$$

(from 1st law of thermodynamics)

$$\Rightarrow dU = -PdV$$

For an ideal gas, $dU = C_v dT$
as interactions between molecules are neglected.

$$\Rightarrow C_v dT = -PdV \quad \text{--- (1)}$$

Also from ideal gas equation,
 $PV = RT$, we have

$$P = \frac{RT}{V} \quad \text{--- (2)}$$

Using (2) in (1)

$$\Rightarrow C_v dT = -\frac{RT}{V} dV$$

$$\Rightarrow \frac{C_v}{R} \left(\frac{dT}{T} \right) = - \left(\frac{dV}{V} \right) \quad \text{--- (3)}$$

Now, we know for an ideal gas,

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$$C_p - C_v = R \quad \& \quad \frac{C_p}{C_v} = \gamma$$

$$\Rightarrow (\gamma - 1) = \frac{R}{C_v} \quad \text{--- (4)}$$

Using in (4)

$$\left(\frac{dT}{T}\right) = -\left(\frac{R}{C_v}\right) \frac{dv}{v} = -(\gamma - 1) \left(\frac{dv}{v}\right)$$

Integrating, we have

$$\ln T = -(\gamma - 1) \ln v + A$$

$$\Rightarrow T v^{\gamma - 1} = e^A = \text{const}$$

$$\Rightarrow \boxed{T v^{\gamma - 1} = \text{const}}$$

(Q3)

For an isothermal gas, temperature remains const. $\therefore$ for an ideal gas internal energy does not change

$$\Rightarrow \text{From 1st law of thermodynamics}$$

$$dQ = \overset{0}{dU} + PdV = PdV$$

Hence energy required is equal to work done by the system.

$$dW_{\text{isothermal}} = PdV$$

$$= \frac{nRT}{V} dV$$

$$\Rightarrow W = \int dW = nRT \ln \left(\frac{V_2}{V_1} \right)$$

$$= \frac{1}{2} \times 8.314 \times 300 \ln \left(\frac{25}{100} \right)$$

$$= \frac{-13457.69}{2} \text{ J}$$

$$= \underline{\underline{-1728.8 \text{ J}}}$$

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Hence energy has to be removed instead.

Q4)

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For a reversible expansion of an ideal gas from volume V_1 to V_2 , the gas is maintained at equilibrium with the surroundings at all intermediate points. This can be done by placing very small weights on the top of piston and slowly removing them one by one, maintaining equilibrium at all points.

However for an irreversible process of expansion, all weights are removed at once such that gas expands at once from V_1 to V_2 against const. external pressure P_2 .

The work done by gas in reversible process is

$$W_R = \int_{V_1}^{V_2} P dV$$

In irreversible process, Work done is

$$W_{ir} = P_2 (V_2 - V_1)$$

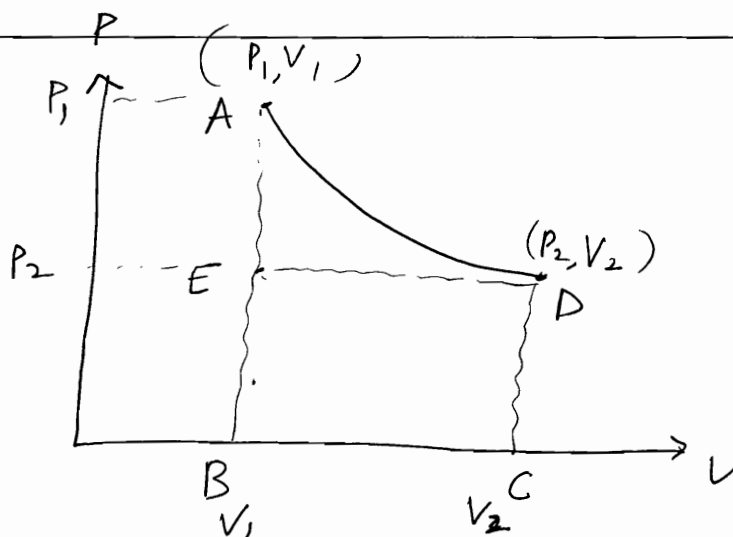
The difference in the two values of works is seen in following PV curve.

$$W_R = \text{Area (ABCD)}$$

$$W_{ir} = \text{Area (EBCD)}$$

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Clearly $W_R > W_{ic}$

Q5)

Any system ~~the~~ which constitutes matter is made up of a large number of molecules ^{state of} which are in constant motion and hence possess internal kinetic energy.

Due to interactions of molecules and intermolecular attraction, they also possess an internal potential energy. The sum total of the two energies of all molecules of the body is called internal energy of the system.

First law of thermodynamics states that if an amount of heat dQ is supplied to a system, it can be spent in two ways :

- (i) increase the internal energy of system, dU
- (ii) Work done by the system, dW

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$$\text{ie. } dQ = du + dw$$

The law is simply another form of Conservation of energy theorem i.e. energy can neither be destroyed nor be created but converted from one form to another form.

<Q3>

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Q6)

<Q1>

3 engineering example of adiabatic process include:

- (i) Heat Engine : Adiabatic ^{Expansion} ~~Compression~~ stroke does useful work.
- (ii) Refrigerator
- (iii) Adiabatic Demagnetization
- (iv) Liquification of gases via adiabatic expansion of gases which causes cooling

In engineering uses Adiabatic Process is performed by either

- (i) thermally insulating the system from surroundings or
- (ii) performing the process very rapidly so that there is no enough time for heat transfer to take place.

Q7)

Process 1

It's an isothermal Process.

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From 1st law of thermodynamics

$$dQ = dU + P dV \quad (\text{Isothermal})$$

$$\Rightarrow dQ = P dV \quad \text{--- (1)}$$

$$= \frac{RT}{V} dV \quad (\text{from ideal gas eqn})$$

$$\Rightarrow Q = \int dQ = RT \int \frac{dV}{V} = RT \ln\left(\frac{V_2}{V_1}\right) \quad \text{--- (2)}$$

Hence from (1) and (2), Heat absorbed
by the gas is equal to work done by
the system equal to $RT \ln\left(\frac{V_2}{V_1}\right)$.

Process 2

Adiabatic expansion

For adiabatic expansion, $TV^{\gamma} = \text{const.}$

For diatomic molecule, $\gamma = (7/5)$

$$\Rightarrow T_2 V_2^{\gamma} = T_1 V_1^{\gamma}$$

$$\Rightarrow T_2 \cdot (2)^{\frac{2}{5}} = T_1$$

$$\Rightarrow T_2 = \frac{T_1}{(2)^{2/5}} = \underline{\underline{0.75 T_1}}$$

Q8)

Second law of thermodynamics can be
written in two forms:

(1)

Kelvin - Planck Statement

It is impossible to ~~make~~ ^{Construct} a device
whose sole job is to take heat from a
source at higher Temperature and convert

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it into equivalent amount of work
without producing any changes in other
parts of the system.

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(2)

Claussius - Ostwald Statement

It is impossible to construct a
device whose ~~sole~~ job is to take heat
from lower temperature and release
it to higher temperature.

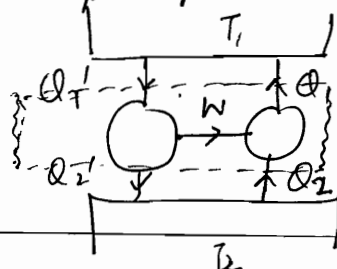
In terms of entropy, the second
law of thermodynamics can be states as
"For an isolated system, entropy cannot
decrease"

In order to Prove the Carnot
theorem, let us use method of
contradiction and assume that
efficiency of an arbitrary engine is
more than efficiency of Carnot engine.

$$\text{i.e. } \eta' > \eta_c \Rightarrow \eta' = \frac{W'}{Q_1'} \quad \& \quad \eta_c = \frac{W}{Q_1}$$

$$\text{as } \eta' > \eta_c \Rightarrow Q_1' < Q_1 \quad \text{Also, } Q_1' - Q_2' = W = Q_1 - Q_2$$

Since Carnot engine is based on
reversible cycle, it can applica in
reverse and can behave as a refrigerator
or heat pump.



We assume that
engines are
constructed to
give the same
amount of
work W.

Better to give
calculation
before & show
coupling afterwards

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Let us couple the two engine as shown in
adjacent figure.

Since we have assumed $\eta' > \eta_c$

$$\Rightarrow 1 - \frac{Q_2'}{Q_1'} > 1 - \frac{Q_2}{Q_1}$$

$$\Rightarrow \frac{W}{Q_1'} > \frac{W}{Q_1}$$

$$\Rightarrow Q_1 > Q_1'$$

$$\text{Hence } (Q_1 - Q_1') > 0$$

$\Rightarrow$ Net heat is being supplied to a
system at higher temperature; but
this violates 2nd law (Clausius-
Ostwald statement).

Hence our assumption is wrong.
No engine operating between 2 temp.
can be more efficient than Carnot
engine.

Q 9)

As given, $P = \alpha V$

Work done in the process $dW = PdV$

$$\Rightarrow dW = \alpha V dV$$

$$\begin{aligned} \Rightarrow W &= \int dW = \alpha \int_V^{mV} V dV \\ &= \alpha \cdot \frac{(m^2 - 1)V^2}{2} \\ &= \underline{\underline{\frac{\alpha (m^2 - 1)}{2} V^2}} \end{aligned}$$

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From ideal gas equation, we know

$$\frac{P_1 V_1}{T_1} = \frac{P_2 V_2}{T_2}$$

$$\Rightarrow \frac{\alpha V_1^2}{T_1} = \frac{\alpha V_2^2}{T_2} = \frac{\alpha m_1^2 V_1^2}{T_2}$$

$$\Rightarrow \boxed{T_2 = m_1^2 T_1}$$

For an ideal, change in internal energy, dU is given as

$$dU = C_v dT$$

$$\Rightarrow \Delta U = C_v (T_2 - T_1)$$

$$\boxed{\Delta U = C_v (m^2 - 1) T_1}$$

From 1st law of thermodynamics,
heat supplied dQ ,

$$dQ = dU + P dv$$

$$\Rightarrow \Delta Q = \Delta U + W$$

$$= C_v T_1 (m^2 - 1) + \frac{\alpha}{2} (m^2 - 1) V_1^2$$

$$= C_v \cdot \frac{P_1 V_1}{R} (m^2 - 1) + \frac{\alpha}{2} (m^2 - 1) V_1^2$$

$$= \frac{C_v \cdot \alpha (m^2 - 1) V_1^2}{R} + \frac{\alpha}{2} (m^2 - 1) V_1^2$$

$$= \alpha (m^2 - 1) \left[\frac{1}{2} + \frac{C_v}{R} \right]$$

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Now, ~~specific~~ ^{capacity} heat for the process C'
is given as

$$C' = \left(\frac{\Delta Q}{\Delta T} \right)$$

$$= \frac{V^2 \alpha (m^2 - 1) \left[\frac{1}{2} + \frac{C_v}{R} \right]}{(m^2 - 1) \cdot T}$$

$$= \frac{V^2 \alpha}{\frac{\alpha V^2}{R}} \left[\frac{1}{2} + \frac{C_v}{R} \right]$$

$$= \underline{\underline{\left[\frac{R}{2} + C_v \right]}}$$

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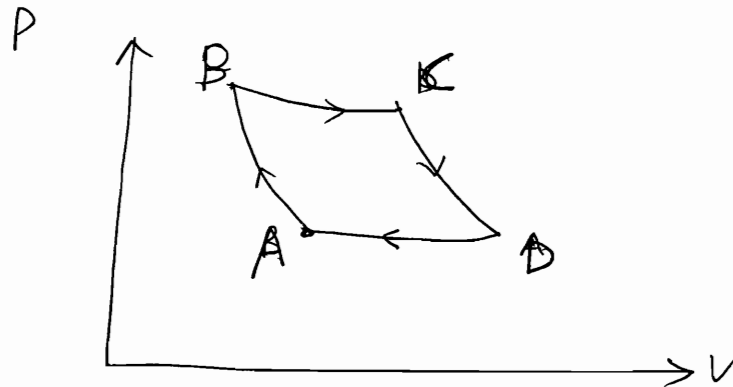
Test Sheet 21 : Carnot Cycle & Entropy

Q1)

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A Carnot engine use a Carnot Cycle
which is shown below.

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Carnot Cycle consists of Adiabatic
Compression (AB) followed by isothermal
expansion (BC), then adiabatic expansion
(CD), finally followed by isothermal
compression (DA).

For an isothermal process,

$$dQ = dU + PdV \quad (1^{st} \text{ law})$$

$$= PdV \quad (dU=0)$$

$$= \frac{RT}{V} \cdot dV \quad (\because dT=0 \text{ for isothermal})$$

(for ideal gas)

$$Q = \int dQ = RT \ln\left(\frac{V_2}{V_1}\right)$$

Now $Q_{AB} = 0$

$$Q_{BC} = RT_B \ln\left(\frac{V_C}{V_B}\right)$$

$$Q_{CD} = 0$$

$$Q_{DA} = -RT_D \ln\left(\frac{V_D}{V_A}\right)$$

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$$\Rightarrow \text{Efficiency} = \frac{W}{Q_{in}} = \frac{Q_1 + Q_2 + Q_3 + Q_4}{Q_{BC}}$$

$$= 1 - \frac{T_D \ln\left(\frac{V_D}{V_A}\right)}{T_B \ln\left(\frac{V_C}{V_B}\right)}$$

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Also for an adiabatic process,
 $TV^{r-1} = \text{const}$

$$\Rightarrow T_A \cdot V_A^{r-1} = T_B V_B^{r-1}$$

$$T_C V_C^{r-1} = T_D V_D^{r-1}$$

$$\Rightarrow \frac{T_C V_C^{r-1}}{T_B V_B^{r-1}} = \frac{T_D V_D^{r-1}}{T_A V_A^{r-1}} \quad \left(\begin{array}{l} T_C = T_B \\ T_D = T_A \end{array} \right)$$

$$\Rightarrow \left(\frac{V_C}{V_B} \right) = \left(\frac{V_D}{V_A} \right)$$

$$\Rightarrow \text{Efficiency} = \eta = 1 - \left(\frac{T_D}{T_B} \right)$$

$$= 1 - \left(\frac{T_2}{T_1} \right)$$

T_2 : Temperature of sink
 T_1 : Temperature of source

For given Carnot Engine $\eta = 1 - \frac{73}{273} = 73.2\%$

Q2)

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For an isothermal process,

$$dQ = dU + PdV \quad (1^{st} \text{ law})$$

$$= PdV \quad (\text{as } T = \text{const.})$$

$$\Rightarrow dS = \text{Change in entropy}$$

$$= \frac{dQ}{T}$$

$$= \frac{PdV}{T}$$

For an ideal gas $\frac{P}{T} = \frac{R}{V}$

$$\Rightarrow dS = R \frac{dV}{V}$$

$$\Rightarrow \Delta S = \int dS = R \int \frac{dV}{V} = R \ln\left(\frac{V_2}{V_1}\right)$$

$$= R \ln\left(\frac{4V}{V}\right)$$

$$= \boxed{R \ln 4}$$

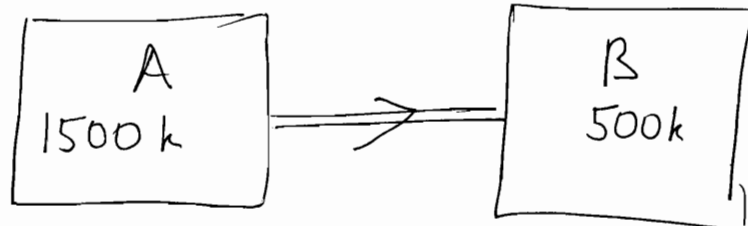
Q3)

L for condensation $= 3.36 \times 10^5 \text{ J/kg}$

$$\Rightarrow dQ = mL = 3.36 \times 10^5 \text{ J}$$

$$dS = \frac{dQ}{T} = \frac{mL}{T} = \frac{3.36 \times 10^5}{273} = \underline{\underline{1230.76 \text{ J/K}}}$$

Q4)



In 1 sec, $dQ = 10^4 \text{ J}$

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$$dS_A = -\frac{dQ}{T_A} \quad \left(\begin{array}{l} \text{negative as heat} \\ \text{flows from heat Temp} \\ \text{to lower temp} \end{array} \right)$$

$$dS_B = \frac{dQ}{T_B}$$

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$$\begin{aligned} \Rightarrow dS_{\text{Total}} &= dS_A + dS_B \\ &= dQ \left[\frac{1}{T_B} - \frac{1}{T_A} \right] \\ &= 10^4 \left[\frac{1}{500} - \frac{1}{1500} \right] \\ &= 10^4 \cdot \frac{1000}{500 \times 1500} \\ &= 13.33 \text{ J/K} \end{aligned}$$

$$\Rightarrow \text{Rate of entropy change} = 13.33 \text{ J/K/s}$$

Q5) Entropy is a measure of the disorder of the system. Greater the ~~disorder~~ disorder, more is the entropy. Entropy (S) is a state variable and thus for a process between initial and final states, Entropy change (dS) is independent of path (or the process)

If the system is at temperature T , and an amount dQ_{rev} of heat is added reversibly to the system, entropy change dS is defined as

$$dS = (dQ_{\text{rev}}/T)$$

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From 1st law,

$$\begin{aligned}d\phi &= dU + P dV \\&= C_v dT + P dV \quad (\text{for ideal gas})\end{aligned}$$

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$$\Rightarrow dS = \frac{d\phi}{T} = C_v \frac{dT}{T} + \frac{P}{T} dV$$

$$dS = C_v \frac{dT}{T} + \frac{R}{V} dV \quad \text{--- (1)}$$

$$\Rightarrow \Delta S = \int dS = C_v \ln\left(\frac{T_2}{T_1}\right) + R \ln\left(\frac{V_2}{V_1}\right)$$

For isothermal change $T_1 = T_2$

$$\Rightarrow \Delta S = \underline{R \ln\left(\frac{V_2}{V_1}\right)}$$

For isochoric change $V_2 = V_1$

$$\Rightarrow \Delta S = \underline{C_v \ln\left(\frac{T_2}{T_1}\right)}$$

To calculate isobaric change, we know

$$PV = RT$$

$$\Rightarrow P dV + V dP = R dT$$

$$\Rightarrow dV = \frac{R dT}{P} - \frac{V dP}{P}$$

Using in (1),

$$dS = C_v \frac{dT}{T} + \frac{R}{P} \frac{R}{V} dT - \frac{R}{V} \frac{V}{P} dP$$

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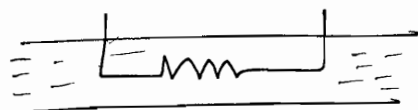
$$\begin{aligned}\Rightarrow dS &= C_v \frac{dT}{T} + \frac{R}{T} \cdot dT - \frac{R}{P} \cdot dP \\ &= (C_v + R) \frac{dT}{T} - \frac{R}{P} dP \\ &= C_p \frac{dT}{T} - \frac{R}{P} dP\end{aligned}$$

$$\Rightarrow \Delta S = C_p \ln\left(\frac{T_2}{T_1}\right) - R \ln\left(\frac{P_2}{P_1}\right)$$

For isobaric process, $P_1 = P_2$

$$\Rightarrow \Delta S = C_p \ln\left(\frac{T_2}{T_1}\right)$$

Q 6)



Given Resistance $R = 10 \Omega$

Current $I = 3 \text{ A}$

Temperature $T = 300 \text{ K}$

The energy obtained from the circuit (the source of EMF) is dissipated in water. For resistor, there is no net loss or gain of heat dQ

$$\text{i.e. } dQ_R = 0$$

$$\Rightarrow dS_R = \frac{dQ_R}{T} = 0$$

Since the Surrounding gains heat dQ at temperature 300 K , change in entropy dS

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In the process, state of resistor remains ~~constant~~ unchanged, as it is maintained at Temp 300 K . $\therefore$ entropy being a state function, change in entropy = 0

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is given as

$$dS_{\text{sur}} = \frac{dq}{T} = \frac{I^2 R dt}{T}$$

$$\Rightarrow \left(\frac{dS}{dt} \right)_{\text{sur}} = \frac{I^2 R}{T} = \frac{9 \times 10}{360} = 0.3 \text{ J/Ks}$$

Hence change in entropy per second is 0.3 J/K

$$\begin{aligned} \text{Since } \Delta S_{\text{univ}} &= \Delta S_{\text{system}} + \Delta S_{\text{surrounding}} \\ &= \Delta S_{\text{surrounding}} \end{aligned}$$

$$\Rightarrow \Delta S_{\text{universe per second}} = 0.3 \text{ J/K}$$

(Q7)

<Q5> <Q5>

Q8

<Q5>

For a process, specific heat is C_p given as

$$C_p = a + bT + cT^2$$

$$\text{Now } dq = m C_p dT$$

where m is mass of the system

Now, dS (change in entropy) is given as

$$dS = \frac{dq}{T} = \frac{m C_p}{T} dT$$

$$= \left(\frac{ma}{T} + mb + mcT \right) dT$$

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$$\Delta S = \int_{T_1}^{T_2} ds = m a \int_{T_1}^{T_2} \frac{dT}{T} + m b \int_{T_1}^{T_2} dT + m c \int_{T_1}^{T_2} T dT$$

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$$= m \left[a \ln \left(\frac{T_2}{T_1} \right) + b (T_2 - T_1) + c (T_2^2 - T_1^2) \right]$$

Q9)

Carnot theorem states that all reversible engines working between same heat reservoirs are equally efficient. Thus, any reversible heat engine operating between temperature T_1 and T_2 must have same efficiency. i.e. efficiency is a function of only temperatures

$$\eta = 1 - \frac{Q_C}{Q_H} \Rightarrow \frac{Q_C}{Q_H} = f(T_C, T_H)$$

$$\text{Let } T_3 > T_2 > T_1$$

$$f(T_3, T_1) = \frac{Q_1}{Q_3} = \left(\frac{Q_1}{Q_2} \right) \left(\frac{Q_2}{Q_3} \right) = f(T_1, T_2) f(T_2, T_3)$$

But left hand side is not a function of T_2 .

$\Rightarrow$ in order to achieve this

$f(T_1, T_2)$ should be of form

$$\frac{g(T_1)}{g(T_2)}, \text{ such that}$$

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$$f(T_1, T_2) f(T_2, T_3) = \frac{g(T_1)}{g(T_2)} \frac{g(T_2)}{g(T_3)}$$

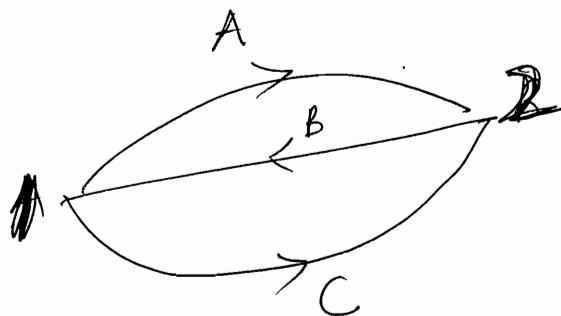
$$= \frac{g(T_1)}{g(T_3)} = f(T_1, T_3)$$

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As a matter of convenience, we choose
 $g(T) = T$. Choosing one temperature
as reference, we establish Thermodynamic
Temperature Scale. We choose Triple
point of water as 273.16 K and
Absolute Zero as 0 K .

Q 10)

Let us consider a cyclic path AB
and other ~~BP~~. Let both be
reversible.



From Clausius Inequality,

$$\oint_{\text{rev}} \frac{dQ}{T} = 0$$

$$\Rightarrow \oint \frac{dQ}{T} = 0 \Rightarrow \int_A^B \frac{dQ}{T} + \int_B^A \frac{dQ}{T} = 0 \quad \text{--- (1)}$$

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$$\text{Also, } \oint_{CB} \frac{dq}{T} = 0$$

$$\Rightarrow \int_C \frac{dq}{T} + \int_B \frac{dq}{T} = 0$$

$$\Rightarrow \int_C \frac{dq}{T} + \left(- \int_A \frac{dq}{T} \right) = 0 \quad \text{--- (from ①)}$$

$$\Rightarrow \boxed{\int_C \frac{dq}{T} = \int_A \frac{dq}{T}}$$

Hence path independent.

Entropy function (S) is defines a state function such that change in entropy dS for a reversible process is given as

$$\boxed{dS = \frac{dq_{rev}}{T}}$$

$$\Delta S = \int dS = \int_{rev} \frac{dq}{T}$$

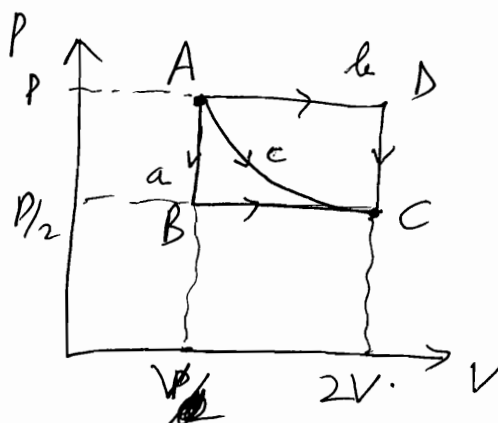
(For path a)

$$dS = dS_1 + dS_2$$

$$dS_1 = \frac{dq}{T}$$

$$= n \frac{dT}{T}$$

$$= n \ln(T_B/T_A)$$



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$$\frac{P_A V_A}{T_A} = \frac{P_B V_B}{T_B} \Rightarrow \frac{P V}{T_A} = \frac{P}{2} \frac{V}{T_B}$$

$$\Rightarrow T_B = \left(\frac{T_A}{2} \right)$$

$$\Rightarrow \Delta S_1 = -C_V \ln 2$$

$$dS_2 = C_P \frac{dT}{T} \Rightarrow \Delta S = C_P \ln \left(\frac{T_C}{T_B} \right)$$

$$\text{Also } \frac{P_B V_B}{T_B} = \frac{P_C V_C}{T_C}$$

$$\Rightarrow \frac{1}{T_B} = \frac{2}{T_C} \Rightarrow T_C = 2T_B = T_A$$

$$\Rightarrow \Delta S_2 = C_P \ln 2$$

$$\Rightarrow \Delta S_a = (C_P - C_V) \ln 2 = R \ln 2$$

For path c (isothermal path)

$$dS = \frac{dQ}{T} = \frac{P dV}{T} = \frac{R}{V} dV$$

$$\Rightarrow \Delta S = R \ln \left(\frac{V_C}{V_A} \right) = \underline{\underline{R \ln 2}}$$

Similarly for path b $\Delta S = R \ln 2$

Q11)

$$\frac{1}{4} = 1 - \frac{T}{(267 + 273)}$$

$$\Rightarrow T = \frac{3}{4} \times 540 = 405 \text{ K} = \underline{\underline{132^\circ \text{C}}}$$

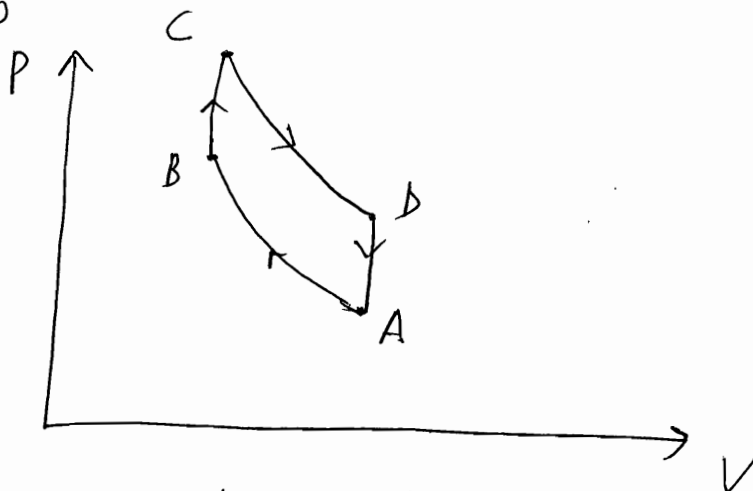
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Q12)
Q13)
Q14)

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Q-3
Q-1

Otto Cycle is a cycle used in Otto Engines that are used in Cars all over the world. The cycle is shown below



Otto Cycle starts with adiabatic compression (AB) followed by isochoric heating (BC), then adiabatic expansion (CD) [here useful work is done] and finally isochoric cooling (DA)

Efficiency of Otto Cycle

$$\begin{aligned} AB : dQ &= 0 \\ BC : dQ &= dW = C_V (T_C - T_B) \\ CD : dQ &= 0 \\ DA : dQ &= -C_V (T_D - T_A) \end{aligned}$$

$$\Rightarrow \eta = 1 + \frac{1/Q_{BC}}{Q_{DA}} = 1 - \frac{T_C - T_B}{T_D - T_A}$$

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Also,

$$V_B = V_C$$

$$V_D = V_A$$

$$T_C V_C^{\gamma-1} = T_D V_D^{\gamma-1}$$

$$T_A V_A^{\gamma-1} = T_B V_B^{\gamma-1}$$

$$\Rightarrow \frac{T_B V_B^{\gamma-1}}{T_C V_C^{\gamma-1}} = \frac{T_A V_A^{\gamma-1}}{T_D V_D^{\gamma-1}}$$

$$\Rightarrow \frac{T_B}{T_C} = \frac{T_A}{T_D}$$

$$\Rightarrow \eta = 1 - \frac{1/T_C (1 - \frac{T_B}{T_C})}{T_D (1 - \frac{T_A}{T_D})} = 1 - \frac{1}{\left(\frac{T_C}{T_D}\right)} = 1 - \left(\frac{T_D}{T_C}\right)$$

$$= 1 - \left(\frac{V_C}{V_D}\right)^{\gamma-1}$$

$$= 1 - \left(\frac{V_B}{V_A}\right)^{\gamma-1} \checkmark$$

$$\boxed{\eta = 1 - \left(\frac{1}{r}\right)^{\gamma-1}}$$

$$r: \text{Compression ratio} \\ = \left(\frac{V_A}{V_B}\right)$$

In Otto cycle, highest temperature is
for T_C and lowest for T_A

(Obtained by comparing isothermal process and

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adiabatic process slopes)

$$\Rightarrow \eta_{\text{Carnot}} = 1 - \frac{T_A}{T_E}$$

$$\text{Now } T_D > T_A$$

$$\Rightarrow \frac{T_D}{T_E} > \frac{T_A}{T_E}$$

$$\Rightarrow \left(1 - \frac{T_D}{T_E}\right) < \left(1 - \frac{T_A}{T_E}\right)$$

$$\Rightarrow \eta_{\text{Otto}} < \eta_{\text{Carnot}}$$

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Q15)

For isothermal process,

$$(i) dW = 0$$

$$(ii) dS = \frac{dQ}{T} = \frac{P dV}{T} = \frac{R}{V} dV =$$

$$\Rightarrow \Delta S = \int dS = R \ln \left(\frac{V_2}{V_1} \right)$$

$$\begin{aligned} (iii) \text{ Change in enthalpy} &= dH \\ &= dW + P dV + V dP \\ &= T dS + V dP \\ &= \frac{TR}{V} dV + \frac{RT}{P} dP \end{aligned}$$

$$\Rightarrow \Delta H = \int dH = TR \ln \left(\frac{V_2}{V_1} \right) + RT \ln \left(\frac{P_2}{P_1} \right)$$

$$= \cancel{RT \ln \left(\frac{V_2}{V_1} \right) + RT \ln \left(\frac{P_2}{P_1} \right)}$$

$$\begin{aligned} &= RT \ln \left(\frac{V_2}{V_1} \cdot \frac{P_2}{P_1} \right) \\ &= RT \ln 1 \\ &= 0 \end{aligned}$$

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For an adiabatic process,

$$dQ = 0$$

$$dU = -P dV$$

$$= -\frac{\alpha}{V^r} dV$$

$$(PV^r = \alpha)$$

$$\Rightarrow \Delta U = \int dU = -\alpha \int \frac{dV}{V^r}$$

$$= -\alpha \cdot (V_2^{-r+1} - V_1^{-r+1}) / (-r+1)$$

~~$$= -\alpha \cdot (V_2^{-r+1} - V_1^{-r+1}) / (-r+1)$$~~

$$= - \left[\frac{\alpha V_1^{-r+1} - \alpha V_2^{-r+1}}{r-1} \right]$$

★ Use

$$\Delta H = \Delta U + \Delta(PV)$$

$$= \Delta U + (P_2 V_2 - P_1 V_1)$$

$$= - \left[\frac{P_1 V_1^r \cdot V_1^{-r+1} - P_2 V_2^r \cdot V_2^{-r+1}}{r-1} \right]$$

$$= - \left[\frac{P_1 V_1 - P_2 V_2}{r-1} \right]$$

$$ds = \frac{dQ}{T} = 0 \quad (\text{as } dQ = 0)$$

$$dH = T ds + V dP$$

$$= V dP$$

$$= \beta \frac{dP}{P^{1/r}}$$

$$[P^{1/r} \cdot \alpha = \beta]$$

$$\Rightarrow \Delta H = \int dH = \beta \int \frac{dP}{P^{1/r}}$$

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★ Remember only for ideal gas $U=f(T)$
otherwise $U=f(T,V)$

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$$= \beta \left(P_2^{-\frac{1}{\gamma} + 1} - P_1^{-\frac{1}{\gamma} + 1} \right) \left[-\frac{1}{\gamma} + 1 \right]$$

$$= \frac{P_2^{\frac{1}{\gamma}} \cdot V_2 P_2^{-\frac{1}{\gamma} + 1} - P_1^{\frac{1}{\gamma}} V_1 P_1^{-\frac{1}{\gamma} + 1}}{-\frac{1}{\gamma} + 1}$$

$$= \frac{\gamma [P_2 V_2 - P_1 V_1]}{\gamma - 1}$$

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Q16)

From 1st law of thermodynamics,

$$dq = du + P dv$$

$$= C_v dT + P dv \quad (\text{for ideal gas}) \quad \text{--- (1)}$$

Now change in entropy ds is defined
as $ds = \left[\frac{dq_{rev}}{T} \right] \quad \text{--- (2)}$

Let us consider a mole of ideal gas
undergoing a reversible process s.t.

$$ds = \frac{dq}{T} = \frac{C_v dT}{T} + \frac{P}{T} dv$$

$$= C_v \left(\frac{dT}{T} \right) + \frac{R}{V} dv \quad \text{--- (3)}$$

⇒ Integrating

$$s = C_v \ln T + R \ln V + s_0$$

Now, for an ideal gas $Pdv + VdP = R dT$

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Using in (3),

$$dS = \frac{C_v}{T} \left(\frac{P dV + V dP}{R} \right) + \frac{R}{V} dV$$

$$= dV \left[\frac{C_v}{V} + \frac{R}{V} \right] + \frac{C_v V}{RT} \cdot dP$$

$$= C_p \frac{dV}{V} + C_v \frac{dP}{P}$$

⇒ Upon integration, we get

$$S = C_p \ln V + C_v \ln P + S_0$$

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Q17)

1st of all we need to find out the final condition of equilibrium.
Let us suppose whole ice melts and finally water is at $T^\circ\text{C}$.

$$\Rightarrow \text{Heat lost by water} = 10 \times 4.2 \times (30 - T) \\ = 42(30 - T) \text{ kJ}$$

$$\text{Heat gained by ice} = 1 \times [336 + 4.2 \times T] \\ = 336 + 4.2T$$

for equilibrium

$$42(30 - T) = 336 + 4.2T$$

$$\Rightarrow \boxed{T = 20^\circ\text{C}}$$

Since $0 < T < 30$, hence our assumption is correct.

Now,

$$\Delta S_{\text{system}} = \Delta S_{\text{water}} + \Delta S_{\text{ice}}$$

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$$\Delta S_{\text{water}} = \int ds = + \int mc \frac{dT}{T} = +mc \ln \left(\frac{T_2}{T_1} \right)$$

$$= 10 \times 4.2 \times \ln \left(\frac{273+20}{273+30} \right)$$

$$= \underline{\underline{-1.409 \text{ kJ/k}}}$$

$$\Delta S_{\text{ice}} = \int ds_1 + \int ds_2$$

$$= \int \frac{mL}{T} + \int mc \frac{dT}{T}$$

$$= \frac{336}{273} + 1 \times 4.2 \times \ln \left(\frac{273+20}{273} \right)$$

$$= +1.527 \text{ kJ/k}$$

$$\Rightarrow \Delta S_{\text{total}} = 0.118 \text{ kJ/k}$$

$$= \underline{\underline{118 \text{ J/k}}}$$

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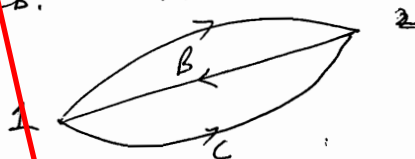
Q 18)

Proof of Entropy Change ≥ 0

Case I : Isolated System

Consider any system.

Let the system undergo a cyclic process, AB and CB.



st. A, B are reversibly executed and C is irreversible.

From Clausius Inequality,

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$$\oint_{rev} \frac{dq}{T} = 0$$

$$\Rightarrow \int_A \frac{dq}{T} + \int_B \frac{dq}{T} = 0$$

But from definition of entropy

$$\int_A \frac{dq}{T} = S_2 - S_1$$

$$\Rightarrow \int_B \frac{dq}{T} = - \int_A \frac{dq}{T} = S_1 - S_2$$

Also from Clausius Inequality,

$$\oint_{irr.} \frac{dq}{T} \leq 0$$

$$\Rightarrow \int_C \frac{dq}{T} + \int_B \frac{dq}{T} \leq 0$$

$$\Rightarrow \int_C \frac{dq}{T} \leq - \int_B \frac{dq}{T} = S_2 - S_1$$

$$\Rightarrow \boxed{S_2 - S_1 \geq \int_C \frac{dq}{T}} \quad \leftarrow \text{General derivation for any system}$$

Now for isolated system, $dQ = 0$

$$\Rightarrow \boxed{S_2 - S_1 \geq 0} \quad \leftarrow \text{specific for isolated system}$$

Case II : Open system

Let heat flows from body at temp T_1 to

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another body at temperature T_2

$$dS = \frac{dq}{T_2} - \frac{dq}{T_1} = dq \left[\frac{1}{T_2} - \frac{1}{T_1} \right]$$

Since heat flows from body at high temperature to body at lower temperature

$$\Rightarrow T_1 > T_2$$

$$\Rightarrow \frac{1}{T_1} < \frac{1}{T_2} \Rightarrow \frac{1}{T_2} - \frac{1}{T_1} > 0$$

$$\Rightarrow dS > 0$$

$$\Rightarrow \boxed{S_2 - S_1 > 0}$$

As proved, for a spontaneous system,

$$dS > \frac{dq}{T}$$

$$\Rightarrow T dS > dq = du + PdV$$

for constant Pressure $du + PdV \Leftrightarrow d(U + PV)$

$$\Rightarrow T dS > d(U + PV) = dH$$

$$\Rightarrow dH - T dS < 0$$

For temperature const, $dH - T dS \Leftrightarrow d(H - TS)$

$$\Rightarrow d(H - TS) < 0$$

$$\Rightarrow \boxed{dG < 0} \quad \text{--- (1)}$$

Now $G = H - TS$

$$\Rightarrow dG = VdP - SdT \quad (\text{Temp. and Pressure being const})$$

$$\Rightarrow \boxed{dG = 0} \quad \text{--- (2)}$$

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From ① and ②, we see that
for a spontaneous reaction, G can ~~only~~
only decrease and finally becomes
minimum at equilibrium state where
Pressure and Temperature become constant.

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Q19)

Final temperature of the system will
become equal to temperature of the
Reservoir

$$\Rightarrow T_2(\text{system}) = T_2$$

$$T_1(\text{system}) = T_1$$

$$dS_{\text{system}} = \frac{dq_{\text{gain}}}{T} = C_p \cdot \frac{dT}{T}$$

$$\Rightarrow \Delta S_{\text{system}} = C_p \ln\left(\frac{T_2}{T_1}\right)$$

$$\Delta S_{\text{reservoir}} = \frac{dq_{\text{lost}}}{T_2} = -\frac{C_p(T_2 - T_1)}{T_2}$$

$$\text{Now } \Delta S_{\text{universe}} = \Delta S_{\text{system}} + \Delta S_{\text{reservoir}}$$

$$= C_p \ln\left(\frac{T_2}{T_1}\right) + C_p \left(\frac{T_1}{T_2}\right) - C_p$$

$$= C_p \left[\frac{T_1}{T_2} - \ln\left(\frac{T_1}{T_2}\right) \right] - C_p$$

Claim $\Delta S > 0$

Proof let $\eta = \frac{T_2}{T_1} > 1$ (as $T_2 > T_1$)

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$$\text{let } f(\eta) = \ln \eta + \frac{1}{\eta} - 1$$

$$\text{Now } f(1) = 0$$

$$\text{Also } f'(\eta) = \frac{1}{\eta} - \frac{1}{\eta^2} = \left(\frac{\eta - 1}{\eta} \right)$$

$$\text{As } \eta > 1$$

$$\Rightarrow f'(\eta) > 0$$

i.e. $f(\eta)$ is monotonically increasing
function of η .

$$\Rightarrow f(\eta) > f(1)$$

$$\Rightarrow f(\eta) > 0$$

$$\Rightarrow \ln \eta + \frac{1}{\eta} - 1 > 0$$

$$\Rightarrow \Delta S = C_p f(\eta) > 0$$

Hence Proved.

The result is commensurate with
the fact that the entropy of the universe
always increases in an irreversible
process. (The process is irreversible as
the difference of temperature across
which heat flows is finite, i.e.
heat cannot flow in reverse direction)

Q20

Q & Tut 20

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Q 21) (ii)

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<Q 5> (iii) Q-18

(ii) Boltzmann Relation derivation

The disorder is expressed in terms of "thermodynamic probability" Ω which gives the number of microstates possible for a given configuration of particles. The entropy is, in fact, directly related to the disorder, greater the disorder, greater the entropy. To formulate relation between S and Ω , we consider the following:

Considering two independent systems with entropies S_1 and S_2 & thermodynamic probabilities Ω_1 and Ω_2 , we see that entropy being an extensive property, the total entropy S is

$$S = S_1 + S_2 \quad \text{--- (1)}$$

whereas for each microstate, in system 1 Ω_1 microstates are possible in system 2,

$\therefore$ total microstates = total thermodynamic probability $= \Omega = \Omega_1 \Omega_2$

$$\text{Let } S = f(\Omega)$$

$$\therefore S_1 = f(\Omega_1) \quad S_2 = f(\Omega_2)$$

$$S_f = f(\Omega_1, \Omega_2)$$

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$$\text{And } f(\Omega_1) + f(\Omega_2) = f(\Omega_1, \Omega_2)$$

taking partial derivative w.r.t. Ω_1 and
 Ω_2 ,

$$f'(\Omega_1) = \Omega_2 f'(\Omega_1, \Omega_2)$$

$$f'(\Omega_2) = \Omega_1 f'(\Omega_1, \Omega_2)$$

$$\Rightarrow \Omega_1 f'(\Omega_1) = \Omega_2 f'(\Omega_2) = \Omega_1 \Omega_2 f'(\Omega_1, \Omega_2)$$

$$\Rightarrow \Omega f'(\Omega) = \text{const} = k_B \quad (\text{say})$$

$$\Rightarrow df = k_B \frac{d\Omega}{\Omega} \quad \left(k_B \text{ is called Boltzmann Constant} \right)$$

$$\Rightarrow f(\Omega) = k_B \ln \Omega$$

$\Rightarrow$

$$S = k_B \ln \Omega$$

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Tut Sheet 22 : Thermodynamic Relationships

Q1)

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Consider a phase change process in which L amount of heat is absorbed, and volume of the phase, is V_1 and phase₂ be V_2 .

From Maxwell Relation

$$\left(\frac{\partial P}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T$$

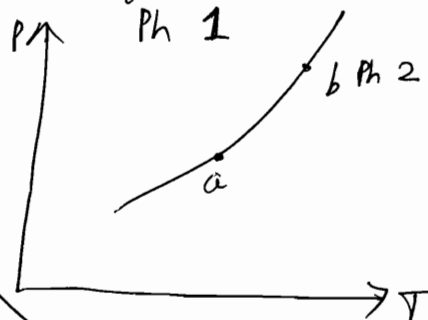
$$\Rightarrow R.H.S. = \left(\frac{dS}{dV}\right)_T = \left(\frac{T dS}{T dV}\right)_T$$

$$= \left(\frac{dQ}{T dV}\right)_T \approx \frac{L}{T \Delta V}$$

$$= \frac{L}{T(V_2 - V_1)}$$

$$\Rightarrow \boxed{\frac{dP}{dT} = \frac{L}{T(V_2 - V_1)}}$$

It can also be shown from
Gibbs' Relation



For transition
from a to b,
if both phases are
in equilibrium

$$\Rightarrow dg_1 = dg_2$$

$$\Rightarrow (V_1 dP - S_1 dT) = (V_2 dP - S_2 dT)$$

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$$\Rightarrow (S_2 - S_1) dT = (V_2 - V_1) dP$$

$$\Rightarrow \frac{dP}{dT} = \frac{S_2 - S_1}{(V_2 - V_1)} = \frac{L}{T(V_2 - V_1)}$$

Hence, Proved.

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For two substance in reversible equilibrium at constant temperature at pressure, the gibbs function remains constant.

Hence for a change of phase (phase i to phase f) at T and P, the specific gibbs functions are equal for two states

$$g^{(i)} = g^{(f)}$$

Similarly for a phase change at $T + dT$, $P + dP$

$$g^{(i)} + dg^{(i)} = g^{(f)} + dg^{(f)}$$

Subtracting,

$$dg^{(i)} = dg^{(f)}$$

$$\Rightarrow -S^{(i)} dT + V^{(i)} dP = -S^{(f)} dT + V^{(f)} dP$$

$$\Rightarrow S^{(f)} - S^{(i)} dT = V^{(f)} - V^{(i)} dP$$

$$\Rightarrow \left(\frac{dP}{dT} \right) = \frac{S^{(f)} - S^{(i)}}{V^{(f)} - V^{(i)}} = \frac{L}{T(V^{(f)} - V^{(i)})}$$

★ Clausius Clapeyron equation relates change in Pressure dP & change in Boiling / Melting Point dT .

Q2)

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Let us consider entropy S as a function
of Temperature T and Volume V
i.e. let $S = f(T, V)$

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$$\Rightarrow dS = \left(\frac{\partial S}{\partial T}\right)_V dT + \left(\frac{\partial S}{\partial V}\right)_T dV$$

Differentiating with Temperature at
Constant Pressure,

$$\left(\frac{\partial S}{\partial T}\right)_P = \left(\frac{\partial S}{\partial T}\right)_V + \left(\frac{\partial S}{\partial V}\right)_T \left(\frac{\partial V}{\partial T}\right)_P$$

$$\Rightarrow \left(\frac{\partial S}{\partial T}\right)_P - \left(\frac{\partial S}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T \left(\frac{\partial V}{\partial T}\right)_P$$

Multiplying by Temperature T

$$\Rightarrow T \left(\frac{\partial S}{\partial T}\right)_P - T \left(\frac{\partial S}{\partial T}\right)_V = T \left(\frac{\partial S}{\partial V}\right)_T \left(\frac{\partial V}{\partial T}\right)_P$$

From Maxwell relation $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$

Also using $T \left(\frac{\partial S}{\partial T}\right)_P = C_P$ and

$T \left(\frac{\partial S}{\partial T}\right)_V = C_V$, we get

$$C_P - C_V = T \left(\frac{\partial P}{\partial T}\right)_V \left(\frac{\partial V}{\partial T}\right)_P$$

Q3)

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Joule Thomson or Joule-kelvin
Coefficient is defined as $\left(\frac{\partial T}{\partial P}\right)_H$ i.e.
rate of change of Temperature with
change in Pressure at constant enthalpy.
Mathematically, expression for Joule
kelvin Coefficient comes as

$$\mu = \left(\frac{\partial T}{\partial P}\right)_H = \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T}\right)_P - V \right]$$

<derivation>

where C_p : specific heat capacity at
const. Pressure

For an ideal gas,

$$\left(\frac{\partial V}{\partial T}\right)_P = \frac{R}{P}$$

$$\begin{aligned} \Rightarrow \mu &= \frac{1}{C_p} \left[T \cdot \frac{R}{P} - V \right] \\ &= \frac{1}{C_p} \left[\frac{PV}{P} - V \right] = \frac{1}{C_p} [V - V] = 0 \end{aligned}$$

<derivation>

expressed as

Let enthalpy H be a function of T and
 P i.e. $H = f(T, P)$

$$\Rightarrow dH = \left(\frac{\partial H}{\partial P}\right)_T dP + \left(\frac{\partial H}{\partial T}\right)_P dT$$

For an isenthalpic process $dH = 0$

$$\Rightarrow \left(\frac{\partial H}{\partial P}\right)_T dP = - \left(\frac{\partial H}{\partial T}\right)_P dT$$

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Joule kelvin
effect says that
when a real
gas experiences
change of
Pressure at
constant
enthalpy,
it is accompanied
by a change in
Temperature.
Joule-kelvin
effect is
expressed in
terms of
Joule-kelvin
Coefficient, μ
 $\mu = \left(\frac{\partial T}{\partial P}\right)_H$

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$$\Rightarrow \left(\frac{dP}{dT} \right)_H = - \frac{\left(\frac{\partial H}{\partial P} \right)_T}{\left(\frac{\partial H}{\partial T} \right)_P} \quad \text{--- (1)}$$

we know $H = U + PV$, by definition

$$\Rightarrow dH = Tds + VdP$$

Replacing dH in (1), we get

$$\Rightarrow \left(\frac{dP}{dT} \right)_H = - \frac{1}{\left(\frac{Tds + VdP}{dT} \right)_P} \cdot \left(\frac{Tds + VdP}{dP} \right)_T$$

$$= - \frac{1}{\left(\frac{Tds}{dT} \right)_P} \cdot \left[T \left(\frac{\partial s}{\partial P} \right)_T + V \right]$$

$$= - \frac{1}{C_p} \left[T \cdot \left(\frac{\partial V}{\partial T} \right)_P + V \right]$$

(from Maxwell equation)

$$= \underline{\underline{\frac{+1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_P - V \right]}}$$

Q4)

From Clausius Clapeyron equation

$$\frac{dP}{dT} = \left(\frac{L}{T(V_2 - V_1)} \right)$$

$$\Rightarrow \frac{dP}{dT} = \frac{L}{\left(\frac{1}{d_2} \right) - \left(\frac{1}{d_1} \right)} \quad \left(\text{as } V \text{ is volume per unit mass} \right)$$

$$\Rightarrow dT = \frac{dP}{L} \cdot \left(\frac{1}{d_2} - \frac{1}{d_1} \right) T$$

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$$[\text{density}] = 1 \text{ g/cm}^3 = 10^{-3} \text{ kg/cm}^3 = 1 \text{ kg/l} = 1000 \text{ kg/m}^3$$

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$$= -\frac{6}{76} \times \frac{10^5}{2.3 \times 10^6} \left[\frac{1}{0.6} - \frac{1}{1000} \right]$$

$$= \underline{\underline{-2.13 \text{ k}}}$$

273
~~273~~
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For boiling, temperature of boiling drops when pressure is reduced. Hence difficult to cook at hills as high Temp can't be reached.

Q5) <Refer copy>

Q6)

$$\mu = \left(\frac{dT}{dP} \right)_H = \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_P - V \right]$$

<derivation>

$$= \frac{1}{8.21 \times 4.2} \left[T \cdot \frac{V-b}{T} \left(1 + \frac{2a}{RTV} \right) - V \right]$$

$$= \frac{1}{8.21 \times 4.2} \left[\frac{2a}{RT} - b \right]$$

$$= \text{[scribbled out]}$$

$$= \frac{1}{8.21 \times 4.2} \left[\frac{2 \times 1.39 \times 10^{-6} \times 10^5}{8314 \times 293} - 3.92 \times 10^{-5} \right]$$

$$= \underline{\underline{2.172 \times 10^{-6} \text{ k/Pascal}}}$$

<derivation>

To prove $\left(\frac{\partial V}{\partial T} \right)_P$ for real Vander Wall

$$\text{gas} = \frac{V-b}{T} \left[\text{[scribbled out]} \right]$$

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For Vander Walls Gas,

$$\left(P + \frac{a}{V^2}\right) (V - b) = RT$$

$$\left(\frac{\partial V}{\partial T}\right)_P = ?$$

$$P dV - \frac{a}{V^2} \cdot dV + \frac{2ab}{V^3} dV = R dT$$

$$\Rightarrow \left(\frac{\partial V}{\partial T}\right)_P = \frac{R}{P - \frac{a}{V^2} + \frac{2ab}{V^3}}$$

$$= \frac{R}{\left(P + \frac{a}{V^2}\right) - \frac{2a}{V^2} + \frac{2ab}{V^3}}$$

$$= \frac{R}{\frac{RT}{V-b} + \frac{2a}{V^2} \left[\frac{b}{V} - 1\right]}$$

$$= \frac{R}{\left(\frac{RT}{V-b}\right) \left[1 - \frac{\frac{2a}{V^2} \left(\frac{b}{V} - 1\right)}{\left(\frac{RT}{V-b}\right)}\right]}$$

$$\approx \frac{V-b}{T} \left[1 + \frac{2a}{V^2} \frac{V-b}{RT}\right]$$

$$\approx \frac{V-b}{T} \left[1 + \frac{2a}{RTV}\right]$$

Q7)

To prove $\left(\frac{\partial V}{\partial T}\right)_T = 0$ for ideal gas
 $\neq 0$ for real gas

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From 1st law

$$dU = dQ - PdV$$
$$= Tds - PdV$$

$$\Rightarrow \left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial s}{\partial V} \right)_T - P$$

$$\underline{\left(\frac{\partial U}{\partial V} \right)_T} = \underline{T \left(\frac{\partial P}{\partial T} \right)_V} - P \quad \text{(from Maxwell relation)}$$

For ideal gas

$$\left(\frac{\partial P}{\partial T} \right)_V = \left(\frac{R}{V} \right)$$

$$\Rightarrow \left(\frac{\partial U}{\partial V} \right)_T = T \cdot \frac{R}{V} - P = P - P = 0$$

Hence, internal energy independent of volume.

For a real gas, eg. Vanderwall gas

$$\left(\frac{\partial P}{\partial T} \right)_V = \left(\frac{R}{V-b} \right)$$

$$\Rightarrow \left(\frac{\partial U}{\partial V} \right)_T = T \cdot \frac{R}{V-b} - P = P + \frac{a}{V^2} - P$$
$$= \left(\frac{a}{V^2} \right) \neq 0$$

Hence internal energy depends upon volume.

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This is concomitant with the fact that for ideal gas, the molecules are point particles, they do not occupy volume and they do not interact. Their energy solely depends on temperature while the real gas molecules occupy volume and interact.

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Q8)

Physical significance of drop in Helmholtz free energy and Gibbs Free energy

Decrease in Helmholtz Free Energy (F),
 dF is given as

$$F = U - TS$$

$$\Rightarrow dF = dU - TdS - SdT$$
$$= -PdV - SdT$$

For constant temperature i.e. isothermal process

$$dF = -PdV$$

= Expansive Work done on the system

$$\Rightarrow -dF = PdV$$

$\therefore$ the drop in Helmholtz free energy is the amount of work done (expansive) by the system.

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From application of Clausius Inequality,
we know, for any process

$$ds \geq \left(\frac{dq}{T} \right)$$

$$\Rightarrow Tds \geq dq$$

This is the relation
to get inequalities
in thermodynamic
systems.....

Using 1st law,

~~First law of thermodynamics~~

$$dw = dq - du \leq Tds - du$$

define $F = U - TS$ [Helmholtz Free Energy]

$\therefore$ at constant temperature, $dF = du - Tds$

$$\Rightarrow dw \leq -dF$$

The maximum amount of work that
can be done in an isothermal
process is equal to decrease in the
Helmholtz free energy F .

$$dq \leq Tds$$

$$dw + pdv + d'w' \leq Tds$$

At const. P

$$dH + dw' \leq Tds$$

at const. T

$$dG + dw' \leq 0$$

$$dw' \leq (-dG)$$

$$\text{let } dw = dA + pdv$$

where dA is the non expansive work.

$$\Rightarrow dA = dq - dw - pdv$$

$$\leq Tds - du - pdv$$

$$\text{define } G = U + PV - TS$$

at const. temp, Pressure,

$$dG = dw + pdv - Tds$$

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Maximum
non expansive
work

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$$\Rightarrow dA \leq -dG$$

$\therefore$ the maximum amount of non expansive (non PdV) work that can done by the system at constant temp T and pressure P is equal to the decrease in the Gibbs Free Energy, G .

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Now F is defined as

$$F = U - TS \quad \text{--- (1)}$$

$$\begin{aligned} dF &= dU - Tds - SdT \\ &= -PdV - SdT \end{aligned}$$

$$\Rightarrow \left(\frac{\partial F}{\partial T} \right)_V = -S \quad \text{--- (2)}$$

Using (2) in (1)

$$F = U + T \cdot \left(\frac{\partial F}{\partial T} \right)_V$$

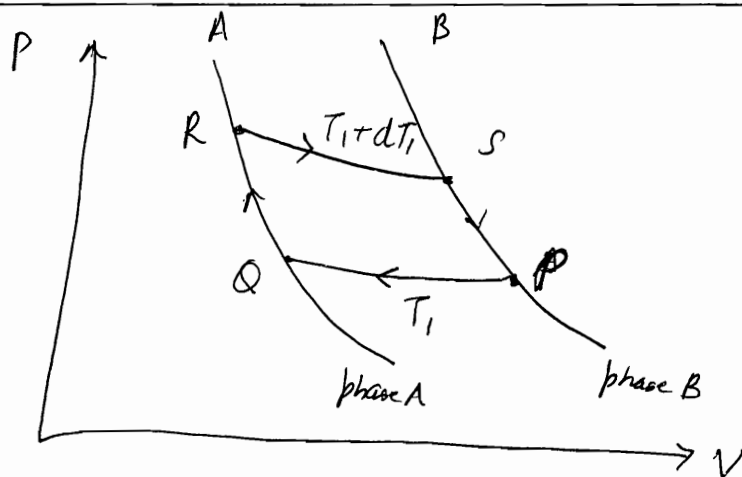
$$\Rightarrow \boxed{U = F - T \left(\frac{\partial F}{\partial T} \right)_V}$$

Q9)

Let the two phases be represented in adjoining figure as isocentric lines A and B for the two phases A and B.

Let phase change occur at Temp T_1 (isothermal process PQ) and the reverse

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Change of phase occurs at temperature $T + dT$. Let the latent heat for 1st phase change be L and second one be $L + dL$. The 4 process form Carnot Engine as can be seen PQRS. From Carnot Engine,

$$\frac{Q_1}{Q_2} = \frac{T_1}{T_2}$$

$$\Rightarrow \frac{L + dL}{L} = \frac{T_1 + dT_1}{T_1}$$

$$\Rightarrow \left(\frac{dL}{L} \right) = \left(\frac{dT_1}{T_1} \right)$$

$$\Rightarrow dL = \frac{L}{T_1} dT_1$$

$$\text{Also Work Done} = Q_1 - Q_2 = \Delta P \Delta V$$

$$\Rightarrow dL \approx P dV$$

$$\Rightarrow dP(dV) = \frac{L}{T_1} dT_1$$

$$\Rightarrow \frac{dP}{dT} = \frac{L}{T_1 dV} = \frac{L}{T(V_2 - V_1)}$$

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From the given data

$$L = T \cdot \Delta S = 273 \text{ K} \times 22.2 \times 10^3 \text{ J/K} \\ = 6.06 \times 10^6 \text{ J/K-mole}$$

1 kilo mole of ice = 18 kg

$$\Rightarrow \boxed{L = 336 \text{ kJ}}$$

Density of water = $\frac{1000 \text{ kg}}{\text{m}^3}$

Density of ice = $\frac{9}{10} \times 1000 = \frac{900 \text{ kg}}{\text{m}^3}$

$$dP = P_2 - P_1 = 2 \times 10^5 - 1 \times 10^5 = 10^5 \text{ N/m}^2$$

$$dT = ?$$

Now, from the derived Clausius Clapeyron equation,

$$\frac{10^5}{dT} = \frac{336 \times 10^3}{273 \left[\frac{1}{1000} - \frac{1}{900} \right]}$$

$$\Rightarrow \boxed{dT = -9.02 \times 10^{-3} \text{ K}}$$

Q10)

Let the entropy of liquid and saturated vapour be S_L and S_g . The change in entropy for the phase change can be written as

$$\Rightarrow S_g - S_L = \left(\frac{L}{T} \right)$$

$\Rightarrow$ Differentiating w.r.t. Temp T .

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$$\frac{ds_s}{dT} - \frac{ds_l}{dT} = \frac{1}{T} \left(\frac{dL}{dT} \right) - \left(\frac{L}{T^2} \right)$$

Multiplying by T ,

$$\left(\frac{T ds_s}{dT} \right) - \left(\frac{T ds_l}{dT} \right) = \left(\frac{dL}{dT} \right) - \left(\frac{L}{T} \right)$$

$$\Rightarrow \left(\frac{ds_s}{dT} \right) - \left(\frac{ds_l}{dT} \right) = \left(\frac{dL}{dT} \right) - \left(\frac{L}{T} \right)$$

$$\Rightarrow \boxed{C_s - C_l = \left(\frac{dL}{dT} \right) - \left(\frac{L}{T} \right)}$$

$$C_s = C_l + \left(\frac{dL}{dT} \right) - \left(\frac{L}{T} \right)$$

$$= 1.01 \text{ cal/k.g} + 0.64 \text{ cal/k.g} - \left(\frac{539}{373} \right) \text{ cal/k.g}$$

$$= \underline{\underline{-1.075 \text{ cal k}^{-1} \text{gm}^{-1}}}$$

Negative value of C_s implies that heat needs to be removed in order to increase the temperature. This is due to the fact that for increasing the temperature of saturated vapour, it needs to be compressed i.e. pressure must be increased. If compression is done adiabatically, the saturated vapour would become super heated vapour. $\therefore$ to convert to saturated vapour at higher temp, some amount of heat needs to be extracted out while compressing $\Rightarrow dQ < 0 \Rightarrow \boxed{C_s < 0}$

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Q11)

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$$\langle Q2 \rangle : C_p - C_v = T \left(\frac{\partial p}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P \quad - (1)$$

We can mathematically write Volume V as
a function of Temp T and Pressure P

$$P = f(T, V)$$

$$dP = \left(\frac{\partial P}{\partial T} \right)_V dT + \left(\frac{\partial P}{\partial V} \right)_T dV$$

For constant volume $dV = 0$

$$\Rightarrow 0 = \left(\frac{\partial P}{\partial T} \right)_V dT + \left(\frac{\partial P}{\partial V} \right)_T dV$$

$$\Rightarrow \left(\frac{\partial P}{\partial V} \right)_T dV = - \left(\frac{\partial P}{\partial T} \right)_V dT$$

$$\Rightarrow \left(\frac{dP}{dT} \right)_V = - \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial P} \right)_T$$

$$= - \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial P} \right)_T \quad - (2)$$

Putting (2) in (1)

$$C_p - C_v = -T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P \left(\frac{\partial P}{\partial V} \right)_T$$

$$\Rightarrow T \left(\frac{\partial V}{\partial T} \right)_P$$

$$= -T \left(\frac{\partial p}{\partial T} \right)_V \left(\frac{\partial p}{\partial T} \right)_V \left(\frac{\partial v}{\partial p} \right)_T$$

$$= -T \cdot \frac{\partial p}{\partial T}$$

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$$C_p - C_v = T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P \quad \text{--- (1) (Q2)}$$

Let us write Volume V as function of Temperature and Pressure

$$V = f(T, P)$$

$$dV = \left(\frac{\partial V}{\partial T} \right)_P dT + \left(\frac{\partial V}{\partial P} \right)_T dP$$

For constant volume $dV = 0$

$$\Rightarrow \left(\frac{\partial V}{\partial T} \right)_P dT = - \left(\frac{\partial V}{\partial P} \right)_T dP$$

$$\begin{aligned} \Rightarrow \left(\frac{\partial P}{\partial T} \right)_V &= - \left(\frac{\partial V}{\partial T} \right)_P / \left(\frac{\partial V}{\partial P} \right)_T \\ &= - \left(\frac{\partial V}{\partial T} \right)_P \cdot \left(\frac{\partial P}{\partial V} \right)_T \quad \text{--- (2)} \end{aligned}$$

Using (2) in (1), we get

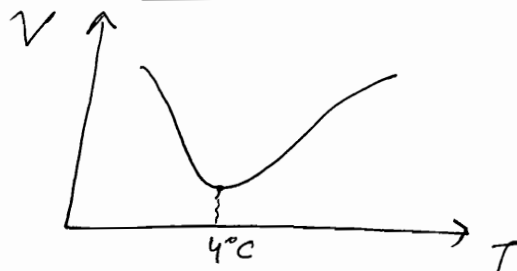
$$\Rightarrow C_p - C_v = - T \cdot \left(\frac{\partial V}{\partial T} \right)_P \left(\frac{\partial P}{\partial V} \right)_T \cdot \left(\frac{\partial V}{\partial T} \right)_P$$

$$C_p - C_v = - T \cdot \left(\frac{\partial V}{\partial T} \right)_P^2 \cdot \left(\frac{\partial P}{\partial V} \right)_T$$

At 4°C ,

$$\left(\frac{dV}{dT} \right) = 0$$

as the density is maximum



$$\Rightarrow \underline{C_p = C_v}$$

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For an ideal gas

$$\begin{aligned} C_p - C_v &= -T \cdot \left(\frac{\partial V}{\partial T} \right)_P^2 \left(\frac{\partial P}{\partial V} \right)_T \\ &= +T \cdot \left(\frac{R}{P} \right)^2 \left(\frac{P}{V} \right) \\ &= T \cdot \frac{R^2}{PV} = \underline{\underline{R}} \end{aligned}$$

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For a real gas, say Van der Waals' gas

$$\left(\frac{\partial V}{\partial T} \right)_P = \frac{V-b}{T} \cdot \left[1 + \frac{2a}{RTV} \right] \quad \text{Q-6}$$

$$\left(\frac{\partial P}{\partial T} \right)_V = \left(\frac{R}{V-b} \right)$$

$$\Rightarrow C_p - C_v = \cancel{T} \cdot \frac{R}{\cancel{V-b}} \cdot \frac{\cancel{V-b}}{\cancel{T}} \cdot \left[1 + \frac{2a}{RTV} \right]$$

$$C_p - C_v = R + \left(\frac{2a}{VT} \right)$$

$$R + \frac{2a}{VT} > R$$

$$\Rightarrow (C_p - C_v)_{\text{real gas}} > (C_p - C_v)_{\text{ideal gas}}$$

Q12)

Q3

Q13)

From Q2)

$$C_p - C_v = T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P$$

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Let $U = f(T, V)$

$$dU = \left(\frac{\partial U}{\partial T}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV$$

$$= \left(\frac{dQ - PdV}{dT}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV$$

$$= T \left(\frac{dS}{dT}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV$$

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$$\Rightarrow C_p - C_v = T \left(\frac{\partial S}{\partial V}\right)_T \cdot \left(\frac{\partial V}{\partial T}\right)_P$$

$$= \left(\frac{dU + PdV}{dV}\right)_T \left(\frac{\partial V}{\partial T}\right)_P$$

$$= \left[\left(\frac{\partial U}{\partial V}\right)_T + P\right] \left[\frac{\partial V}{\partial T}\right]_P$$

$$\text{R.H.S.} = \left(p + \frac{a}{V^2}\right) \cdot \left(\frac{V-b}{T}\right) \left[1 + \frac{2a}{RTV}\right]$$

$$= R \left[1 + \frac{2a}{RTV}\right]$$

$$\Rightarrow C_p = C_v + R + \left(\frac{2a}{TV}\right)$$

Q14) From Clausius Clapeyron equation,

$$\frac{dP}{dT} = \frac{L}{T(V_2 - V_1)}$$

$$\Rightarrow \frac{dP}{-0.91} = \frac{-80 \times 4.2}{(0.091) \times 10^{-6} \times 273}$$

$$\Rightarrow dP = 123 \text{ atm}$$

★ water $\rightarrow$ ice
 $\Rightarrow L$ negative
 $V_2 - V_1 > 0$

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Tut Sheet 23 Especial Topics

Q1)

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According to Curie's law for Paramagnetic Materials, the susceptibility χ of a Paramagnetic Material is inversely proportional to temperature

$$\text{i.e. } \chi = \frac{C}{T} \quad \text{where } C: \text{const.}$$

$$\Rightarrow \chi_1 T_1 = \chi_2 T_2$$

At very low temperatures, it's very difficult to measure Temperature by any ordinary means. Here this magnetic relation is used to define Temperature

$$T_1 = \frac{\chi_{\text{ref}} T_{\text{ref}}}{\chi_1}$$

χ_1 is measured through measurement of inductance.

When an external magnetic field is applied, dipoles are aligned in line with it. As soon as demagnetization is done adiabatically, the aligned dipole's potential energy is lost & dipoles are de-aligned. The loss of internal energy results in drop of Temperature.

In order to find Mathematical expression for drop in Temperature, Consider entropy as function of Temperature T and applied Magnetic Field H

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$$\text{ie. } S = f(T, H)$$

$$\Rightarrow dS = \left(\frac{\partial S}{\partial T}\right)_H dT + \left(\frac{\partial S}{\partial H}\right)_T dH$$

$$\begin{aligned} \Rightarrow TdS &= \left(\frac{\partial Q}{\partial T}\right)_H dT + T \left(\frac{\partial S}{\partial H}\right)_T dH \\ &= C_H dT + T \left(\frac{\partial S}{\partial H}\right)_T dH \end{aligned}$$

where C_H is heat capacity at constant external magnetic field.

Also Work done PdV is equivalent to Work Magnetic $-HdM$

$\therefore$ writing equivalent Maxwell equation

$$\left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P \quad \text{as}$$

$$\left(\frac{\partial S}{\partial H}\right)_T = +\left(\frac{\partial M}{\partial T}\right)_H$$

$$\Rightarrow TdS = C_H dT + T \left(\frac{\partial M}{\partial T}\right)_H dH$$

From Curie law $M = \frac{C}{T}$

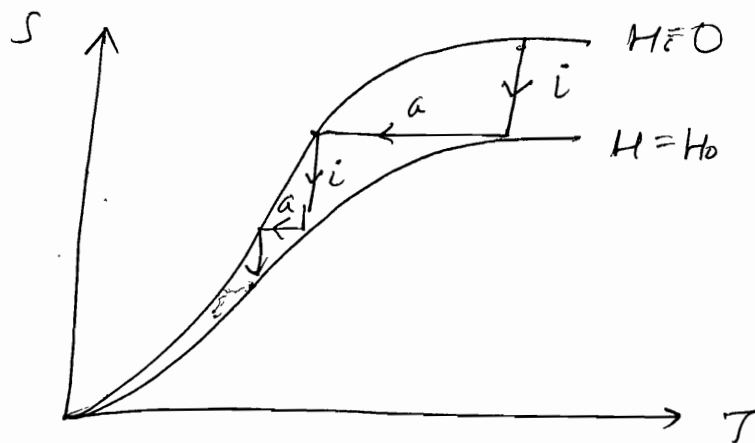
$$\Rightarrow \left(\frac{\partial M}{\partial T}\right)_H = -\frac{CH}{T^2}$$

$$\Rightarrow TdS = C_H dT - \frac{CH}{T} dH$$

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This is
derived
from Gibbs
Equation

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a : adiabatic demagnetization
 i : isothermal magnetization

For adiabatic demagnetization, $dS=0$

$$\Rightarrow C_H dT - \frac{C_H}{T} dH = 0$$

$$\Rightarrow C_H dT = \frac{C_H dH}{T}$$

$$\Rightarrow T dT = \left(\frac{C}{C_H}\right) H dH$$

$\Rightarrow$ Integrating both sides

$$\int_{T_1}^{T_2} T dT = \frac{C}{C_H} \int_H^0 H dH$$

$$\Rightarrow \frac{T_2^2 - T_1^2}{2} = \left(\frac{C}{C_H}\right) \cdot \frac{-H^2}{2}$$

$$\Rightarrow T_2 - T_1 \approx -\frac{C}{C_H} \frac{H^2}{2T} \quad (T_1 + T_2 \approx 2T)$$

$$\Rightarrow T_2 = T_1 - \frac{C}{C_H} \frac{H^2}{2T}$$

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At higher Temperature, drop in
temperature $\frac{CH^2}{C_H \cdot 2T}$ is not significant

Hence this cooling works better for
lower temperatures.

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Q2,3,4)

Q1

Q5)

Correction for a Vander Waal gas
need to be two fold

① Finite size of molecule

Since volume of real gas particles $\neq 0$

$\Rightarrow$ space for their movement or
interaction is reduced

$$\Rightarrow V_{\text{real}} = (V_{\text{available}} - b)$$

b: correction for volume of molecules.

② Intermolecular Attraction

Vander Wall found out that $\left(\frac{\partial U}{\partial V}\right)$ for
a real gas is not zero,

$$\text{Rather } \left(\frac{\partial U}{\partial V}\right) \approx \frac{A}{r^6} - \frac{B}{r^{12}}$$

$$\text{He consider } \left(\frac{\partial U}{\partial V}\right) = \frac{a}{V^2}$$

$$\text{Note that } \left(\frac{\partial U}{\partial V}\right) = \frac{\partial U}{A \cdot \partial r} = \frac{F}{A} = P$$

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$$\Rightarrow \text{Pressure Correction} = P_{\text{observed}} + \frac{a}{V^2}$$

After applying the corrections,

$$\left(P + \frac{a}{V^2}\right) (V - b) = RT$$

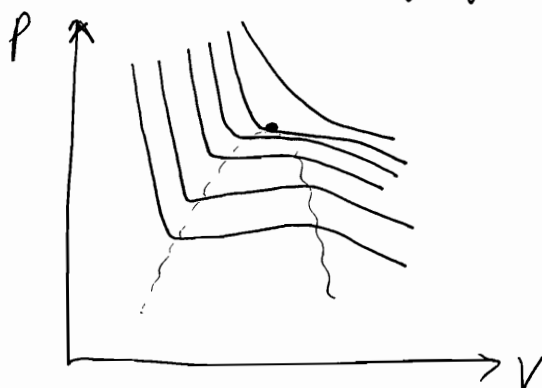
Expanding Vander Wall Equation,

$$(PV^2 + a)(V - b) = RTV^2$$

$$PV^3 - (Pb + RT)V^2 + aV - ab = 0$$

$$\Rightarrow V^3 - \left(b + \frac{RT}{P}\right)V^2 + \frac{a}{P}V - \frac{ab}{P} = 0 \quad \text{--- (1)}$$

Also from liquefaction of gases curve



At critical point all the three roots of Vander Wall gas converge at V_c

$\Rightarrow (V - V_c)^3 \equiv 0$ is equivalent equation of cubic equation (1)

$$\Rightarrow V^3 - V_c^3 - 3VV_c(V - V_c) = 0$$

$$\Rightarrow V^3 - 3V^2V_c + 3V_c^2V - V_c^3 = 0 \quad \text{--- (2)}$$

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$$V_c^3 = \frac{ab}{P_c} \quad - (3)$$

$$3V_c^2 = \frac{a}{P_c} \quad - (4)$$

$$3V_c = b + \frac{RT_c}{P_c} \quad - (5)$$

$$\Rightarrow \frac{V_c}{3} = b \quad \Rightarrow \quad V_c = 3b$$

$$\left[P_c = \frac{3V_c^2}{a} = \frac{3 \cdot 9b^2}{a} = \frac{27b^2}{a} \right] \Rightarrow P_c = \frac{a}{27b^2}$$

$$T_c = \frac{(3V_c - b) P_c}{R} = \frac{8b \cdot \frac{a}{27b^2}}{R}$$

$$\Rightarrow T_c = \frac{8a}{27Rb}$$

$$\begin{aligned} \text{Critical Coefficient} &= \frac{RT_c}{P_c V_c} \\ &= \frac{R \cdot \frac{8a}{27Rb}}{\frac{a}{27b^2} \cdot 3b} \\ &= \underline{\underline{\left(\frac{8}{3} \right)}} \end{aligned}$$

Since it's const., it's independent of
the gas taken !!

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Q8)

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(a) Q5

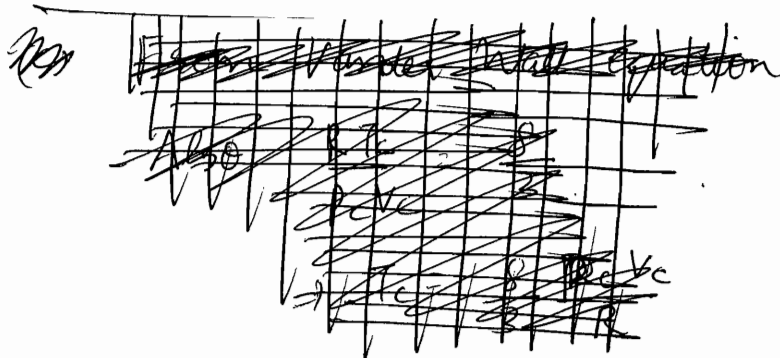
(b) In terms of Critical Parameters,

$$\text{as } V_c = 3b \quad \& \quad P_c = \frac{a}{27b^2}$$

$$\Rightarrow b = \frac{V_c}{3}$$

$$a = \frac{27 \cdot V_c^2 \cdot P_c}{9}$$

$$a = 3V_c^2 P_c$$



$$\text{Also } T_c = \frac{8a}{27Rb}$$

$$\Rightarrow a = \frac{27Rb T_c}{8} \\ = \frac{27R T_c \cdot V_c}{3 \times 8}$$

$$a = \frac{9R T_c V_c}{8}$$

(c) From given Parameters,

$$V_c = \frac{3}{8} \frac{RT_c}{P_c} = \frac{3}{8} \times \frac{80 \times 10^{-6} \times 10^5 \times 154.2}{49.7 \times 10^5} \\ = 9.3 \times 10^{-5} \text{ m}^3$$

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We can also
write b
as $4Nv$
where v is
volume occupied
by 1 molecule
 N : Avogadro's
number.



$$N \times \frac{4V_c}{2} \\ = \frac{N \times 4V_c}{2} \\ = 4NV_c$$

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$$\Rightarrow b = \frac{V_c}{3} = \frac{3.1 \times 10^{-5} \text{ m}^3}{3}$$

$$\begin{aligned} a &= 3 V_c^2 p_c \\ &= 3 \times 9.3^2 \times 10^{-15} \times 49.7 \times 10^5 \\ &= 1.28 \times 10^{-6} \text{ m}^6 \text{ Pascal} \end{aligned}$$

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Q8

For a closed PVT system, from the combined 1st and 2nd law, we get

$$dU = Tds - PdV \quad \text{--- (1)}$$

$$\Rightarrow U \equiv U(S, V)$$

$$\therefore dU = \left(\frac{\partial U}{\partial S} \right)_V dS + \left(\frac{\partial U}{\partial V} \right)_S dV \quad \text{--- (2)}$$

From (1) & (2),

$$\Rightarrow \left(\frac{\partial U}{\partial S} \right)_V = T \quad \left(\frac{\partial U}{\partial V} \right)_S = -P$$

The internal energy U is an extensive property and is proportional to number of moles included in the system.

When the system is open, so that we can add ~~to~~ remove material, internal energy becomes a function of 'n' along with S and V

$$dU = \left(\frac{\partial U}{\partial S} \right)_{V,n} dS + \left(\frac{\partial U}{\partial V} \right)_{S,n} dV + \left(\frac{\partial U}{\partial n} \right)_{S,V} dn$$

Tut Sheet 24 : Specific Heat of Solids

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The coefficient of dn is defined as
chemical potential μ ,

$$\mu = \left(\frac{\partial U}{\partial n} \right)_{S, V}$$

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You can
do this
question
better
by
taking
G instead
of
U

Thus Chemical Potential is change of
internal energy per mole. Hence, it is
intensive as it is per mole.

Since the change is at constant S and
 V , it is dependent only on remaining
thermodynamic parameters viz, Pressure
 P and Temperature T .

$$\Rightarrow \mu = \left(\frac{\partial G}{\partial n} \right)_{P, T}$$

Since μ is
defined for a
particular
value of
 P and T ,
Hence it is
a function
of these
two parameters.

Q1)

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Thermal Energy E , is given as

$$E = \int_0^{\gamma_m} \frac{h\nu g(\nu) d\nu}{e^{\frac{h\nu}{k_B T}} - 1}$$

$$= \int_0^{\gamma_m} \frac{h\nu \cdot 6N h^2 \nu}{\left(e^{\frac{h\nu}{k_B T}} - 1\right) (k_B \theta_D)^2} d\nu$$

$$\text{Put } \frac{h\nu}{k_B T} = x$$

$$\text{Also } \theta_D = \frac{\gamma_m h}{k_B}$$

$$= \frac{6N h^3}{(k_B \theta_D)^2} \left(\frac{k_B T}{h}\right)^3 \int_0^{\frac{\theta_D}{T}} \frac{x^2 dx}{(e^x - 1)}$$

$$= \frac{6N \cdot k_B^3 T^3}{k_B^2 \cdot \theta_D^2} \int_0^{\frac{\theta_D}{T}} \frac{x^2}{e^x - 1} dx$$

$$E = 6R \cdot \frac{T^3}{\theta_D^2} \int_0^{\frac{\theta_D}{T}} \frac{x^2}{e^x - 1} dx$$

$$\text{for } T \ll \theta_D \quad \frac{\theta_D}{T} \rightarrow \infty$$

$$\Rightarrow E \approx 6R \frac{T^3}{\theta_D^2} \int_0^{\infty} \frac{x^2}{e^x - 1} dx$$

$$= 2.404 \cdot 6R \left(\frac{T^3}{\theta_D^2}\right)$$

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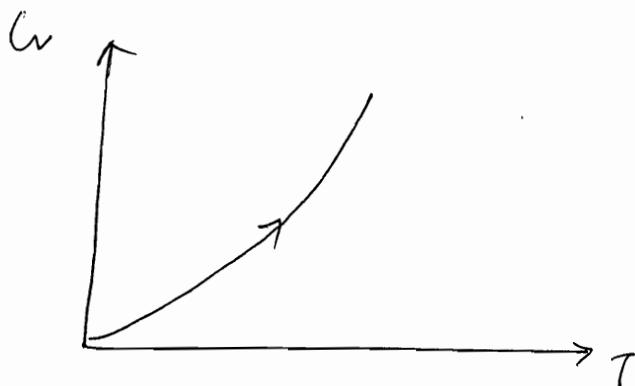
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$$\Rightarrow E = \left[\frac{14.424 R}{\theta_D^2} \right] \cdot T^3$$

$$\text{Now } C_v = \frac{dE}{dT}$$

$$= \frac{14.424 R}{\theta_D^2} \cdot 3T^2$$

$$= \left(\frac{43.272}{\theta_D^2} \right) T^2$$



Parabolic variation. (low temperature)

Q2)

(30 marks)

Einstein's theory of specific heat was a major improvement over classical theory of Dulong and Petit, as it explained that C_v drops with temperature.

Instead of a classical approach, Einstein followed quantum theory to derive specific ~~heat~~ heat. He assumed:

①

All molecules are Planck Oscillators vibrating quantum mechanically, with

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a frequency ν , characteristic of a material.

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(2) Average energy per mode of vibration is given as $\left(\frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} \right)$

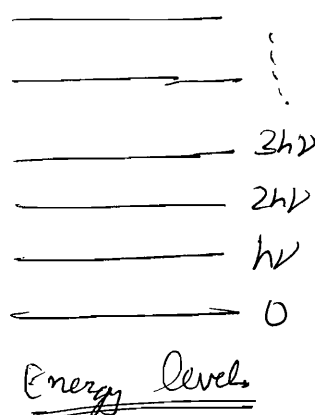
(3) Total number of modes of vibration are $3N$. as each molecule has 3 degrees of freedom.

(4) Maxwell distribution of energy levels hold. Energy levels are discrete in units of $h\nu$.

Claim $\langle E \rangle = \left(\frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} \right)$

Proof

According to Plank's theory, energy levels are discrete in unit of $h\nu$.



Also from ~~Plank's~~ Maxwell distribution, Molecules in n^{th} energy level are ~~given~~

$$N_n = N_0 e^{-\left(\frac{nh\nu}{kT}\right)}$$

$$\Rightarrow N = N_0 + N_1 + N_2 + \dots = N_0 \left[1 + e^{-\frac{h\nu}{kT}} + e^{-\frac{2h\nu}{kT}} + \dots \right]$$

$$= \left(\frac{N_0}{1 - e^{-\frac{h\nu}{kT}}} \right)$$

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Also, total energy is

$$E = N_0 E_0 + N_1 E_1 + N_2 E_2 + \dots$$

$$= N_0 h\nu e^{-\frac{h\nu}{kT}} + 2N_0 h\nu e^{-\frac{2h\nu}{kT}} + \dots$$

This is AG series,

$$\Rightarrow E = \frac{N_0 h\nu \cdot e^{-\frac{h\nu}{kT}}}{\left(1 - e^{-\frac{h\nu}{kT}}\right)^2}$$

Now the average energy per molecule per degree of freedom is

$$\Rightarrow \langle E \rangle = \frac{E}{N}$$

$$= \frac{h\nu e^{-\frac{h\nu}{kT}}}{1 - e^{-\frac{h\nu}{kT}}}$$

$$\boxed{\langle E \rangle = \left[\frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} \right]}$$

Using these assumptions, Einstein proposed total energy of system of molecules E , is given as for 1 gram mole of solid

$$E = 3N \cdot \langle E \rangle$$

$$= \frac{3N h\nu}{\left(e^{\frac{h\nu}{kT}} - 1\right)}$$

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New specific heat of solid Cr is defined
as,

$$C_v = \left(\frac{dE}{dT} \right)$$

$\therefore C_v$ in Einstein theory is given as,

$$C_v = + \frac{3N h \nu}{\left(e^{\frac{h\nu}{kT}} - 1 \right)^2} \cdot e^{\frac{h\nu}{kT}} \cdot \frac{h\nu}{kT^2}$$

Put $\frac{h\nu}{k} = \theta_E$: Einstein Temperature

$$\Rightarrow C_v = 3Nk \cdot \left(\frac{h\nu}{k^2 T^2} \right) \cdot \frac{e^{\frac{\theta_E}{T}}}{\left(e^{\frac{\theta_E}{T}} - 1 \right)^2}$$

$$C_v = 3R \left(\frac{\theta_E}{T} \right)^2 \frac{e^{\frac{\theta_E}{T}}}{\left(e^{\frac{\theta_E}{T}} - 1 \right)^2}$$

For $T \gg \theta_E$

$$\frac{\theta_E}{T} \ll 1 \Rightarrow e^{\frac{\theta_E}{T}} - 1 \approx \left(\frac{\theta_E}{T} \right)$$

$$\Rightarrow C_v = 3R \cdot \left(\frac{\theta_E}{T} \right)^2 \cdot \frac{\left(1 + \frac{\theta_E}{T} \right)}{\left(\frac{\theta_E}{T} \right)^2}$$

$$C_v \approx 3R$$

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∴ for higher Temperature, Einstein specific heat is equal to Dulong & Petit specific heat.

At lower temp $T \ll \Theta_E$

$$\frac{\Theta_E}{T} \gg 1$$

$$\begin{aligned} \Rightarrow C_v &= 3R \cdot \left(\frac{\Theta_E}{R}\right)^2 \cdot \frac{e^{-\frac{\Theta_E}{R}}}{\left(e^{-\frac{\Theta_E}{R}}\right)^2} \\ &= 3R \cdot \left(\frac{\Theta_E}{R}\right)^2 \frac{1}{1 + \left(\frac{\Theta_E}{T}\right) + \left(\frac{\Theta_E}{T}\right)^2 + \dots} \\ &\approx 0 \end{aligned}$$

Hence for lower temperature, $C_v \rightarrow 0$.

Shortcomings

- ① Einstein assumed all molecules to have same frequency. This is not a good assumption.
- ② Θ_E or Einstein Temperature is difficult to measure.
- ③ At low temperature C_v is found to vary as T^3 with Temperature T . This is not proved by Einstein theory.

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In Einstein theory, C_v falls much faster as compared to observed behaviour.

Q3)

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<Q2>

$$C_v = 3R \left(\frac{\theta_E}{T} \right)^2 \cdot \frac{e^{\frac{\theta_E}{T}}}{\left(e^{\frac{\theta_E}{T}} - 1 \right)^2}$$

For $T = T_E/2$ or $(\theta/2)$, we have

$$C_v = 3R \cdot (2)^2 \cdot \frac{e^2}{(e^2 - 1)^2} = \underline{\underline{2.17 R}}$$

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Q4)

The expression for specific heat are as:

Einstein: $C_v = 3R \left(\frac{\theta_E}{T} \right)^2 \frac{e^{\left(\frac{\theta_E}{T} \right)}}{\left[e^{\left(\frac{\theta_E}{T} \right)} - 1 \right]^2}$

Debye: $C_v = 9R \left[4 \left(\frac{T}{\theta_D} \right)^3 \int_0^{\frac{\theta_D}{T}} \frac{x^3}{e^x - 1} dx \right.$

$\left. - \left(\frac{\theta_D}{T} \right)^4 \frac{1}{e^{\theta_D/T} - 1} \right]$

1st expression
साब बड़ा बड़ा
4
Power: 3
गुंदा integral

2nd expression
- no integer
- no Power
- no integral

Remember:

$$E_{\text{ein}} = 3N \cdot \frac{h\nu}{\left(e^{\frac{h\nu}{kT}} - 1 \right)}$$

$$E_{\text{debye}} = \frac{9RT^4}{\theta_D^3} \int_0^{\left(\frac{\theta_D}{T} \right)} \frac{x^3}{e^x - 1} dx$$

$$x = \left(\frac{h\nu}{kT} \right)$$

$$\theta_E = \left(\frac{h\nu}{k} \right)$$

$$\theta_D = \left(\frac{h\nu_m}{k} \right)$$

Difference between Debye and Einstein
Theory are:

①
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Einstein assume single frequency with
with all molecules are said to be vibrating
in harmonic oscillations.
while

②

Debye considered a spectrum of frequency
of vibration varying from 0 to ν_m .

② No. of degrees of freedom for Einstein are 3
while

In Debye, no. of modes of vibration are
frequency dependent.

For lower temperature region, Debye's
specific heat is more accurate.

For lower temperature $T \ll \Theta_D$

$$\Rightarrow \frac{\Theta_D}{T} \gg 1$$

$$\Rightarrow C \approx 9R \left[4 \left(\frac{T}{\Theta_D} \right)^3 \int_0^\infty \frac{x^3}{e^x - 1} dx - \left(\frac{\Theta_D}{T} \right) \frac{1}{e^{\frac{\Theta_D}{T}}} \right]$$

$$\approx 9R \left[4 \cdot \frac{T^3}{\Theta_D^3} \cdot \lambda \right]$$

$$\text{where } \lambda = \int_0^\infty \frac{x^3}{e^x - 1} dx$$

$$\approx \frac{36R\lambda}{\Theta_D^3} T^3$$

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③ While Einstein assumed, vibrations of atoms are independent of each other, Debye postulate that the vibrations of atomic oscillations are strongly coupled with each other.

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$$\begin{aligned}\Delta Q_{\text{required}} &= \int C_v dT \\ &= \frac{36R\lambda}{\theta_D^3} \cdot \int_1^{T_2} T^3 dT \\ &= \frac{36R\lambda}{\theta_D^3} \cdot \frac{(T_2^4 - T_1^4)}{4}\end{aligned}$$

$$\Delta Q = \frac{9R\lambda}{\theta_D^3} (T_2^4 - T_1^4)$$

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Q5) Debye's theory was an improved upon Einstein's theory, Debye assumed that:

i) Molecular Oscillations are not independent of each other rather strongly coupled with each other.

ii) There is no single frequency but rather a continuum of frequencies with which the molecules vibrate with a maxima of γ_m .

iii) Total number of modes of vibration are $3N$
From phase space, we know no. of modes of vibration for momentum p , is given as

$$g(p) dp = \frac{2\pi V}{h^3} 4\pi p^2 dp$$

Assuming $p = \frac{h\nu}{v}$, we have

$$g(\nu) d\nu = \frac{2\pi V}{h^3} 4\pi \left[\frac{\nu^2}{v_l^2} + \frac{2}{v_t^2} \right] d\nu = \underline{\underline{A\nu^2 d\nu}}$$

where v_l and v_t are longitudinal &

★ Note that both Einstein & Debye theory use maxwell boltzmann distribution !!!

of elastic waves

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transverse velocities. — Factor of 2 due to 2 modes of polarization.

But total modes of vibration are $3N$

$$\Rightarrow 3N = \int_0^{\gamma_m} A \cdot \gamma^2 d\gamma$$

$$\Rightarrow 3N = \frac{A \cdot \gamma_m^3}{3} \Rightarrow A = \frac{9N}{\gamma_m^3}$$

$$\text{Hence } g(\gamma) d\gamma = \frac{9N}{\gamma_m^3} \gamma^2 d\gamma$$

Total energy is given as $\int g(\gamma) d\gamma E(\gamma)$

Now from Planck theory and Maxwell distribution,

$$E(\gamma) = \frac{h\gamma}{(e^{\frac{h\gamma}{kT}} - 1)}$$

$$\Rightarrow E = \int_0^{\gamma_m} \frac{9N}{\gamma_m^3} \gamma^2 \frac{h\gamma}{e^{\frac{h\gamma}{kT}} - 1} d\gamma$$

$$= \frac{9Nh}{\gamma_m^3} \int_0^{\gamma_m} \frac{\gamma^3 d\gamma}{e^{\frac{h\gamma}{kT}} - 1}$$

$$\text{Put } \boxed{\frac{h\gamma}{kT} = x}$$

$$\text{And } \boxed{\frac{h\gamma_m}{k} = \theta_D}$$

$$= \frac{9Nh}{\gamma_m^3} \left(\frac{kT}{h} \right)^4 \int_0^{\frac{\theta_D}{T}} \frac{x^3 dx}{e^x - 1} = \frac{9RT^4}{\theta_D^3} \int_0^{\frac{\theta_D}{T}} \frac{x^3}{(e^x - 1)} dx$$

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Now to calculate C_v , we must know
Leibnitz rule

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$$\frac{d}{dx} \int_{g(x)}^{f(x)} \lambda(y) dy = \lambda(f(x)) f'(x) - \lambda(g(x)) g'(x) \quad \text{--- (1)}$$

$$\Rightarrow C_v = \frac{dE}{dT} = \frac{qR}{\theta_D^3} \left[4T^3 \int_0^{\frac{\theta_D}{T}} \frac{x^3}{e^x - 1} dx + T^4 \cdot \frac{d}{dT} \int_0^{\frac{\theta_D}{T}} \frac{x^3}{e^x - 1} dx \right]$$

Using (1)

$$= \frac{qR}{\theta_D^3} \left[4T^3 \cdot \int_0^{\frac{\theta_D}{T}} \frac{x^3}{e^x - 1} dx + T^4 \cdot \frac{\left(\frac{\theta_D}{T}\right)^3}{e^{\frac{\theta_D}{T}} - 1} \cdot \frac{-\theta_D}{T^2} \right]$$

$$= qR \left[4 \left(\frac{T}{\theta_D}\right)^3 \int_0^{\frac{\theta_D}{T}} \frac{x^3}{e^x - 1} dx - \left(\frac{\theta_D}{T}\right) \left(\frac{1}{e^{\frac{\theta_D}{T}} - 1}\right) \right]$$

At $\boxed{T \gg \theta_D}$, we have

$$\frac{\theta_D}{T} \ll 1 \quad e^{\frac{\theta_D}{T}} = 1 + \left(\frac{\theta_D}{T}\right)$$

$$\begin{aligned} \Rightarrow C_v &\approx \frac{qR}{\theta_D^3} \left[4 \left(\frac{T}{\theta_D}\right)^3 \int_0^{\frac{\theta_D}{T}} x^2 dx - 1 \right] \\ &= qR \left[\frac{4}{3} \cdot \frac{T^3}{\theta_D^3} \cdot \left(\frac{\theta_D}{T}\right)^3 - 1 \right] \\ &= qR \cdot \frac{1}{3} = \underline{\underline{\frac{3R}{3}}} \end{aligned}$$

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Hence at higher temperature, it
converges to dubong & petit law.

At $T \ll \theta_D$

$$\frac{\theta_D}{T} \gg 1$$

$$\Rightarrow U \approx 9R \left[4 \left(\frac{T}{\theta_D} \right)^3 \cdot \int_0^{\infty} \frac{x^3}{e^x - 1} dx - \frac{1}{\infty} \right]$$

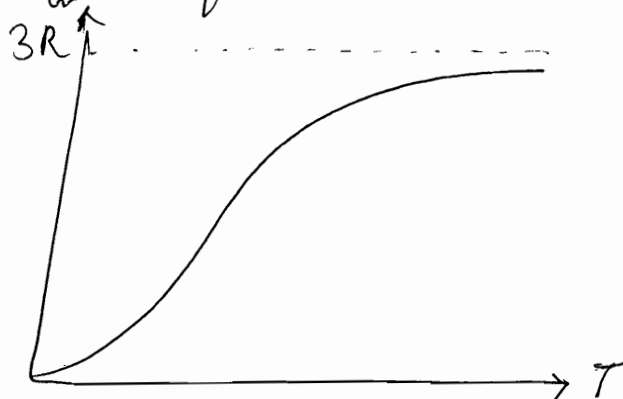
$$= 9R \cdot 4 \cdot \frac{T^3}{\theta_D^3} \cdot \lambda$$

$$= \left(\frac{36 R \lambda}{\theta_D^3} \right) \cdot T^3$$

$$\approx AT^3$$

Hence at lower temperature, the
expression converges to T^3 law

$\therefore$ for both high and low temperature,
debye's theory is concomitant with the
observed facts.



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However there are some shortcomings in the
theory :

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- ✓ ① electronic contribution is not taken
- ✓ ② There is no theoretical prediction of θ_D , and moreover it is not const. with temperature, but varies widely with temperature, except at very low temperature.
- ✓ ③ The theory is valid only for isotropic solids.

Q6) Given $C(T) = aT + bT^2$

$$\begin{aligned}\text{Avg. specific heat} &= \frac{\int_0^T C(T) dT}{T} \\ &= \frac{\int_0^T (aT + bT^2) dT}{T} \\ &= \frac{\left(a \frac{T^2}{2} + b \frac{T^3}{3} \right)}{T} \\ &= \left(\frac{aT}{2} \right) + \left(\frac{bT^2}{3} \right)\end{aligned}$$

At mid temperature $\frac{T}{2}$, $C\left(\frac{T}{2}\right) = \left(\frac{aT}{2}\right) + \left(\frac{bT^2}{4}\right)$

$\therefore$ avg. specific heat is more than
mid point specific heat

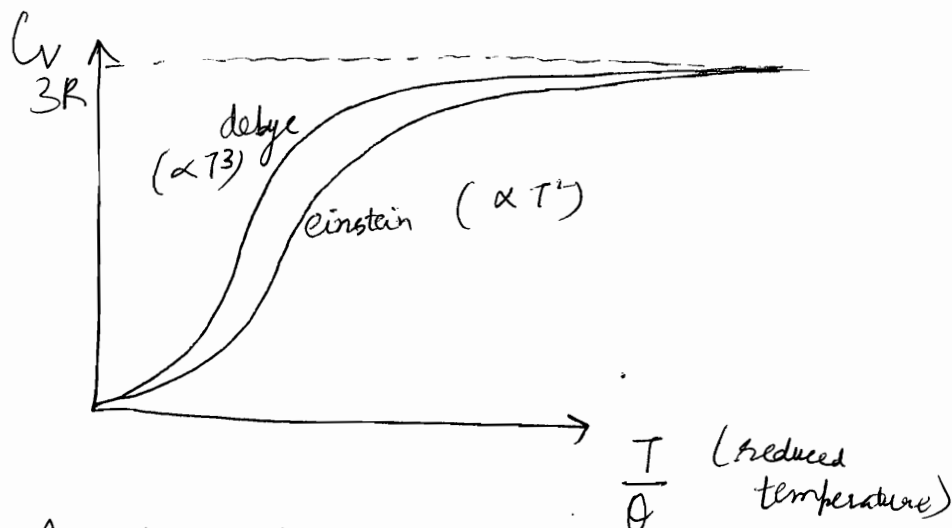
Q7)

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Q8)

Q-2)

Q4)



- At higher temperature both the curves converge at $3R$, which is the Dulong & Petit law for higher temperature.
- For lower temperature Einstein specific heat drops exponentially while the Debye's specific heat varies cubically, T^3 .

"Law of corresponding states" says that while the plots of C_v versus T may be different for different metals, however, if we plot C_v versus reduced temperature $\frac{T}{\Theta}$ [$\frac{T}{\Theta_E}$: Einstein and $\frac{T}{\Theta_D}$: Debye], the graph is strikingly similar for almost all metals. This is due to the fact that both Debye and Einstein specific heats are of form $C_v = f\left(\frac{T}{\Theta}\right)$.

2 different solids at different temperature T_1, T_2 but different reduced temperature $\frac{T_1}{\Theta_1} = \frac{T_2}{\Theta_2}$, will have same C_v .

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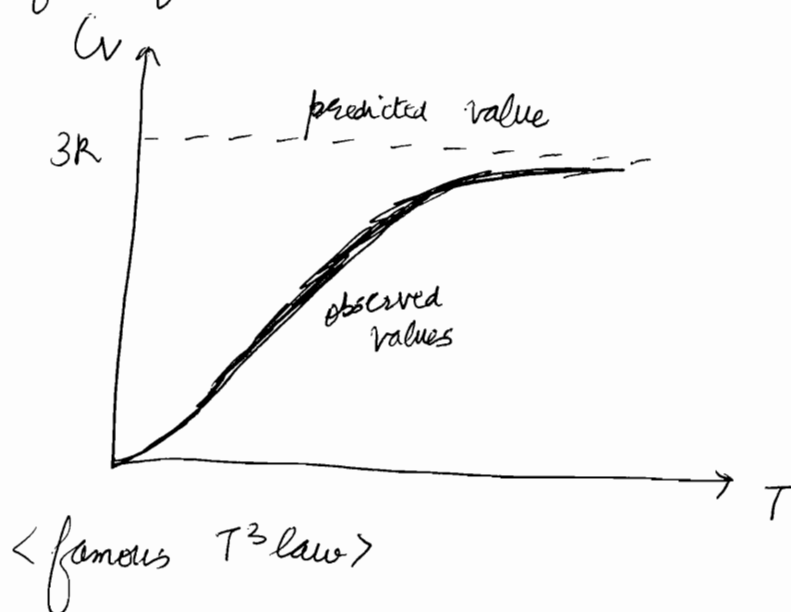
Q9)

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Dulong and Petit's law states that molar specific heat of all solids is given by $C_v = 3R$

where R is universal gas constant

While Dulong and Petit's law agrees with experimental values at high temperatures it fails for lower temperatures.



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Q10)

<Q27

<Q5>

Significance of Debye's Temperature

① give upper bound to the frequency spectrum as $\omega_D = \frac{h\gamma_m}{k} \Rightarrow \gamma_m = \left(\frac{\omega_D k}{h} \right)$

② Law of Correspondence states, curves of C_v vs T for different solids are different by a single curve of C_v vs $\left(\frac{T}{\omega_D} \right)$ fits for all solids.

Q1)

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For a maxwell distribution

Occupation index $f(\epsilon) = A e^{-\left(\frac{\epsilon}{kT}\right)}$;

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(60 marker) density of states $g(\epsilon) = \frac{2\pi V}{h^3} \cdot (2m)^{\frac{3}{2}} \epsilon^{\frac{1}{2}}$

$\Rightarrow$ No. of particles with energy between ϵ and $d\epsilon$, $n(\epsilon) d\epsilon$ is given by

$$\begin{aligned} n(\epsilon) d\epsilon &= f(\epsilon) g(\epsilon) d\epsilon \\ &= A e^{-\frac{\epsilon}{kT}} \frac{2\pi V}{h^3} \cdot (2m)^{\frac{3}{2}} \epsilon^{\frac{1}{2}} d\epsilon \\ &= B \epsilon^{\frac{1}{2}} e^{-\frac{\epsilon}{kT}} d\epsilon \end{aligned}$$

(where B is a const)

But total no. of particles = N

$$\begin{aligned} \Rightarrow N &= \int_0^{\infty} n(\epsilon) d\epsilon \\ &= \int_0^{\infty} B \epsilon^{\frac{1}{2}} e^{-\frac{\epsilon}{kT}} d\epsilon \end{aligned}$$

Put $\frac{\epsilon}{kT} = x$

$$\begin{aligned} &= B (kT)^{\frac{3}{2}} \int_0^{\infty} x^{\frac{1}{2}} e^{-x} dx \\ &= B (kT)^{\frac{3}{2}} \Gamma\left(\frac{3}{2}\right) \\ &= B \cdot (kT)^{\frac{3}{2}} \cdot \frac{1}{2} \sqrt{\pi} \end{aligned}$$

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$$\Rightarrow B = \frac{2N}{(kT)^{3/2}\sqrt{\pi}} = \frac{2\pi N}{(\pi kT)^{3/2}}$$

$$\Rightarrow n(\epsilon) d\epsilon = \frac{2\pi N}{(\pi kT)^{3/2}} \epsilon^{1/2} e^{-\frac{\epsilon}{kT}} d\epsilon$$

Now for non relativistic case, the energy of molecules in motion with velocity v is given as

$$\epsilon = \frac{1}{2} m v^2$$

Hence no. of particles with velocity between v and $v+dv$ is given as

$$n(v) dv = \frac{2\pi N}{(\pi kT)^{3/2}} \left(\frac{1}{2} m v^2\right)^{1/2} \cdot e^{-\left(\frac{m v^2}{2kT}\right)} m v dv$$

$$n(v) dv = NC v^2 e^{-\left(\frac{m v^2}{2kT}\right)} dv$$

where C is a const. $C = \frac{2\pi}{(\pi kT)^{3/2}} \frac{m^{3/2}}{\sqrt{2}}$

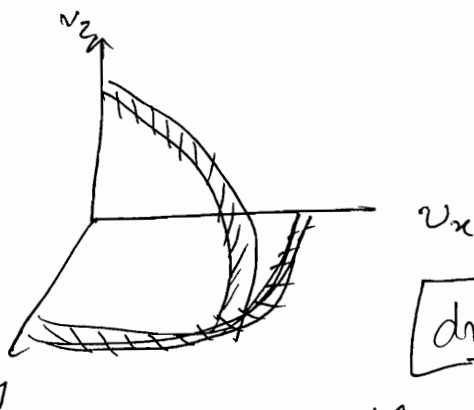
now in a phase space, we can consider velocity v to $v+dv$ can be broken into components v_x to (v_x+dv_x) , v_y to (v_y+dv_y) and v_z to (v_z+dv_z)

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$$\Rightarrow P(v) dv = \frac{n(v) dv}{N} = C v^2 e^{-\frac{mv^2}{2kT}} dv$$

$$\begin{aligned} \text{Also } P(v) dv &= P(v_x) dv_x P(v_y) dv_y P(v_z) dv_z \\ &= P(v_x) P(v_y) P(v_z) dv_x dv_y dv_z \end{aligned}$$



$$dv_x dv_y dv_z = 4\pi v^2 dv$$

$\Rightarrow$ Replace $4\pi v^2 dv$ by $dv_x dv_y dv_z$

Also since Probability of v_x, v_y, v_z are same $\Rightarrow P(v_x) = P(v_y) = P(v_z)$

Now we can write

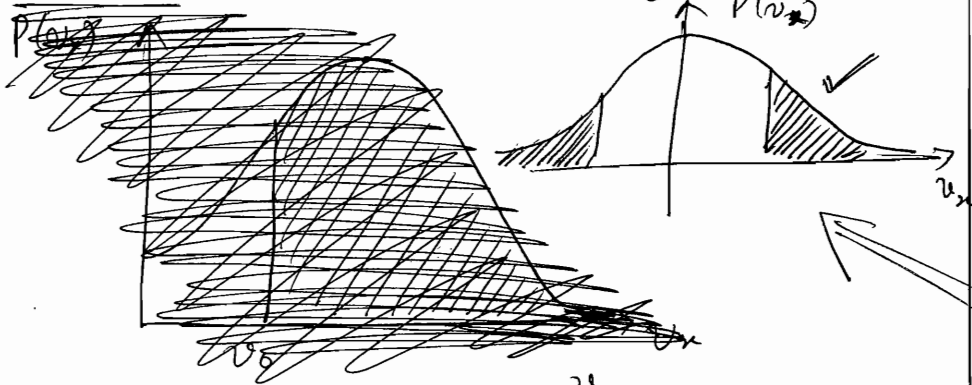
$$\begin{aligned} P(v) dv &= \left[\sqrt{\frac{m}{2\pi kT}} e^{-\frac{mv_x^2}{2kT}} \right] \left[\sqrt{\frac{m}{2\pi kT}} e^{-\frac{mv_y^2}{2kT}} \right] \left[\sqrt{\frac{m}{2\pi kT}} e^{-\frac{mv_z^2}{2kT}} \right] 4\pi v^2 dv \\ &= \left[\frac{1}{2\sqrt{\pi}} \cdot \frac{m^{3/2}}{(kT)^{3/2}} \right] e^{-\frac{mv^2}{2kT}} 4\pi v^2 dv \\ &= \frac{m}{\sqrt{2\pi kT}} e^{-\frac{mv^2}{2kT}} \end{aligned}$$

$$\begin{aligned} \Rightarrow P(v_x) dv_x &= \sqrt{\frac{m}{2\pi kT}} e^{-\frac{mv_x^2}{2kT}} dv_x \\ &= D e^{-\frac{mv_x^2}{2kT}} dv_x \end{aligned}$$

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$$\text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-u^2} du$$

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important
curve

$$P(v_x > v_0) = \frac{1 - 2 \int_0^{v_0} P(v_x) dv_x}{2}$$

in +ve direction

$$= \frac{1}{2} \left[1 - 2 \int_0^{v_0} \frac{\sqrt{m}}{\sqrt{2\pi kT}} e^{-\frac{mv_x^2}{2kT}} dv_x \right]$$

Put $\sqrt{\frac{mv_x^2}{2kT}} = x$

$$\Rightarrow dv_x = \sqrt{\frac{2kT}{m}} \cdot dx$$

$$= \frac{1}{2} \left[1 - 2 \int_0^{\sqrt{\frac{mv_0^2}{2kT}}} \sqrt{\frac{m}{2\pi kT}} \sqrt{\frac{2kT}{m}} e^{-x^2} dx \right]$$

$$= \frac{1}{2} \left[1 - \frac{2}{\sqrt{\pi}} \int_0^{\sqrt{\frac{mv_0^2}{2kT}}} e^{-x^2} dx \right]$$

$$P(v_x > v_0) = \frac{1}{2} \left[1 - \text{erf} \left(\sqrt{\frac{mv_0^2}{2kT}} \right) \right]$$

Q2)

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r.m.s. velocity for a gas molecule, v_{rms}
is given as

$$v_{rms} = \sqrt{\frac{3kT}{m}} \quad - (1)$$

Escape velocity for a particle be v_{esc}

$$\Rightarrow \frac{1}{2} m v_{esc}^2 = + \frac{GMm}{R}$$

$$\Rightarrow v_{esc} = \sqrt{\frac{2GM}{R}}$$

We also know $mg = \frac{GMm}{R^2}$

$$\Rightarrow GM = gR^2$$

$$\Rightarrow v_{esc} = \sqrt{\frac{2gR^2}{R}} = \sqrt{2gR} \quad - (2)$$

Since $v_{esc} = v_{rms}$

$$\Rightarrow \sqrt{2gR} = \sqrt{\frac{3kT}{m}}$$

$$\Rightarrow T = \left(\frac{2gR \cdot m}{3k} \right)$$

$$= \frac{2 \times 9.8 \times 6370 \times 10^3 \times 23.24 \times 10^{-27} \times 2}{3 \times 1.38 \times 10^{-23}}$$

$$= \boxed{1.4 \times 10^5 \text{ K}}$$

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do not
forget to
convert g_m
to kg.



$$\eta = \frac{F/A}{\left(\frac{dv}{dz}\right)}$$

Q3)

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From transport phenomenon, viscosity η is given as

$$\eta = \frac{1}{3} m n \bar{v} \lambda$$

$$\bar{v} = \sqrt{\frac{8kT}{\pi m}}$$

$$\Rightarrow \lambda = \frac{3\eta}{\frac{m n}{V} \sqrt{\frac{8kT}{\pi m}}} = \frac{3\eta V \sqrt{\pi}}{N \sqrt{8kTm}}$$

$$= \frac{3\eta RT \sqrt{\pi}}{N \cdot P \sqrt{8RTm}} \quad (\text{assuming ideal gas})$$

$$\lambda = \frac{3\eta}{P} \sqrt{\frac{T \pi R}{8 N^2 m}}$$

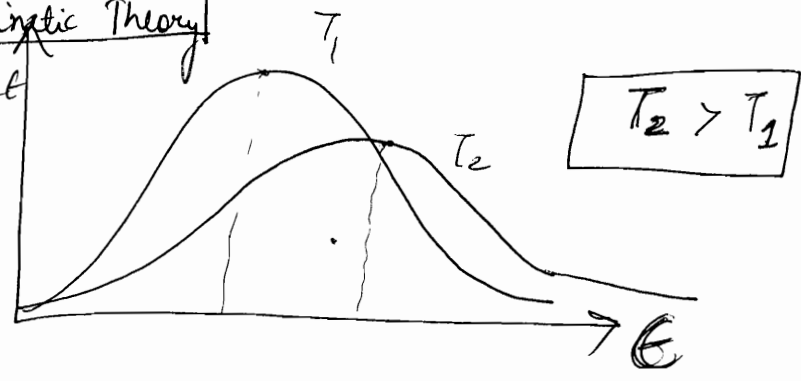
$$= \frac{3.974 \times 10^{-1} \text{ m}}{\sqrt{2}}$$

$$= 2.77 \times 10^{-7}$$

Remember,
its average speed $\bar{v}$
 $= \sqrt{\frac{8kT}{\pi m}}$

$$\frac{\text{kg m}^2}{\text{s}^2 \text{ m}^2} \left(\frac{\text{kg}}{\text{m}^3} \right)$$

Q4) From kinetic Theory



From kinetic Model

$$PV = \frac{1}{3} m n v^2 = \frac{2}{3} \frac{1}{2} m n v^2$$

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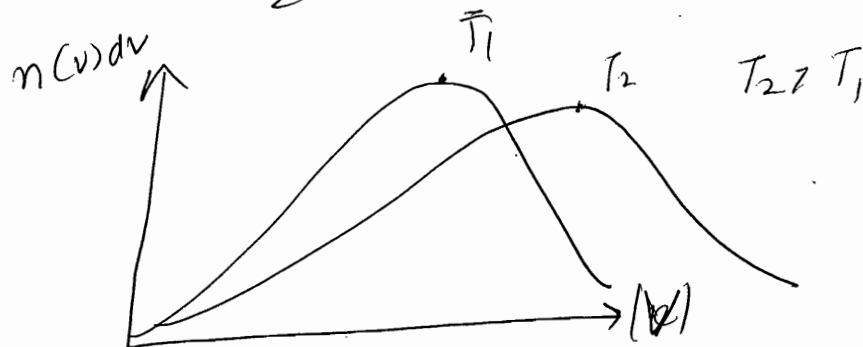
$$= \frac{2}{3} E$$

$$\Rightarrow E = \frac{3}{2} PV = \frac{3}{2} RT \quad (\text{from ideal gas})$$

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$\therefore$ more the temperature more the energy
Hence the curve moves towards right.

$$\text{Also } E = \frac{1}{2} mv^2$$



Also the area under the curve is
equal to number of particles, hence
constant, therefore peak shifts
downwards also on increasing Temp.

From distribution

The peak in the Maxwell distribution
curve corresponds to the most probable
speed. Now,

$$n(v) dv = N \left(\sqrt{\frac{m}{2\pi kT}} \right)^3 4\pi v^2 e^{-\frac{mv^2}{2kT}} dv$$

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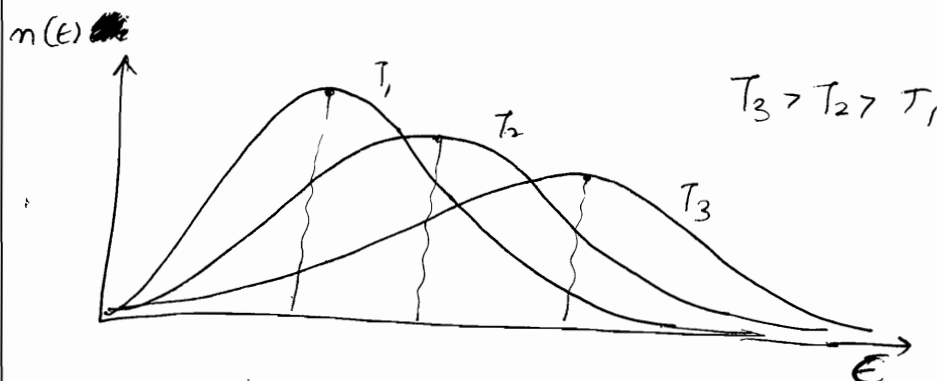
For most probable speed, $\frac{d n(v)}{d v} = 0$

$$\Rightarrow \left[2v e^{-\frac{mv^2}{2kT}} + v^2 e^{-\frac{mv^2}{2kT}} \left(-\frac{2mv}{2kT} \right) \right] = 0$$

$$\Rightarrow 1 = v^2 \cdot \frac{m}{2kT} \Rightarrow v_{M.P.} = \sqrt{\frac{2kT}{m}}$$

$$\Rightarrow v_{M.P.} \propto \sqrt{T}$$

When temperature is increase
 $v_{M.P.}$ also increases, hence curve
shifts rightwards.



Q5)

$$\lambda = \frac{1}{\sqrt{2} \pi d^2 n}$$

Mean free path, λ , is defined as the
average distance after which a
molecules undergoes collision!!

Assuming ideal gas

$$n = \frac{N}{V} = \frac{NP}{RT}$$

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$$\Rightarrow \lambda = \frac{RT}{\sqrt{2} \pi (2 \times 10^{-10})^2 \times N \times P}$$

$$= \underline{\underline{2.27 \times 10^{-7} \text{ m}}}$$

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Q6)

<Q1> <Q4>

$$n(v) dv = N \left(\sqrt{\frac{m}{2\pi kT}} \right)^3 4\pi v^2 e^{-\frac{mv^2}{2kT}} dv$$

$$= A v^2 e^{-\frac{mv^2}{2kT}} dv$$

where $A = N \left(\sqrt{\frac{m}{2\pi kT}} \right)^3 4\pi$

$$\Rightarrow \langle v^3 \rangle = \frac{\int_0^\infty v^3 dn}{\int_0^\infty dn}$$

$$= \frac{\int_0^\infty v^5 e^{-\frac{mv^2}{2kT}} dv}{\int_0^\infty v^2 e^{-\frac{mv^2}{2kT}} dv}$$



$$= \frac{1}{2 \left(\frac{m}{2kT} \right)^{\frac{6}{2}}} \cdot \sqrt{\frac{6}{2}}$$

$$\frac{1}{2 \left(\frac{m}{2kT} \right)^{\frac{3}{2}}} \cdot \sqrt{\frac{3}{2}}$$

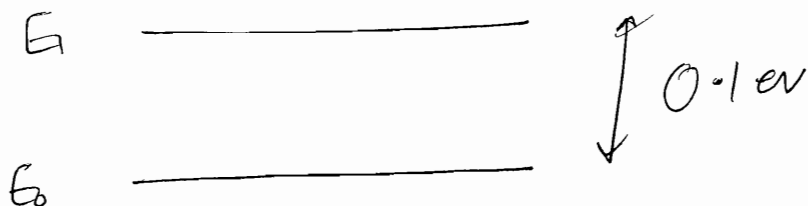
$$= \frac{1}{\left(\frac{m}{2kT} \right)^{\frac{3}{2}}} \cdot \frac{2}{\frac{1}{2} \cdot \sqrt{\pi}}$$

do not
use denomina
as N.
This will
avoid
confusion
with constants
like A

$$\langle C^3 \rangle = \frac{4}{\sqrt{\pi}} \cdot \left(\frac{2kT}{m} \right)^{\frac{3}{2}}$$

Q7)

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Assuming that the assembly follows
Maxwell distribution,

let no. of particle in '0' state be n_0

$\Rightarrow$ no. of particles in '1' state, n_1 , is
given by

$$n_1 = n_0 e^{-\frac{\Delta E}{kT}} = n_0 e^{-\left(\frac{0.1 \text{ eV}}{kT}\right)}$$

Also $n_0 + n_1 = N$

$$\Rightarrow n_0 \left[1 + e^{-\frac{\Delta E}{kT}} \right] = N$$

$$\Rightarrow n_0 = \frac{N}{1 + e^{-\left(\frac{\Delta E}{kT}\right)}}$$

$$\Rightarrow n_1 = \frac{N}{1 + e^{-\frac{\Delta E}{kT}}} e^{-\frac{\Delta E}{kT}}$$

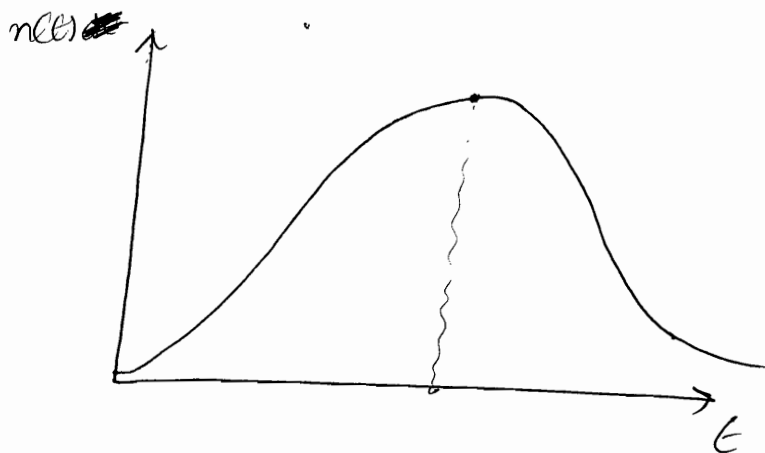
$$\Rightarrow \left(\frac{n_1}{N} \right) = \frac{e^{-\left(\frac{\Delta E}{kT}\right)}}{1 + e^{-\left(\frac{\Delta E}{kT}\right)}}$$

$$e^{-\frac{\Delta E}{kT}} = e^{-\frac{0.1 \times 1.6 \times 10^{-19}}{1.38 \times 10^{-23} \times 300}} = 0.02$$

$$\Rightarrow \frac{n_1}{N} = \frac{0.02}{1.02} \approx 0.02 = 2\%$$

Q8)

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$$n(E) dE = \frac{2\pi N}{(\pi kT)^{3/2}} E^{1/2} e^{-\frac{E}{kT}} dE$$

This distribution peaks for $\left. \frac{dn}{dE} \right|_{E=E_0} = 0$

$$\Rightarrow \frac{1}{2} E^{-1/2} \cdot e^{-\frac{E}{kT}} + E^{1/2} e^{-\frac{E}{kT}} \left(-\frac{1}{kT} \right) = 0$$

$$\Rightarrow \frac{1}{2\sqrt{E}} = \frac{\sqrt{E}}{kT} \Rightarrow \boxed{E = \frac{kT}{2}}$$

Do not forget to mention that at $E = \frac{1}{2} kT$, $\frac{d^2 n}{dE^2} < 0$

The mean energy of molecules is given as

$$\bar{E} = \frac{\int E \, dn}{\int dn} = \frac{\int E^{3/2} e^{-\frac{E}{kT}} dE}{\int E^{1/2} e^{-\frac{E}{kT}} dE}$$

Put $\frac{E}{kT} = x$, we have

$$\bar{E} = \frac{(kT)^{5/2} \int x^{3/2} \cdot e^{-x} dx}{(kT)^{3/2} \int x^{1/2} e^{-x} dx} = (kT) \frac{\int x^{3/2} e^{-x} dx}{\int x^{1/2} e^{-x} dx}$$

$$\frac{\int x^{3/2} e^{-x} dx}{\int x^{1/2} e^{-x} dx} = \frac{5}{2} = \frac{3}{2} kT$$

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$$\bar{E} = \frac{3}{2} \times \frac{1.38 \times 10^{-23} \times 300}{1.6 \times 10^{-19}} \text{ eV}$$

$$= \underline{\underline{0.0388 \text{ eV}}}$$

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Q9) Average speed $\bar{v}$ for a molecule is
given as

$$\bar{v} = \sqrt{\frac{8}{\pi} \frac{kT}{m}}$$

$$\Rightarrow \left(\frac{\bar{v}_1}{\bar{v}_2} \right) = \frac{\sqrt{\frac{T_1}{m_1}}}{\sqrt{\frac{T_2}{m_2}}}$$

if $\bar{v}_o = \bar{v}_H$

$$\Rightarrow \frac{T_o}{m_o} = \frac{T_H}{m_H}$$

$$\Rightarrow T_H = T_o \cdot \frac{m_H}{m_o} = 350 \times \frac{2}{32}$$

$$= \underline{\underline{21.875 \text{ K}}}$$

Q10) using $\frac{dN_{nc}}{N} = - \frac{dx}{\lambda}$

N.C. : non collided

$$\Rightarrow \boxed{N = N_o e^{-\left(\frac{x}{\lambda}\right)}}$$

for $x = 2\lambda$

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$$N = N_0 e^{-2} = \frac{6 \times 10^3}{e^2} = \underline{\underline{812}}$$

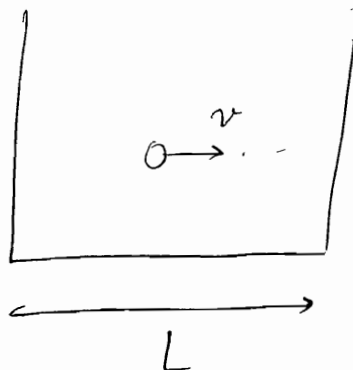
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Q11)

$$\Delta p = 2 m v_x$$

$$\Delta t = \left(\frac{2L}{v} \right)$$

$$\Rightarrow \frac{\Delta p}{\Delta t} = \left(\frac{m v_x^2}{L} \right)$$



$$\Rightarrow P = \frac{F}{A} = \frac{m v_x^2}{L}$$

$$\text{Total Pressure by all the molecules} = \frac{N m v_x^2}{L}$$

$$= \left(\frac{m N v_x^2}{V} \right)$$

$$\text{Now } v_{rms}^2 = v_x^2 + v_y^2 + v_z^2 = 3 v_x^2$$

$$\Rightarrow v_x^2 = \frac{v^2}{3} \quad \& \quad \frac{N}{V} = n \Rightarrow$$

$$\Rightarrow P = \frac{1}{3} n m v_{rms}^2$$

$$\text{Assuming } \langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle$$

$$\langle v_x^2 \rangle = \frac{\langle v^2 \rangle}{3} = \left(\frac{v_{rms}^2}{3} \right)$$

From Boltzmann distribution,
No. of molecules with velocity between
 v and $v+dv$

$$n dv = N \left(\frac{m}{2\pi kT} \right)^{3/2} 4\pi v^2 e^{-\frac{mv^2}{2kT}} dv$$

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$$\begin{aligned}
 \langle v^2 \rangle &= \frac{\int v^2 n(v) dv}{\int n(v) dv} \\
 &= \frac{\int v^4 e^{-\frac{mv^2}{2kT}} dv}{\int v^2 e^{-\frac{mv^2}{2kT}} dv} \\
 &= \frac{1}{\left(\frac{m}{2kT}\right)} \cdot \frac{\sqrt{5}}{\sqrt{3}} \\
 &= \frac{3}{2} \cdot \frac{2kT}{m} = \left(\frac{3kT}{m}\right)
 \end{aligned}$$

$$\Rightarrow P = \frac{1}{3} m n \frac{3kT}{m}$$

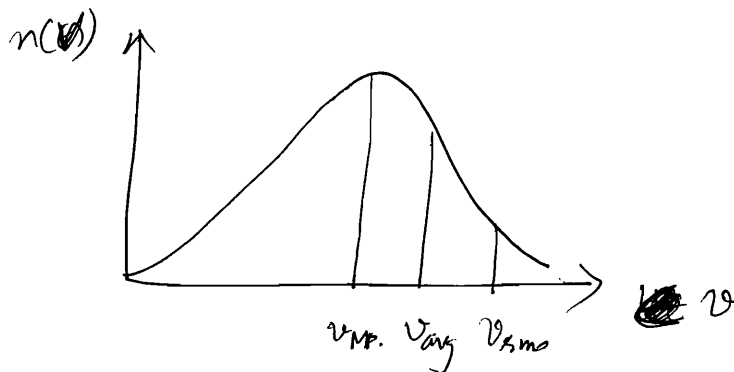
$$P = nkT$$

Q12)

(Q1)

Q13)

From the Maxwell distribution, we get



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$$v_{M.P.} = \sqrt{\frac{2RT}{M}}$$

$$v_{avg} = \sqrt{\frac{8RT}{\pi M}}$$

$$v_{rms} = \sqrt{\frac{3RT}{M}}$$

$$\Rightarrow v_{M.P.} = \sqrt{\frac{2 \times 8.314 \times 273}{2 \times 10^{-3}}} = 1506.5 \text{ m/s}$$

$$v_{avg} = \sqrt{\frac{8}{\pi \cdot 2}} \quad v_{M.P.} = 1700 \text{ m/s}$$

$$v_{rms} = \sqrt{\frac{3}{2}} v_{M.P.} = 1845.07 \text{ m/s}$$

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Q14) X

Q15) Energy of a particle is, in general,
a function of a number of parameters.
Each of these parameters is called
degree of freedom. If energy E ,
associated with such a particle can
be described in terms of a parameter
such that

$$E = aZ^2$$

where $a = \text{const.}$, then the average
energy per particle associated with
this parameter Z , will be $\left(\frac{1}{2}kT\right)$ at equilibrium
temperature T .

This is the law of equipartition of energy

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Proof :

Let the particles follow Maxwell
distribution. i.e.

$$n(\epsilon) d\epsilon = \frac{2\pi N}{(\pi kT)^{3/2}} \epsilon^{1/2} e^{-\frac{\epsilon}{kT}} d\epsilon$$

$$= A \epsilon^{1/2} e^{-(\epsilon/kT)} d\epsilon$$

Now if $\epsilon = az^2$

$$\Rightarrow \langle \epsilon \rangle = \frac{\int \epsilon n(\epsilon) d\epsilon}{\int n(\epsilon) d\epsilon}$$

$$= \frac{\int az^2 \cdot (az^2)^{1/2} e^{-\frac{az^2}{kT}} \cdot 2az dz}{\int (az^2)^{1/2} \cdot e^{-\frac{az^2}{kT}} \cdot 2az dz}$$

Cancelling out common terms,

$$= \frac{\int z^4 \cdot e^{-\frac{az^2}{kT}} dz}{\int z^2 e^{-\frac{az^2}{kT}} dz}$$

$$= \frac{1}{\left(\frac{a}{kT}\right)} \cdot \frac{\sqrt{\frac{5}{2}}}{\sqrt{\frac{3}{2}}} =$$

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Proof

Let the particle follow Maxwell Distⁿ

$$n(z) dz = A e^{-\frac{E(z)}{kT}} dz$$

Now if $E = az^2$

$$\langle E \rangle = \frac{\int E n(E) dE}{\int n(E) dE}$$

$$= \frac{\int az^2 \cdot e^{-\frac{az^2}{kT}} dz}{\int e^{-\frac{az^2}{kT}} dz}$$

$$= a \cdot \frac{1}{\left(\frac{a}{kT}\right)} \frac{\Gamma^{3/2}}{\Gamma^{1/2}}$$

$$= \underline{\underline{\frac{1}{2} kT}}$$

Shortcomings:

①

Only for energy that can be expressed
as az^2 i.e. not even gravitation
energy $\left(\frac{GMm}{z}\right)$

②

For lower temperatures, it does not yield

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correct results as quantum effect
take over.

eg. Dulong & Petit law fails for
low temperatures

For other ^{higher} powers integrals do not converge...

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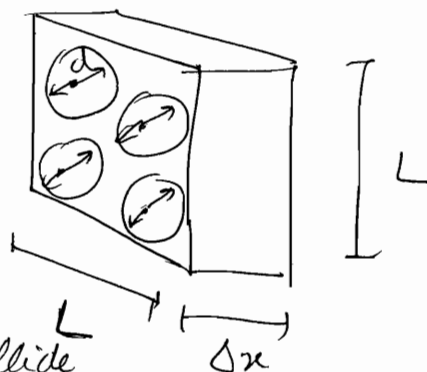
Q16)

<Q4> <Q11>

Q17)

Consider a thin layer of gas of
dimension L, L and Δx as shown.

Let a large number of
molecules N be
projected towards the
layer



Let ΔN molecules collide
with the molecules of the
layer. Let n' be the no. of molecules
per unit volume, and d be their
diameter. On average ratio of
no. of collisions ΔN to the no. of
projected molecules N is equal to the
ratio of the target area of molecules
in the layer to the total area of
face of the layer.

$$\therefore \frac{\Delta N}{N} = \frac{(n' L^2 \Delta x) \pi d^2}{L^2}$$

$$4\pi r^2 = \pi d^2$$

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$$= n' \pi d^2 \Delta x$$

In the limit of very small Δx ,

$$\left(\frac{dN}{N}\right) = + n' \pi d^2 dx$$

$$\text{define } \lambda = \left(\frac{1}{n' \pi d^2}\right)$$

$$\Rightarrow \left(\frac{dN_c}{N}\right) = \left(\frac{dx}{\lambda}\right)$$

N_c : collided molecules

$$\Rightarrow \lambda = \frac{dx}{\left(\frac{dN}{N}\right)}$$

i.e. λ is the ratio of distance travelled by the molecules to the collision probability $\left(\frac{dN}{N}\right)$; λ is called mean free path.

Other theories give correction in mean free path as,

$$\text{Clausius : } \lambda = \frac{3}{4} \left(\frac{1}{n' \pi d^2}\right)$$

$$\text{Maxwell : } \lambda = \frac{1}{\sqrt{2}} \frac{1}{n' \pi d^2}$$

Using Maxwell's result,

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$$\lambda = \frac{1}{\sqrt{2}} \frac{V}{n N_A} \cdot \frac{1}{\pi d^2}$$

$$= \frac{1}{\sqrt{2}} \frac{1}{\pi d^2} \frac{RT}{P} \cdot \frac{1}{N_A}$$

$$\Rightarrow N_A = \left(\frac{RT}{\sqrt{2} \pi d^2 P \lambda} \right)$$

$$N_A = 6.64 \times 10^{23}$$

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Q18)

$$10^5 * 10^{-3} = n \cdot 8.314 \times 293$$

$$\Rightarrow n = 0.041$$

(from ideal
gas equation)

$$\Rightarrow 0.041 \times 6.023 \times 10^{23} \times m_{H_2} = 0.09 \times 10^{-3}$$

$$\Rightarrow m_{H_2} = 3.64 \times 10^{-27} \text{ kg}$$

From Maxwell distribution, we know

$$v_{RMS} = \sqrt{\frac{3kT}{m}} \quad v_{MP} = \sqrt{\frac{2kT}{m}} \quad \bar{v} = \sqrt{\frac{8kT}{\pi m}}$$

$$\Rightarrow v_{RMS} = \sqrt{\frac{3 \times 1.38 \times 10^{-23} \times 300}{3.64 \times 10^{-27}}} = \underline{\underline{1847 \text{ m/s}}}$$

$$\bar{v} = \sqrt{\frac{8 \times 1}{\pi \times 3}} \quad v_{RMS} = \underline{\underline{1702.27 \text{ m/s}}}$$

$$v_{MP} = \sqrt{\frac{2}{3}} \times v_{RMS} = \underline{\underline{1508.06 \text{ m/s}}}$$

NTP (1 atm, 20°C)

STP (1 atm, 0°C)

Q19)

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$$v_{RMS} = \sqrt{\frac{3kT}{m}}$$

$$\frac{v_{Ar}}{v_{O_2}} = \frac{\sqrt{\frac{T_{Ar}}{m_{Ar}}}}{\sqrt{\frac{T_{O_2}}{m_{O_2}}}} = \frac{\sqrt{\frac{313}{40}}}{\sqrt{\frac{273}{32}}} = 0.95$$

$$\Rightarrow v_{Ar} = 460 \times 0.95 = \underline{\underline{440.54 \text{ m/s}}}$$

$$\frac{v_{T_2}}{v_{T_1}} = \sqrt{\frac{T_2}{T_1}} \Rightarrow 2 = \sqrt{\frac{T}{293}}$$

$$\Rightarrow T = 1172 \text{ K}$$

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Q20) <Q-167

Q21) <Q-1>

Q22) Assuming Maxwell Boltzmann distⁿ $2kT$ ————— N_2
and density of state gles to be equal kT ————— N_1
for all levels, ————— N_0

$$\Rightarrow N_1 = N_0 e^{-\frac{kT}{kT}} = \left(\frac{N_0}{e} \right)$$

$$N_2 = N_0 e^{-\frac{2kT}{kT}} = \left(\frac{N_0}{e^2} \right)$$

$$\Rightarrow N_0 \left[1 + \frac{1}{e} + \frac{1}{e^2} \right] = N \Rightarrow N_0 = \left(\frac{N}{1.5} \right) = 0.66 N$$

$$= \underline{\underline{66}}$$

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$$N_1 = N_0 e^{-1} = 25$$

$$N_2 = N_0 e^{-2} = 9$$

$$\Rightarrow E_{\text{Total}} = E_0 N_0 + E_1 N_1 + E_2 N_2$$

$$= 25 kT + 9 * 2 kT$$

$$E = 43 kT$$

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Q23)

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Let the no. of degree of freedom of
gas be f

$$\Rightarrow \text{Internal Energy} = f * \frac{1}{2} kT = \frac{f}{2} kT$$

$$\Rightarrow dU = \frac{f}{2} k dT$$

But we know $dU = C_V dT$

$$\Rightarrow C_V = \frac{f}{2} k \quad \text{for 1 particle}$$

For 1 mole,

$$C_V = \frac{f}{2} k \cdot N_A = \frac{f}{2} R$$

Also for ideal gas

$$C_p - C_V = R$$

$$\Rightarrow C_p = R + C_V = \left(\frac{f+2}{2} \right) R$$

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$$\Rightarrow \gamma = \frac{C_p}{C_v} = \frac{(f+2)}{f}$$

For diatomic, $f = 5 \Rightarrow \gamma = 7/5$

For diatomic, $f = 6 \Rightarrow \gamma = 8/6 = 4/3$

Mention : diatomic $\rightarrow$ 3 D.O.F. due to
motion of COM

+
2 DOF due to
rotation along axis $\perp$ to
internuclear axis.

triatomic $\rightarrow$ 3 DOF for motion
of COM

+
3 DOF due to
rotation of molecule
along 3 principal axes

(Q24) Let number of particle in $(-E)$
state be N_1 and total number of
particles in (E) state be N_2

Now ; $N_1 + N_2 = N$ (fixed)

$-EN_1 + EN_2 = E$ (fixed) ... given

$\Rightarrow N_1 - N_2 = -\frac{E}{E}$

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$$\Rightarrow N_1 = \left(\frac{N - \frac{E}{\epsilon}}{2} \right)$$

$$N_2 = \left(\frac{N + \frac{E}{\epsilon}}{2} \right)$$

Entropy of Microcanonical ensemble S ,
is defined as

$$S = k \ln \Omega$$

where Ω is thermodynamic probability
which in this case is,

$$\begin{aligned} N_{C_{N_1}} &= N_{C_{N_2}} = \frac{N!}{N_1! N_2!} \\ &= \frac{N!}{\left(\frac{N - \frac{E}{\epsilon}}{2} \right)! \left(\frac{N + \frac{E}{\epsilon}}{2} \right)!} \end{aligned}$$

$$\Rightarrow S = k \left[\ln N! \right] - k \ln \left[\left(\frac{N - \frac{E}{\epsilon}}{2} \right)! \right] - k \left[\ln \left(\frac{N + \frac{E}{\epsilon}}{2} \right)! \right]$$

By Sterling approximation

$$\ln n = n \ln n - n$$

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$$\Rightarrow S = k \ln N - N - \cancel{k \ln(N+E)} \\ - \frac{k}{2} \left(N - \frac{E}{2} \right) \ln \frac{1}{2} \left(N - \frac{E}{2} \right) \\ - \frac{k}{2} \left(N + \frac{E}{2} \right) \ln \frac{1}{2} \left(N + \frac{E}{2} \right) \\ - N$$

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Tut 26 : FD and BE statistics

Q1)

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Bose Einstein Distribution is the energy distribution of integral spin particles like photons, deuterons.

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(30 marks)

In order to find the distribution, we need to calculate the thermodynamic probability Ω , i.e. the number of possible microstates in a macrostate or a configuration

let there be $(N_1, N_2, N_3, \dots, N_N)$ particles in Energy states $(E_1, E_2, E_3, \dots, E_N)$

Since the particles are indistinguishable, there is one way to choose which $N_1, N_2, \dots$ particles in $E_1, E_2, \dots$ states.

Now we need to distribute N_i particles in g_i cells such that

- ① Particles are indistinguishable
- ② Any of the g_i cells can hold any of the N_i particles.
- ③ A priori probability of all microstates is same.

No. of such possible combinations is given by

$$\Omega_i = \frac{g_i^{N_i + g_i - 1}}{g_i^{N_i}} = \frac{N_i + g_i - 1}{N_i^{g_i - 1}}$$

$$\Rightarrow \Omega = \prod_i \Omega_i = \prod_i \frac{N_i + g_i - 1}{N_i^{g_i - 1}}$$

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Now at equilibrium entropy is maximum
 $\Rightarrow d(\ln \Omega) = 0$ $[\because ds = k(d \ln \Omega)]$

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$$\Rightarrow d \left[\sum \ln \frac{N_i + g_i - 1}{N_i} - \ln N_i - \ln (g_i - 1) \right] = 0 \quad (1)$$

Note that according to sterling approximation,

$$\ln n = n \ln n - n$$

$$\Rightarrow \underline{d(\ln \Omega)} = \left(n \cdot \frac{1}{n} + \ln n - 1 \right) dn = \underline{\ln n \, dn} \quad (2)$$

Using (2) in (1),

$$\sum \ln(N_i + g_i - 1) \, dN_i - \ln(N_i) \, dN_i = 0$$

$$\Rightarrow \sum \ln \left(\frac{N_i + g_i - 1}{N_i} \right) dN_i = 0 \quad (3)$$

$$\text{Also } \sum dN_i = N$$

$$\Rightarrow \sum dN_i = 0 \quad (4)$$

$$\text{Also } \sum E_i N_i = E$$

$$\Rightarrow \sum E_i dN_i = 0 \quad (5)$$

Multiply (4) by α , (5) by β , and perform (4) + (3) - (5), we have

$$\sum \alpha dN_i + \sum \ln \left(\frac{N_i + g_i - 1}{N_i} \right) dN_i - \sum \beta E_i dN_i = 0$$

$$\Rightarrow \alpha + \ln \left(\frac{N_i + g_i - 1}{N_i} \right) - \beta E_i = 0$$

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⇒ Approximating $N_i + g_i - 1 \approx N_i + g_i$, we have

$$\frac{N_i + g_i}{N_i} = e^{-(\alpha - \beta \epsilon_i)}$$

$$\Rightarrow \frac{g_i}{N_i} = e^{-(\alpha - \beta \epsilon_i)} - 1$$

$$\Rightarrow f(\epsilon_i) = \frac{N_i}{g_i} = \frac{1}{e^{-(\alpha - \beta \epsilon_i)} - 1}$$

Also, we know

$$S = k \ln \Omega$$

$$\Rightarrow T ds = T k d(\ln \Omega)$$

$$\Rightarrow T ds = k T (\alpha - \beta \epsilon_i) d n_i$$

$$\Rightarrow \sum \epsilon_i d n_i = 0 + k \beta T \sum \epsilon_i d n_i$$

$$\Rightarrow k \beta T = 1$$

$$\Rightarrow \boxed{\beta = 1/kT}$$

$$\Rightarrow f(\epsilon_i) = \frac{1}{e^{\left(\frac{\epsilon_i}{kT} - \alpha\right)} - 1}$$

★ μ is called
chemical potential
of assembly of
bosons

Put $\alpha = \frac{\mu}{kT}$, we have

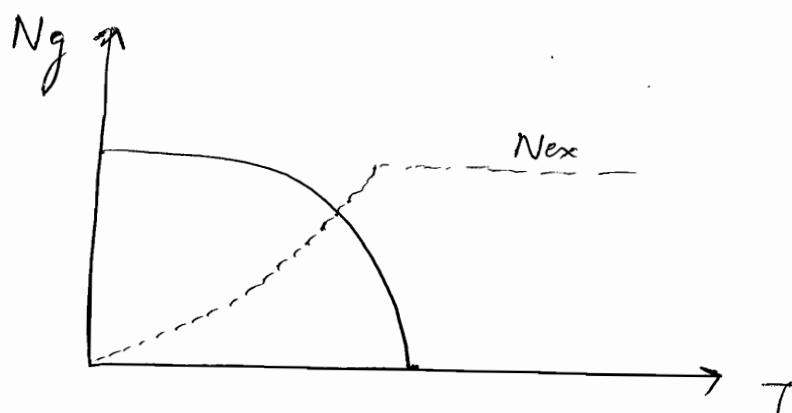
$$\boxed{f(\epsilon_i) = \frac{1}{e^{\left(\frac{\epsilon_i - \mu}{kT}\right)} - 1}$$

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(02)

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Bose Einstein Condensate is a state of matter of dilute gas of weakly interacting bosons at temperatures close to absolute zero. Under such conditions very large fractions of bosons occupy the lowest quantum state at which point quantum effects become apparent on macroscopic scale. This transition occurs below a critical temperature T_B called Condensation Temperature.



Total number of particles for temp. before T_B is given as

$$N = N_g + N_{ex}$$

where N_{ex} have energy more than 0.

For such particles from B.E. statistics

$$N_{ex} = \int f(E) g(E) dE$$

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$$= \int \frac{1}{e^{\frac{\epsilon - \mu}{kT}} - 1} \cdot \frac{2\pi V}{h^3} \cdot (2m)^{3/2} \cdot \epsilon^{1/2} d\epsilon$$

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A very interesting thing for $T = T_B$ is that
chemical potential $\mu = 0$

$$= \int \frac{1}{e^{\frac{\epsilon}{kT}} - 1} \cdot \frac{2\pi V}{h^3} \cdot (2m)^{3/2} \cdot \epsilon^{1/2} d\epsilon$$

Put $\frac{\epsilon}{kT} = x$, we have

$$N_{ex} = (kT)^{3/2} \frac{2\pi V}{h^3} (2m)^{3/2} \int_0^\infty \frac{1}{e^x - 1} x^{1/2} dx$$

Also let $\int_0^\infty \frac{x^{1/2}}{e^x - 1} dx = \lambda$

$$\Rightarrow N_{ex} = \frac{2\pi V}{h^3} (2m)^{3/2} \lambda (kT)^{3/2}$$

At $T = T_B$, $N_{ex} = N$ as $N_g = 0$

$$\Rightarrow N = \frac{2\pi V}{h^3} (2m)^{3/2} \lambda (kT_B)^{3/2}$$

$$\Rightarrow T_B = \frac{1}{k} \left[\frac{Nh^3}{(2\pi V) (2m)^{3/2} \lambda} \right]^{2/3}$$

$$T_B = \frac{h^2}{2mk} \left[\frac{n}{2\pi\lambda} \right]^{2/3}$$

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$$\text{Also, } N_{ex} = (kT)^{3/2} \cdot \frac{N}{(kT_B)^{3/2}}$$

$$= N \cdot \left(\frac{T}{T_B}\right)^{3/2}$$

$$\Rightarrow N_g = N - N_{ex}$$

$$N_g = N \left(1 - \left(\frac{T}{T_B}\right)^{3/2}\right)$$

< draw the curve >

Helium Bose Einstein Condensate

For He, $T_B = \text{~~2.17~~ } 2.17\text{K}$

At atmospheric pressure Helium condenses at 4.2K into a normal liquid. As temp is lowered further, it goes phase transition once again at 2.17K and change into another liquid of vastly different properties.

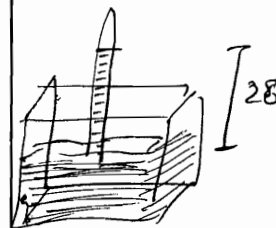
The new fluid has

- (i) Zero viscosity
- (ii) infinite thermal conductivity
- (iii) no bubbles.

Lack of viscosity means superfluid can flow through smallest of cracks without impediment. Another effect that superfluid Helium shows are that of fountain effect & 2nd sound.

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Anomalous
fluid is called
Helium-II

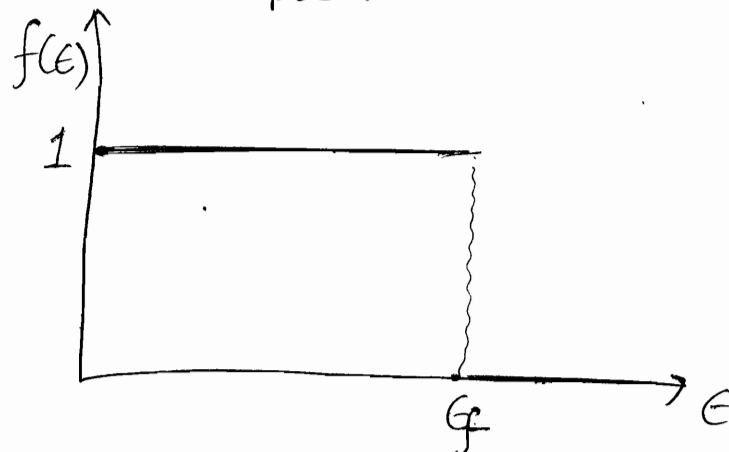


Fountain effect
in capillary
tubes.

Q3)

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Fermi Energy is defined as the energy level for which, at $T=0\text{K}$, all energy states below this energy level are fully occupied, and above which no state is occupied.



For $T=0\text{K}$

$$f(\epsilon) = \begin{cases} 1 & \epsilon < \epsilon_0 \\ 0 & \epsilon > \epsilon_0 \end{cases}$$

$$g(\epsilon) d\epsilon = \frac{2\pi V}{h^3} (2m)^{3/2} \epsilon^{1/2} d\epsilon$$

Also spin degeneracy = 2 (fermions)

$$\begin{aligned} \Rightarrow N &= \int 2f(\epsilon) g(\epsilon) d\epsilon \\ &= \int_0^{\epsilon_0} \frac{4\pi V}{h^3} (2m)^{3/2} \epsilon^{1/2} d\epsilon \\ &= \left(\frac{4\pi V}{2h^3} \right) \cdot (2m)^{3/2} \cdot \frac{\epsilon_0^{3/2}}{(3/2)} \end{aligned}$$

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$$\Rightarrow \epsilon_{f_0} = \left[\frac{3}{2} N \cdot \frac{2h^3}{4\pi V (2m)^{3/2}} \right]^{2/3}$$

Also,

$$\begin{aligned} E &= \int \epsilon_i \, dN_i \\ &= \int \epsilon \cdot \frac{2 \cdot 2\pi V}{h^3} \cdot (2m)^{3/2} \cdot \epsilon^{1/2} \, d\epsilon \\ &= \int_0^{\epsilon_f} \epsilon^{3/2} \cdot \frac{3}{2} \frac{N}{\epsilon_{f_0}^{3/2}} \, d\epsilon \\ &= \frac{3N}{2 \epsilon_{f_0}^{3/2}} \cdot \frac{\epsilon_f^{5/2}}{5/2} \end{aligned}$$

$$E = \frac{3}{5} N \epsilon_{f_0}$$

Hence total energy at $T=0$, is not zero, but proportional to fermi energy and no. of particles.

Q4)

Electrons are fermions i.e. $\frac{1}{2}$ spin particles. They follow Pauli's exclusion principle.

In order to find an expression for distribution of electrons, we need to find out thermodynamic probability Ω i.e. number of microstates in a configuration.

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Let $(N_1, N_2, N_3, \dots, N_N)$ particles be
there in $(E_1, E_2, E_3, \dots)$ energy
states.

Since the particles are indistinguishable,
there is only 1 way to choose $(N_1, N_2, \dots)$
particle in $(E_1, E_2, E_3, \dots)$ states

Now we need to distribute N_i particles
in g_i cells such that :

- ① Particles are indistinguishable
- ② Pauli Exclusion Principle is followed i.e.
1 cell of ~~any~~ g_i cells can take at
maximum 1 fermion.
- ③ A priori Probability of all microstates
is same.

No. of such combinations is given as

$$\Omega_i = \prod_i C_{g_i}^{N_i} = \frac{g_i^{N_i}}{N_i!}$$

Use min one $\Rightarrow$

$$\Rightarrow \ln \Omega = \sum [\ln(N_i) - \ln g_i - \ln(N_i - g_i)] \quad \text{--- (1)}$$

~~From~~

Also we know,

$$\sum n_i = N$$

$$\Rightarrow \sum dn_i = 0 \quad \text{--- (2)}$$

$$\sum E_i n_i = E$$

$$\Rightarrow \sum E_i dn_i = 0 \quad \text{--- (3)}$$

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At equilibrium, Entropy is maximum

$$\Rightarrow dS = 0$$

$$\Rightarrow k d(\ln \Omega) = 0 \quad (S = k \ln \Omega)$$

from ①, we have

$$d \ln \Omega = \sum d [\ln N_i - \ln g_i - \ln (N_i - g_i)] = 0 \quad \text{--- ④}$$

From sterling approximation,

$$\ln n \approx n \ln n - n$$

$$\Rightarrow d \ln n = \left[n \cdot \frac{1}{n} + \ln n - 1 \right] dn$$

$$\boxed{d(\ln n) = \ln n \, dn} \quad \text{--- ⑤}$$

Using ⑤ in ④, we get

$$\begin{aligned} d(\ln \Omega) &= \sum [\ln n_i \, dn_i - \ln (n_i - g_i) \, dn_i] \\ &= \sum \ln \left(\frac{n_i}{n_i - g_i} \right) dn_i \quad \text{--- ⑥} \end{aligned}$$

Multiply ② by α , ③ by β and

performing ② + ⑥ - ③, we get

$$\sum \left(\alpha + \ln \left(\frac{n_i}{n_i - g_i} \right) - \beta \epsilon_i \right) dn_i = 0$$

$$\Rightarrow \alpha + \ln \left(\frac{n_i}{n_i - g_i} \right) - \beta \epsilon_i = 0$$

$$\Rightarrow \left(\frac{n_i}{n_i - g_i} \right) = e^{-(\alpha - \beta \epsilon_i)}$$

$$\Rightarrow 1 - \frac{g_i}{n_i} = e^{-(\alpha - \beta \epsilon_i)} \Rightarrow \frac{n_i}{g_i} = \frac{1}{1 - e^{-(\alpha - \beta \epsilon_i)}}$$

Corrected $\Rightarrow$

$$\frac{n_i}{g_i} = \frac{1}{1 + e^{-(\alpha - \beta \epsilon_i)}}$$

$$= \frac{1}{1 + e^{\left(\frac{\epsilon_i - \epsilon_f}{kT} \right)}}$$

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$$N_i = f(\epsilon_i) g_i$$

$$= \frac{g_i}{e^{\frac{\epsilon_i - \epsilon_f}{kT}} + 1} = \frac{g_i}{A e^{\frac{\epsilon_i}{kT}} + 1}$$

$$\text{where } A = e^{-\frac{\epsilon_f}{kT}}$$

$$\frac{g_i}{N_i} = A e^{\left(\frac{\epsilon_i}{kT}\right) + 1}$$

For Maxwell distribution

$$\frac{g_i}{N_i} = A e^{+\left(\frac{\epsilon_i}{kT}\right)}$$

$\Rightarrow$ they both converge, when

$$\boxed{\frac{g_i}{N_i} \gg 1}$$

We can also find a quantum mechanical condition. For particles to be distinguishable (Maxwell), we should have

$$\lambda_{\text{de Broglie}} < \left(\frac{V}{N}\right)^{1/3}$$

$$\frac{h}{\sqrt{2mE}} \ll \left(\frac{V}{N}\right)^{1/3}$$

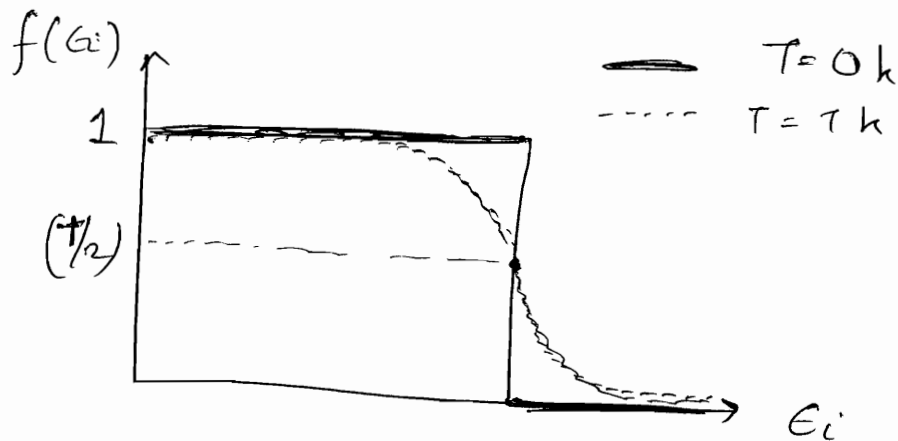
Approximating $E \approx kT$, we have

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$$\left(\frac{N}{V}\right)^{\frac{1}{3}} \frac{h}{\sqrt{2mkT}} < 1$$

$$\Rightarrow T > \left(\frac{N}{V}\right)^{\frac{2}{3}} \frac{h^2}{2mk}$$



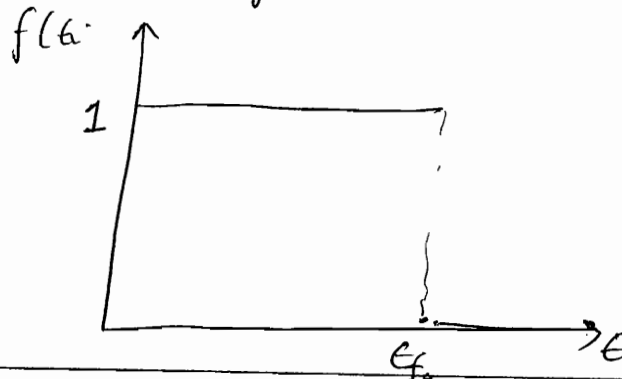
Q5(6) , Q1

Q7) Q4

Q8) Let the fermi momentum be given as p_{f_0} . We have fermi energy ϵ_{f_0} , given as

$$\epsilon_{f_0} = \left(\frac{p_{f_0}^2}{2m} \right)$$

At absolute zero, occupation index f_i for fermion is given as shown below



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Also density of state

$$g(E) = \frac{2\pi V}{h^3} \cdot (2m)^{3/2} E^{1/2}$$

Also spin degeneracy for fermions = 2

⇒ No. of fermions between energy E
and $E+dE$ are given by

$$dn = 2f(E) g(E) dE$$

$$\Rightarrow N = \int dn = \frac{4\pi V}{h^3} \cdot (2m)^{3/2} \int_0^{E_f} E^{1/2} dE$$

$$\Rightarrow N = \frac{4\pi V (2m)^{3/2}}{h^3} \cdot \frac{(E_f)^{3/2}}{(3/2)}$$

$$\text{Now } E_f = \frac{p_f^2}{2m}$$

$$\Rightarrow N = \frac{8\pi V}{3h^3} p_f^3 \quad \text{--- (1)}$$

Now energy of electrons, E is given as

$$E = \int E dn$$

$$= \int \frac{4\pi V (2m)^{3/2}}{h^3} E^{3/2} dE$$

$$= \frac{3}{2} N \cdot (E_f)^{-3/2} \int_0^{E_f} E^{3/2} dE$$

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$$= \frac{3}{2} \cdot N \cdot \frac{E_f^{-3/2}}{\frac{5}{2}} E_f^{5/2} = \frac{3}{5} N E_f$$

$$= \frac{3}{5} \cdot N \cdot \frac{p_{f0}^2}{2m}$$

$$\Rightarrow \boxed{E = \frac{3}{10} \frac{N p_{f0}^2}{m}}$$

$$\text{Also } PV = \frac{2}{5} E$$

$$\Rightarrow P = \frac{2}{5} \cdot \frac{3}{10} \cdot \frac{N}{V} \frac{p_{f0}^2}{m}$$

$$= \frac{2}{5} \cdot \frac{3}{10} \cdot m \frac{p_{f0}^2}{m}$$

From (1)

$$P = \frac{3}{25} \frac{n}{m} \cdot \left(\frac{3 h^3 N}{8 \pi V} \right)^{2/3}$$

$$\Rightarrow A \cdot n^{5/3}$$

Hence

$$\boxed{P \propto n^{5/3}}$$

It's good
habit to
define const.

$$A = \frac{3}{25m} \left(\frac{3 h^3}{8 \pi} \right)^{2/3}$$

(Q9) Q3

(Q10) Given $n_i = g_i e^{-\alpha - \beta \epsilon_i}$

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From equipartition theorem, the average energy per degree of freedom per molecule is $(1/2)kT$. For 3 degrees of freedom of a free molecule, we have

$$\langle E \rangle = \frac{3}{2} kT \quad \text{--- (1)}$$

Also, from classical statistical mechanics,

$$g_i = \frac{2\pi V}{h^3} \cdot (2m)^{3/2} \epsilon_i^{1/2}$$

$$= A \epsilon_i^{1/2}$$

where $A = \frac{2\pi V}{h^3} (2m)^{3/2} = \text{const.}$ for the given system.

$$\Rightarrow \langle E \rangle = \frac{E}{N} = \frac{\int \epsilon_i n_i}{\int n_i} = \frac{\int \epsilon_i A \epsilon_i^{1/2} e^{-\beta \epsilon_i} d\epsilon_i}{\int A \epsilon_i^{1/2} d\epsilon_i}$$

$$= \frac{\int e^{-\beta \epsilon_i} \epsilon_i^{3/2} d\epsilon_i}{\int e^{-\beta \epsilon_i} \epsilon_i^{1/2} d\epsilon_i}$$

$$= \frac{\cancel{\left(\frac{1}{\beta}\right)^{3/2}} \int \cancel{e^{-\beta \epsilon_i}} d\epsilon_i}{\cancel{\left(\frac{1}{\beta}\right)^{1/2}} \int \cancel{e^{-\beta \epsilon_i}} d\epsilon_i}$$

$$= \frac{3 \sqrt{\pi}}{4 \beta^{5/2}} = \frac{3}{4} \cdot \frac{1}{\beta} \cdot 2 = \frac{3}{2} \frac{1}{\beta} \quad \text{--- (2)}$$

$$\left(\frac{1}{\beta}\right)^{3/2} \cdot \frac{\sqrt{\pi}}{2}$$

Comparing (1) and (2), we have

$$\boxed{\beta = \frac{1}{kT}}$$

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Q11)

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This is similar to distribution of
fermions. Its given by

$${}^{15}C_{10} = \underline{\underline{3003}}$$

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Q12)

For fermions,

$$f(E) = \frac{1}{e^{\left(\frac{E-E_f}{kT}\right)} + 1}$$

Now for $E = E_f$, we have

$$f(E) = \frac{1}{e^0 + 1} = \left(\frac{1}{2}\right)$$

$$\Rightarrow \frac{n_i}{g_i} = \frac{1}{2}$$

i.e. half of the energy states are
filled for ~~every~~ energy level equal
to fermi energy.

$\therefore$ Probability of ~~energy~~ occupation of
quantum state is $(1/2)$

Q13)

According to F.D,

$$f(E) = \frac{1}{e^{\left(\frac{E-E_f}{kT}\right)} + 1} < 1$$

$\Rightarrow$ for every cell, occupancy is less than
1 $\Rightarrow$ cell can't have 2 fermions $\Rightarrow$
every fermion in unique energy state

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Q1)

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For a γ ray, Energy = 1 MeV

$$\text{Momentum } p = \frac{E}{c} = \frac{1 \times 10^6}{3 \times 10^8} \text{ eV}$$

$$= \underline{\underline{3.3 \times 10^{-3} \text{ eV/ms}^{-1}}}$$

In order to conserve momentum,

$$\begin{aligned} \text{Momentum of Particle} &= 3.3 \times 10^{-3} \text{ eV/s/m} \\ &= 5.23 \times 10^{-22} \text{ kg m/s} \end{aligned}$$

$$\Rightarrow \text{Energy of the Particle} = \sqrt{p^2 c^2 + m_0^2 c^4}$$

$$= \sqrt{2.5 \times 10^{-26} + 8.1 \times 10^{-19}} \text{ J}$$

$$= \underline{\underline{9 \times 10^{-10} \text{ J}}}$$

$$= \underline{\underline{5.625 \text{ MeV}}}$$

~~$$\begin{aligned} &= \sqrt{(3.3 \times 10^{-3})^2 + (9.1 \times 10^{-31})^2} \times 3 \times 10^8 \\ &= \sqrt{1.089 \times 10^{-5} + 8.281 \times 10^{-62}} \times 3 \times 10^8 \\ &= \sqrt{1.089 \times 10^{-5}} \times 3 \times 10^8 \\ &= 3.3 \times 10^{-3} \times 3 \times 10^8 \\ &= 9.9 \times 10^5 \text{ eV} \end{aligned}$$~~

$$m_0 c^2 \gg p^2 c^2$$

~~$$\Rightarrow T = E - m_0 c^2$$~~

$$\Rightarrow T = \left(\frac{p^2}{2m} \right)$$

~~$$= \sqrt{(m_0 c^2)^2 + (pc)^2} - m_0 c^2$$~~

~~$$= m_0 c^2 \sqrt{1 + \left(\frac{pc}{m_0 c^2} \right)^2} - m_0 c^2$$~~

~~$$= m_0 c^2 + \frac{1}{2} \left(\frac{pc}{m_0 c^2} \right)^2 m_0 c^2$$~~

~~$$= \frac{1}{2} \frac{p^2 c^2}{m_0 c^2} = \frac{1}{2} \frac{p^2}{m_0}$$~~

$$= \underline{\underline{88.88 \text{ eV}}}$$

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(Q2)

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Differential Cross section refers to the number of particles scattered per unit time per solid angle as a fraction of no. of particles incident per unit time per unit area.

$$\sigma(\theta) = \frac{dN}{I d\Omega}$$

where dN : No. of particles scattered per unit time

$d\Omega$: solid angle

I : No. of particles incident per unit area per unit time

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mention
the
symbols
in
the
definition

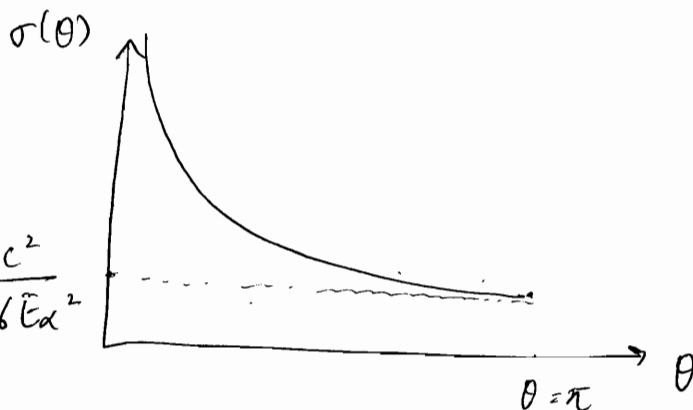
For Rutherford scattering,

$$\sigma(\theta) = \frac{C^2}{16E_\alpha^2} \csc^4\left(\frac{\theta}{2}\right)$$

$$\text{where } C = \frac{Z_1 Z_2 e^2}{4\pi\epsilon_0}$$

θ : scattering angle

E_α : energy of α particle



$$\begin{aligned} \sigma &= \frac{C^2}{16E_\alpha^2} \int_0^\pi \csc^4\left(\frac{\theta}{2}\right) \cdot 2\pi \sin\theta d\theta \\ &= \frac{4\pi C^2}{16E_\alpha^2} \int_0^\pi \frac{1}{\sin^4\left(\frac{\theta}{2}\right)} \sin\left(\frac{\theta}{2}\right) \cos\left(\frac{\theta}{2}\right) d\theta \end{aligned}$$

Also note that
this curve signifies
that scattering
around $\theta=0$
is maximum
 $\Rightarrow$ nucleus is
extremely small.
As most of the
particles
remain
unscattered!!

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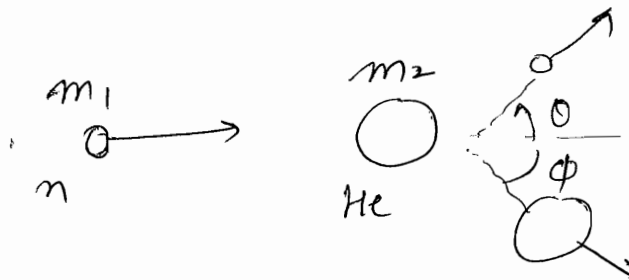
Put $\sin \frac{\theta}{2} = x$, we have

$$\sigma = \frac{4\pi C^2}{16E\alpha^2} \int_0^{\alpha} \frac{2dx}{x^3} \rightarrow \infty$$

The range of the Coulomb force is infinite.
Hence infinite area of cross section where
force is felt.

However in reality, electrons screen the
nucleus and effectively force is zero
beyond a certain distance, and therefore a
finite scattering cross section is there.

(Q3)



velocity of neutron $u_1 = \sqrt{\frac{2E}{m}}$

In the centre of mass system

$$u_1' = u_{1,cm} = u_1 - u_{cm} = u_1 - \frac{m_1 u_1}{m_1 + m_2} = \left(\frac{m_2 u_1}{m_1 + m_2} \right)$$

$$u_2' = u_{2,cm} = u_2 - u_{cm} = 0 - \frac{m_1 u_1}{m_1 + m_2} = \left(-\frac{m_1 u_1}{m_1 + m_2} \right)$$

$$u_1' = -\frac{m_2}{m_1} u_2'$$

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Now from momentum conservation,

$$m_1 v_1' \cos \theta' + m_2 v_2' \cos \phi' = m_1 u_1' + m_2 u_2' = 0$$

$$\Rightarrow m_1 v_1' \cos \theta' = -m_2 v_2' \cos \phi' \quad \text{--- (1)}$$

Also

$$m_1 v_1' \sin \theta' = m_2 v_2' \sin \phi' = 0$$

$$\Rightarrow m_1 v_1' \sin \theta' = m_2 v_2' \sin \phi' \quad \text{--- (2)}$$

From (1) and (2)

$$\boxed{\theta' = \pi - \phi'}$$

Also from energy conservation,

$$m_1 v_1'^2 + m_2 v_2'^2 = m_1 u_1'^2 + m_2 u_2'^2 \quad \text{--- (3)}$$

Replacing v_1' by $\frac{m_2}{m_1} v_2'$ and u_2' by $-\frac{m_2}{m_1} u_1'$, we have

$$\boxed{\begin{aligned} |v_1'| &= |u_1'| \\ |v_2'| &= |u_2'| \end{aligned}}$$

Hence the magnitude of velocities in COM frame before & after collision remain same.

Now in COM frame

$$u_1' = \frac{m_2 u_1}{m_1 + m_2} = \frac{4}{1+4} u_1 = \left(\frac{4}{5}\right) u_1$$

$$u_2' = -\frac{m_1 u_1}{m_1 + m_2} = -\left(\frac{1}{5}\right) u_1$$

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$$\Rightarrow |\text{Momentum Neutron}| = \frac{4}{5} m_N u_1 = \frac{4}{5} \sqrt{2mE}$$

$$= \underline{\underline{1.84 \times 10^{-20} \text{ kg m/s}}}$$

$$|\text{Momentum Helium}| = \frac{1}{5} m_{He} u_1 = \frac{4}{5} m u_1$$

$$= \underline{\underline{1.84 \times 10^{-20} \text{ kg m/s}}}$$

Also the momentum of individual particles before & after the collision remain same.

Q4)

<Q27

$$\sigma(\theta) = A e^{-B(\vec{k}_i - \vec{k}_f)^2}$$

$$|\vec{k}_i - \vec{k}_f|^2 = k_i^2 + k_f^2 - 2k_i k_f \cos \theta$$

$$= 2k^2(1 - \cos \theta)$$

$$= 4k^2 \sin^2\left(\frac{\theta}{2}\right)$$

$$\Rightarrow \sigma(\theta) = \cancel{A} k^2 A e^{-4Bk^2 \sin^2\left(\frac{\theta}{2}\right)}$$

$$\sigma = \int_0^\pi \sigma(\theta) 2\pi \sin \theta d\theta$$

$$= 2\pi A \int_0^\pi e^{-4Bk^2 \sin^2\left(\frac{\theta}{2}\right)} \sin \theta d\theta$$

Put $\sin^2 \frac{\theta}{2} = x \Rightarrow 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2} \cdot \frac{1}{2} d\theta = dx$

$$\sigma = 2\pi A \int_0^1 e^{-4Bk^2 x} 2dx$$

$$\sigma = 4\pi A \int_0^1 e^{-4Bk^2 x} dx = 4\pi A \frac{(1 - e^{-4Bk^2})}{4Bk^2}$$

$$= \underline{\underline{\frac{\pi A}{Bk^2} (1 - e^{-4Bk^2})}}$$

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Q5)

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From Q-3)

$$u_1' = -\frac{m_2}{m_1} u_2'$$

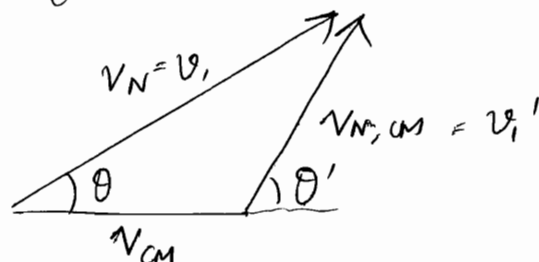
$$v_1' = \frac{m_2}{m_1} v_2'$$

$$|u_1'| = |u_2'|$$

$$|u_2'| = |v_2'|$$

$$\theta' = \pi - \theta$$

Now for neutron, we can see



$$v_1' \cos \theta' + v_{cm} = v_1 \cos \theta$$

$$v_1' \sin \theta' = v_1 \sin \theta$$

$$\text{Also, } \frac{E'}{E} = \frac{v_1'^2}{v_1^2} = \frac{(v_1' \sin \theta')^2 + (v_1' \cos \theta' + v_{cm})^2}{v_1^2}$$

$$= \frac{v_1'^2 + v_{cm}^2 + 2 v_{cm} v_1' \cos \theta'}{v_1^2}$$

$$= \frac{\frac{m_2^2}{m_1^2} u_1^2 + \left(\frac{m_1 u_1}{m_1 + m_2}\right)^2 + 2 \frac{m_1}{m_1 + m_2} u_1^2 \cos \theta'}{u_1^2}$$

$$u_1' = \left(\frac{m_2 u_1}{m_1 + m_2}\right)$$

$$= \left(\frac{m_2}{m_1 + m_2}\right)^2 + \left(\frac{m_1}{m_1 + m_2}\right)^2 + \frac{2 m_1 m_2 \cos \theta'}{(m_1 + m_2)^2}$$

$$\text{Now } m_2 = A m_1$$

$$\therefore \frac{E'}{E} = \left(\frac{A+1}{1+A}\right)^2 + \left(\frac{1}{1+A}\right)^2 + \left(\frac{2A}{(1+A)^2}\right) \cos \theta'$$

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Simply apply
Cosine theorem

$$\cos \theta' = \frac{v_{cm}^2 + v_1'^2 - v_1^2}{2 v_{cm} v_1'}$$

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$$\Rightarrow \frac{E'}{E} = \frac{1 + A^2 + 2A \cos \theta}{(1 + A)^2}$$

$$\Rightarrow \boxed{\frac{E'}{E} = \frac{1 + A^2 + 2AMc}{(1 + A)^2}}$$

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(Q6)

Centre of the mass of a system of particles refers to the point where

- (i) whole mass of the system is considered to be concentrated
- (ii) net force acting on the system is supposed to act.

Let a Consider a system of n particles such that each experience a force $\vec{F}_i$ and has position vector $\vec{r}_i$

$$\begin{aligned} \text{i.e. } \vec{F}_{\text{total}} &= \vec{F}_1 + \vec{F}_2 + \dots + \vec{F}_n \\ &= m_1 \frac{d^2 \vec{r}_1}{dt^2} + m_2 \frac{d^2 \vec{r}_2}{dt^2} + \dots \end{aligned}$$

Now at a point called C.O.M., where Mass is concentrated, and net force is supposed to act.

$$\Rightarrow \vec{F}_{\text{total}} = (M_1 + M_2 + \dots + M_n) \frac{d^2 \vec{R}_{\text{cm}}}{dt^2}$$

$$\Rightarrow \frac{d^2 \vec{R}_{\text{cm}}}{dt^2} = \frac{d^2}{dt^2} \frac{[m_1 \vec{r}_1 + m_2 \vec{r}_2 + \dots + m_n \vec{r}_n]}{[M_1 + M_2 + M_3 + \dots]}$$

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$$\Rightarrow \vec{R}_{CM} = \frac{\sum m_i \vec{r}_i}{\sum m_i}$$

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$$\begin{aligned} \text{velocity of C.O.M} &= \frac{d\vec{R}_{CM}}{dt} \\ &= \frac{d}{dt} \vec{R}_{CM} = \frac{d}{dt} \frac{\sum m_i \vec{r}_i}{\sum m_i} = \sum \frac{m_i \frac{d\vec{r}_i}{dt}}{\sum m_i} \end{aligned}$$

$$\vec{v}_{CM} = \frac{\sum m_i \vec{v}_i}{\sum m_i}$$

$$Q7) \quad \sigma(\theta) = \alpha \operatorname{cosec}^4\left(\frac{\phi}{2}\right) = \frac{d\sigma}{d\Omega}$$

$$\begin{aligned} \therefore N_{\text{new}} = d\sigma &= \sigma(\phi) \cdot 2\pi \sin\phi \, d\phi \\ &= \alpha \cdot \operatorname{cosec}^4\left(\frac{\phi}{2}\right) \cdot 2\pi \sin\phi \, d\phi \end{aligned}$$

$$d\sigma = 4\pi\alpha \frac{\cos(\phi/2)}{\sin^3(\phi/2)} d\phi$$

$$\left[\frac{d\sigma}{d\phi} = \frac{4\pi\alpha \cos(\phi/2)}{\sin^3(\phi/2)} \right]$$

Q8)

$$\text{Thrust} = v_{\text{ex}} \frac{dm}{dt} = mg$$

$$\Rightarrow 10000 = 4.5 \times 10^3 \left(\frac{dm}{dt} \right)$$

$$\Rightarrow \left(\frac{dm}{dt} \right) = 2.22 \text{ kg/s}$$

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$$\text{Thrust} = v \left(\frac{dm}{dt} \right) = 4.5 \times 10^3 \times 2.50$$

$$= 11250 \text{ N}$$

$$\vec{a} = \frac{\text{Thrust} - mg}{m} = 1.25 \text{ m/s}^2$$

$$\Rightarrow \vec{v}_f = \vec{v}_i + \vec{a}t = 0 + 1.25 \times 8$$

$$\Rightarrow \boxed{v_f = 10 \text{ m/s}}$$

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Q9)

Ignoring relativistic effects,

$$E \approx \frac{p^2}{2m} \Rightarrow p = \sqrt{2m \cdot E}$$

$$p_{\text{Nucleus}} = \sqrt{2m_{\alpha} E}$$

$$\Rightarrow E_{\text{Nucleus}} = \frac{2m_{\alpha} E}{2m_{\text{Nucleus}}} = \frac{2 \cdot 4 \cdot E}{2(m_0 - 4)}$$

$$= \frac{4E}{(m_0 - 4)}$$

Energy evolved during disintegration = Q

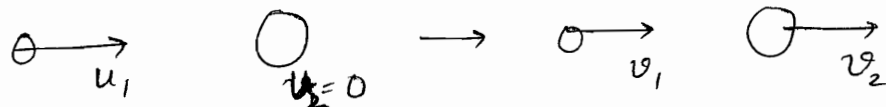
$$= (K.E.)_{\text{product}} - (K.E.)_{\text{reactant}}$$

$$= (K.E.)_{\text{product}}$$

$$= E + \frac{4E}{m_0 - 4} = \frac{(m_0 + 4)E}{(m_0 - 4)}$$

⊗ Note that
there is
nothing wrong
to ignore
relativistic
effects in
nuclear
reactions.

Q10)



From conservation of momentum,

$$m_1 u_1 + m_2 u_2 = m_1 v_1 + m_2 v_2$$

$$m_1 u_1 = m_1 v_1 + m_2 v_2 \quad \text{--- (1)}$$

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Also for an elastic collision Coefficient
of restitution, $e = 1$

$\Rightarrow$ Velocity of Approach of particles =
velocity of their separation

$$\Rightarrow u_1 = v_2 - v_1 \quad \text{--- (2)}$$

From (1) and (2),

$$2m_1 u_1 = m_2 v_2 + m_1 v_2$$

$$\Rightarrow v_2 = \frac{2m_1 u_1}{m_1 + m_2}$$

$$\text{Also } v_1 = v_2 - u_1 = \frac{2m_1 u_1}{m_1 + m_2} - u_1$$

$$\Rightarrow v_1 = \left(\frac{m_1 - m_2}{m_1 + m_2} \right) u_1$$

$$\frac{\text{decrease in k.E}}{\text{k.E initial}} = - \frac{\left[\left(\frac{m_1 - m_2}{m_1 + m_2} \right) u_1^2 - u_1^2 \right]}{u_1^2}$$

$$\Rightarrow \text{fractional decrease} = \frac{4m_1 m_2}{(m_1 + m_2)^2}$$

Now neutron mass = H_{mass}

$$\Rightarrow v_1 = 0$$

if Hydrogen nucleus is used to slow down
neutrons in a nuclear reactor.

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But D_2O is used as :

- ① Deuterium has better cross section of absorption.
- ② Unlike H, Deuterium does not absorb neutrons.

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Q 11)

<Q6>

Proof of Unique Centre of Mass

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Usefulness of centre of mass lies in the fact that when no external force acts, the Centre of Mass remains in const. motion. When we shift the frame of reference to Centre of Mass itself, the net momentum becomes zero, and no. of unknown variables reduce and analysis is easy in case of collisions or scattering or annihilation when net force external = 0

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Also, in order to study motion of a system, all we need to know is motion of Centre of Mass, for however complicated the system is internal forces cancel each other, Centre of mass behaves as if the whole mass is concentrated on it, and the external forces directly apply to it.

Q12)

Take Carbon Atom as Origin,

$$\vec{R}_{cm} = \frac{m_c \vec{r}_c + m_o \vec{r}_o}{m_c + m_o}$$

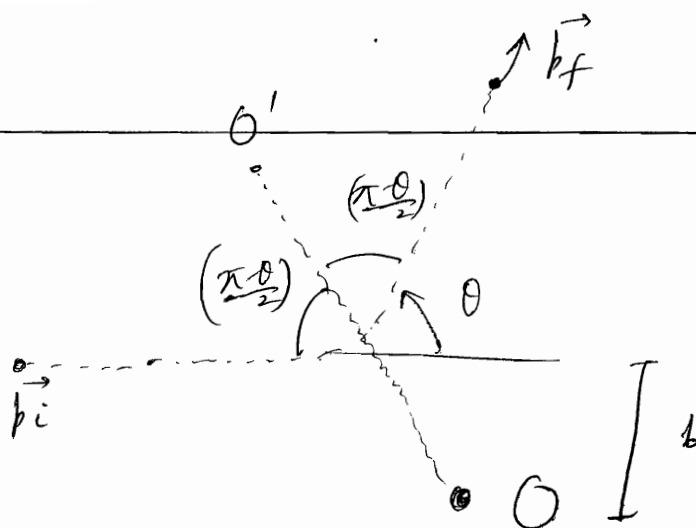
$$= \frac{12 \times 0 + 16 \times 1.2 \text{ \AA}}{28}$$

$$= \underline{0.685 \text{ \AA}} \quad \text{away from Carbon atom}$$

Q13)

Let the Origin of Central force be O. Since the force is central, torque = 0 hence angular momentum remains conserved

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The net force acts along OO' .

Also impulse = $\Delta \vec{p}$

$$\text{i.e. } \int F_{OO'} dt = \vec{p}_f - \vec{p}_i$$

since the collision is elastic, $|\vec{p}_f| = |\vec{p}_i|$

we have

$$\begin{aligned} |\vec{p}_f - \vec{p}_i| &= \sqrt{p_f^2 + p_i^2 - 2p_f p_i \cos \theta} \\ &= 2p \sin\left(\frac{\theta}{2}\right) \end{aligned}$$

$$\Rightarrow \boxed{\int F_{OO'} dt = 2p \sin\left(\frac{\theta}{2}\right)} \quad \text{--- (1)}$$

Also the angular momentum J
remains conserved

$$J_{\text{initial}} = p_i b$$

$$J(r) = mr^2 \left(\frac{d\phi}{dt} \right)$$

$$\Rightarrow p_i b = mr^2 \left(\frac{d\phi}{dt} \right) \quad \text{--- (2)}$$



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Using in (2), we have

$$\int F_{00'} \cdot \frac{m r^2}{b^2} d\phi = 2b \sin\left(\frac{\theta}{2}\right)$$

$$\text{New } F_{00'} = F \cos\theta = \frac{k}{r^2} \cos\theta$$

$$\Rightarrow \int \frac{k}{r^2} \cos\theta \cdot \frac{m r^2}{b^2} \cdot d\phi = 2b \sin\left(\frac{\theta}{2}\right)$$

Also ϕ varies from $-\left(\frac{\pi-\theta}{2}\right)$ to $\left(\frac{\pi-\theta}{2}\right)$

$$\Rightarrow \int_{-\left(\frac{\pi-\theta}{2}\right)}^{\left(\frac{\pi-\theta}{2}\right)} \frac{k m}{b^2} \cdot \cos\phi \cdot d\phi = 2b \sin\left(\frac{\theta}{2}\right)$$

$$\Rightarrow \frac{k m}{b^2} \cdot 2 \cos\left(\frac{\theta}{2}\right) = 2b \sin\left(\frac{\theta}{2}\right)$$

$$\Rightarrow \cot\left(\frac{\theta}{2}\right) = \frac{b^2}{k m} = \frac{m v_0^2 b}{k}$$

Hence Proved!!

Q14) Q57

$$v_1 \cos\theta = v_1' \cos\theta' + v_{cm}$$

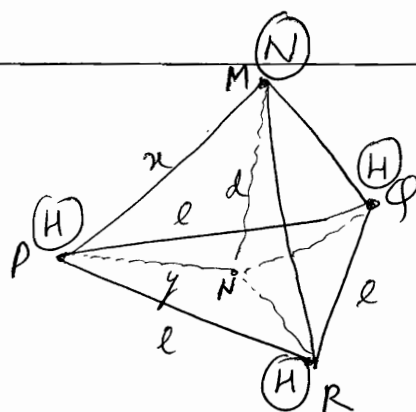
$$v_1 \sin\theta = v_1' \sin\theta'$$

$$\Rightarrow \tan\theta = \frac{\sin\theta'}{\cos\theta' + \frac{v_{cm}}{v_1'}} = \left[\frac{\sin\theta'}{\cos\theta' + \frac{m_1}{m_2}} \right]$$

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Q15)

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$$y = 0.939 \text{ \AA}$$

$$x = 1.014 \text{ \AA}$$

$$d = \sqrt{x^2 - y^2}$$

$$d = 0.38 \text{ \AA}$$

wrt. N atom, we can see from symmetry
that COM will lie on MN

$$\begin{aligned} \Rightarrow y_{\text{CM}} &= \frac{y_N + y_{H_1} + y_{H_2} + y_{H_3}}{m_N + 3m_H} \\ &= \frac{0 + 0.38 \times 3}{14 + 3} = \frac{0.38 \times 3}{17} \\ &= \underline{\underline{0.067 \text{ \AA}}} \text{ from N atom} \end{aligned}$$

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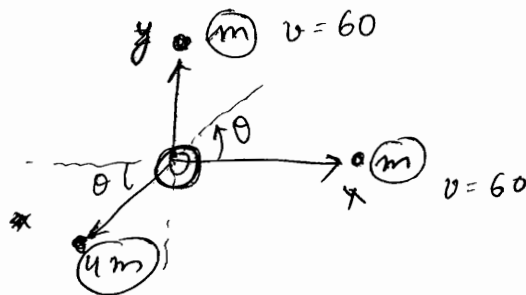
Q16)

From Newton's 2nd law,

$$\vec{F}_{\text{ext}} = \frac{d\vec{p}}{dt}$$

$$\text{if } \vec{F}_{\text{ext}} = 0 \Rightarrow \vec{p} = \text{const.}$$

Hence if net external force = 0, the
total momentum of a system of particles
remains conserved.



Since initially bomb is stationary, the

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momentum $\neq 0$

$\Rightarrow$ Final momentum = 0 (no external force)

let the direction of velocity be $-\theta$ s.t.

$$v_x = -v \cos \theta \hat{i} - v \sin \theta \hat{j}$$

(from figure 1)

To conserve momentum

$$p_x = 0$$

$$\Rightarrow m \cdot 60 - 4m \cdot v \cos \theta = 0$$

$$\Rightarrow v \cos \theta = \frac{60}{4} = 15 \text{ m/s}$$

$$p_y = 0$$

$$-4m v \sin \theta + m \cdot 60 = 0$$

$$\Rightarrow v \sin \theta = 15 \text{ m/s}$$

$$\Rightarrow \boxed{\theta = 45^\circ}$$

$$\boxed{v = 15\sqrt{2} \text{ m/s}}$$

Q 17)

For the rocket motion

$$M = M_0 - m'$$

$$\Rightarrow \boxed{dM = -dm'} \quad \text{--- (1)}$$



$M_0 - m'$: mass after
ejection

m' : mass of gases

Now, we know

$$\vec{F}_{\text{ext}} dt = \frac{d\vec{p}}{dt}$$

$$\Rightarrow d\vec{F}_{\text{ext}} dt = \vec{p}_f - \vec{p}_i$$

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Let dm mass be removed and the
velocity increases by dv . Let u_{exhaust} be the
velocity of removed mass w.r.t. rocket.

$$\Rightarrow d\vec{F}dt = (M - dm)(\vec{v} + d\vec{v}) + dm(\vec{u}_{\text{ex}} + \vec{v} + d\vec{v}) - M\vec{v}$$

$$\Rightarrow d\vec{F}dt = M d\vec{v} - dm \vec{v} + dm \vec{u}_{\text{ex}} + dm \vec{v}$$

$$d\vec{F}dt = M d\vec{v} + dm \vec{u}_{\text{ex}}$$

Taking components along motion of
rocket,

$$-Mg dt = M dv - u_{\text{ex}} dm$$

Also $dm = -dM$, we have,

$$-Mg = M dv + u_{\text{ex}} dM$$

$$\Rightarrow dv = -g dt - u_{\text{ex}} \frac{dM}{M}$$

$$\Rightarrow \int dv = -g t - u_{\text{ex}} \int_{M_0}^M \frac{dM}{M}$$

$$\Rightarrow v_2 = v_1 - gt + u_{\text{ex}} \ln\left(\frac{M_0}{M}\right)$$

Given $-\frac{dM}{dt} = k$

$$\Rightarrow kt = (M_f + M_0) - (M_0)$$

$$\Rightarrow t = \left(\frac{M_f}{k}\right)$$

$$\Rightarrow v_2 = v_1 - \frac{g M_f}{k} + u_{\text{ex}} \ln\left(\frac{M_0 + M_f}{M_0}\right)$$

$$\Rightarrow v_2 = v_1 - \frac{g M_f}{k} + u_{\text{ex}} \ln\left(1 + \frac{M_f}{M_0}\right)$$

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M_f : Mass of fuel
 M_0 : initial Mass of
rocket w/
fuel

Q18)

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Q19)

Q12)

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$$\vec{P}_i = \vec{P}_f$$

$$\Rightarrow 40 * 4 - 60 * 2 = 100 * v$$

$$\Rightarrow \boxed{v = 0.4 \text{ m/s}} \text{ in direction of Ram's travel.}$$

$$\begin{aligned} \text{loss in k.E.} &= \frac{1}{2} \cdot 40 \cdot 16 + \frac{1}{2} \cdot 60 \cdot 4 - \frac{1}{2} \cdot 100 \cdot \frac{16}{100} \\ &= 320 + 120 - 8 \\ &= \underline{\underline{432 \text{ J}}} \end{aligned}$$

✓ Q20)

Scattering Cross section is the number of particles scattered per unit time as a fraction of no. of particles per unit time per unit area. It is a measure of area of cross section responsible for scattering of particles.

< Q13)

$$b = \frac{c}{2E\alpha} \cot\left(\frac{\theta}{2}\right) \Rightarrow \sigma(\theta) = \frac{b}{\sin\theta} \left(\frac{db}{d\theta}\right)$$

~~$$\sigma = \int_0^\pi \sigma(\theta) \cdot 2\pi \sin\theta d\theta$$~~

$$\sigma(\theta) = \frac{c^2}{16E\alpha^2} \operatorname{cosec}^4\left(\frac{\theta}{2}\right)$$

$$\begin{aligned} \Rightarrow \sigma &= \int \sigma(\theta) \cdot 2\pi \sin\theta d\theta \\ &= \infty \end{aligned}$$

Q 21)

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Rocket equation,

$$v_f = v_i - gt + u_{ex} \ln\left(\frac{M_0}{M}\right)$$

$$\frac{dm}{dt} = -\frac{m}{200}$$

$$\Rightarrow \int dm = -\frac{m}{200} \int dt$$

$$\Rightarrow (1 - 0.27)m = \frac{m}{200} t$$

$$\Rightarrow t = 200 * 0.73 = \underline{146.8}$$

★ do not
neglect the
"-gt"

factor

$$\Rightarrow v_f = 0 - 146 * 9.8 + 2500 \ln\left(\frac{m}{0.27m}\right)$$

$$v_f = \underline{1842.5 \text{ m/s}}$$

Q 22)

< Q 13 >

$$b = \frac{C}{2E\alpha} \cdot \cot\left(\frac{\theta}{2}\right)$$

$$= \frac{1}{4\pi\epsilon_0} \cdot \frac{z_1 z_2 e^2}{2E\alpha} \cot\left(\frac{\theta}{2}\right)$$

$$= \frac{9 \times 10^9 \times 22 \times (1.6 \times 10^{-19})^2}{2 \times 6 \times 10^{-13}} \times \cot(15)$$

$$= \underline{1.43 \times 10^{-15}}$$

Q 23)

$$8 \times 10^3 = -9.8 * 30 + 3000 \ln\left(\frac{M_0}{M}\right)$$

$$\Rightarrow \left(\frac{M_0}{M}\right) = 15.87$$

Q 24)

< Q 13 >

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Q27)
Q28)
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Q17)
Let velocity of cm be v_{cm} .

$$\vec{v}_{i,cm} = \vec{v}_i - \vec{v}_{cm}$$

$$\vec{p}_{i,cm} = m_i [\vec{v}_i - \vec{v}_{cm}]$$

$$\begin{aligned} \vec{p}_{system,cm} &= \sum_i m_i [\vec{v}_i - \vec{v}_{cm}] \\ &= \sum_i m_i \vec{v}_i - \sum_i m_i \vec{v}_{cm} \\ &= \sum_i m_i \vec{v}_{cm} - \sum_i m_i \vec{v}_{cm} \\ &= 0 \end{aligned}$$

Q29)

For a conservative force field $\vec{F}$, $\vec{\nabla} \times \vec{F} = 0$

For the given Force

$$\vec{\nabla} \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 2xy+z^3 & x^2 & 3xz^2 \end{vmatrix}$$

$$\begin{aligned} &= \hat{i} \left[\frac{\partial}{\partial y} (3xz^2) - \frac{\partial}{\partial z} (x^2) \right] \\ &+ \hat{j} \left[\frac{\partial}{\partial z} (2xy+z^3) - \frac{\partial}{\partial x} (3xz^2) \right] \\ &+ \hat{k} \left[\frac{\partial}{\partial x} x^2 - \frac{\partial}{\partial y} (2xy+z^3) \right] \end{aligned}$$

$$\begin{aligned} &= \hat{i} [0] + \hat{j} [3z^2 - 3z^2] \\ &+ \hat{k} [2x - 2x] \end{aligned}$$

$$= 0$$

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Hence a conservative force.

For this force, let Potential be V

$$\Rightarrow \vec{F} = -\vec{\nabla} V$$

$$\Rightarrow V = -\int F_x dx + f_1(y, z) \quad \text{--- (1)}$$

$$V = -\int F_y dy + f_2(x, z) \quad \text{--- (2)}$$

$$V = -\int F_z dz + f_3(x, y) \quad \text{--- (3)}$$

From (1),

$$V = -x^2y + z^3x + C_1(y, z) \quad \text{--- (4)}$$

From (2)

$$V = -x^2y + C_2(x, z) \quad \text{--- (5)}$$

from (3)

$$V = -xz^3 + C_3(x, y) \quad \text{--- (6)}$$

From (4), (5), (6), we have

$$\boxed{V = -x^2y - z^3x + C}$$

where C is const.

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(Q1)

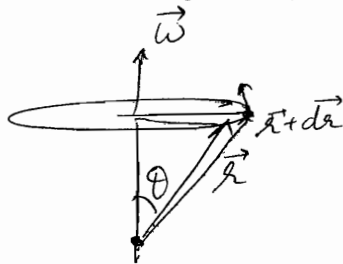
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A rotating frame of reference is a non inertial frame of reference. Therefore in order to validate the use of Newton's laws in such a frame, Pseudo Forces needed to be applied. Coriolis force is one such pseudo force, that is applied on particles that are moving on a rotating body.

$$C.F. = -2m(\vec{\omega} \times \vec{v})$$

where $\vec{\omega}$: angular velocity vector of rotating body

$\vec{v}$: velocity of particle as seen on rotating body.



Now Consider a ~~moving~~ rotating body in the above figure

$$\vec{v} = \frac{d\vec{r}}{dt} = \frac{r \sin \theta \frac{d\phi}{dt}}{dt} = r \sin \theta \left(\frac{d\phi}{dt} \right)$$

$$= \omega r \sin \theta$$

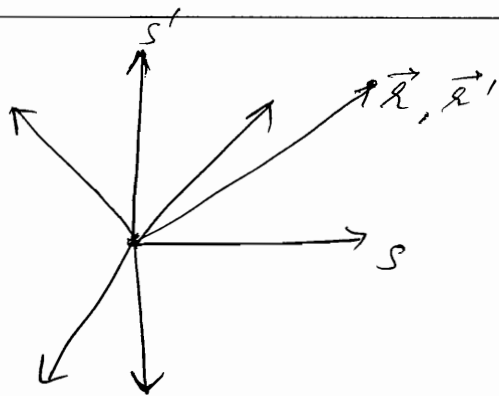
$$= \vec{\omega} \times \vec{r}$$

$\therefore$ for any vector $\vec{r}$ fixed in rotating body's frame of reference

$$\left(\frac{d\vec{r}}{dt} \right) = \vec{\omega} \times \vec{r}$$

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Consider two frames S and S' where
S' is rotating at angular velocity ω .
Now since origin is common

$$\vec{r} = \vec{r}'$$

i.e. $x\hat{i} + y\hat{j} + z\hat{k} = x'\hat{i}' + y'\hat{j}' + z'\hat{k}'$

The velocity vector $\frac{d\vec{r}}{dt}$ is given as

$$\begin{aligned} \frac{d\vec{r}}{dt} &= \left(\frac{dx'}{dt} \hat{i} + \frac{dy'}{dt} \hat{j} + \frac{dz'}{dt} \hat{k} \right) + x' \frac{d\hat{i}'}{dt} \\ &\quad + y' \frac{d\hat{j}'}{dt} + z' \frac{d\hat{k}'}{dt} \\ &= \left(\frac{d\vec{r}'}{dt} \right)_R + x' (\vec{\omega} \times \hat{i}') + y' (\vec{\omega} \times \hat{j}') \\ &\quad + z' (\vec{\omega} \times \hat{k}') \end{aligned}$$

$$\frac{d\vec{r}}{dt} = \left(\frac{d\vec{r}'}{dt} \right)_R + \vec{\omega} \times \vec{r}$$

↑
velocity
measured
from inertial
frame

↑
velocity
measured
from rotating
frame

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Similarly for velocity vector $\vec{v}$,

$$\frac{d\vec{v}}{dt} = \left(\frac{d\vec{v}}{dt} \right)_R + \vec{\omega} \times \vec{v}$$

Now $\vec{v} = \vec{v}' + \vec{\omega} \times \vec{r}$

$$\Rightarrow \frac{d\vec{v}}{dt} = \left(\frac{d(\vec{v}' + \vec{\omega} \times \vec{r})}{dt} \right)_R + \vec{\omega} \times (\vec{v}' + \vec{\omega} \times \vec{r})$$

$$= \vec{a} = \vec{a}' + \left(\frac{d\vec{\omega}}{dt} \times \vec{r} \right)_R + \left(\vec{\omega} \times \frac{d\vec{r}}{dt} \right)_R + \vec{\omega} \times \vec{v}' + \vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$\Rightarrow \vec{a} = \vec{a}' + 2(\vec{\omega} \times \vec{v}') + \vec{\omega} \times (\vec{\omega} \times \vec{r}) + \left(\frac{d\vec{\omega}}{dt} \right) \times \vec{r}$$

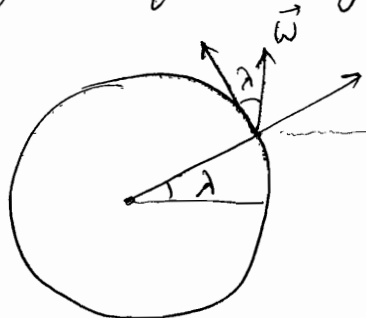
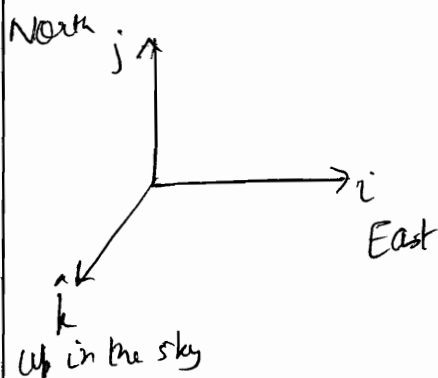
$$\Rightarrow \vec{a}' = \vec{a} - 2(\vec{\omega} \times \vec{v}) - \vec{\omega} \times (\vec{\omega} \times \vec{r}) - \left(\frac{d\vec{\omega}}{dt} \right) \times \vec{r}$$

Q2)

<Q1>

2 examples

① A body falling freely from height H.



$$\vec{\omega} = \omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}$$

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$$\text{Coriolis Force} = -2m (\vec{\omega} \times \vec{v})$$

Now

$$\vec{v} = -gt \hat{k}$$

$$\Rightarrow C.F. = -2m [\omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}] \times -gt \hat{k}$$

$$= 2mgt \omega \cos \lambda \hat{i}$$

$$\Rightarrow \vec{a}_{\text{East}} = 2gt \omega \cos \lambda$$

~~$$\vec{s} = \frac{1}{2} \vec{a} t^2 + \vec{u} t$$~~

$$\Rightarrow \vec{s}_{\text{East}} = \frac{2gt^3 \omega \cos \lambda}{6}$$
~~$$= \frac{1}{2} 2gt \omega \cos \lambda t^2$$~~

$$= \frac{gt^3 \omega \cos \lambda}{3}$$
~~$$\vec{s} = \frac{1}{2} \vec{a} t^2$$~~

$$\text{Also } H = \frac{1}{2} g t^2 \Rightarrow t = \left(\frac{2H}{g} \right)^{1/2}$$

$$\Rightarrow \boxed{s_{\text{East}} = \frac{g \omega \cos \lambda}{3} \cdot \left(\frac{2H}{g} \right)^{3/2}}$$

② Foucault's law

let $\vec{v}$ on the surface be $[v_E \hat{i} + v_N \hat{j}]$

$$\Rightarrow C.F. = -2m [\omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}] \times [v_E \hat{i} + v_N \hat{j}]$$

$$(C.F.)_{ij \text{ plane}} = -2m\omega [\sin \lambda v_E \hat{j} + \sin \lambda v_N \hat{i}]$$

$$\vec{F}_{\text{on earth surface}} = 2m\omega \sin \lambda [v_N \hat{i} - v_E \hat{j}]$$

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$\therefore$ we have 2 observations

① more λ , more Coriolis force.
 $\therefore$ no tropical cyclones on Equator

② in N-H. $\sin \lambda > 0 \Rightarrow$ force towards
right of motion.
in S-H $\sin \lambda < 0 \Rightarrow$ force towards
left of motion

This is called Ferrel's law and explains
direction of winds and Ocean
currents.

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Q3) Q2

Q4) For a rotating frame, we have
Centrifugal pseudo force acting on every
body.

$$\begin{aligned} \text{Cf. force} &= -m \vec{\omega} \times (\vec{\omega} \times \vec{r}) \\ &= -m [\omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}] \times [\\ &\quad [\omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}] \times [r \hat{k}]] \\ &= -m [\omega \cos \lambda \hat{j} + \omega \sin \lambda \hat{k}] \times \\ &\quad [\omega r \cos \lambda \hat{i}] \\ &= +m \omega^2 r \cos^2 \lambda \cdot [\hat{k}] \\ &\quad - m \omega^2 r \cos \lambda \sin \lambda \cdot \hat{j} \end{aligned}$$

Taking the component along earth's
Earth centre

$$\text{Cf. force} = m \omega^2 r \cos^2 \lambda$$

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$$g_{\text{eff}} = \frac{mg - m\omega^2 r \cos^2 \lambda}{m} = g - \omega^2 r \cos^2 \lambda$$

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For $g_{\text{eff}} = 0$, we have

$$g = \omega_1^2 r \cos^2 \lambda$$

$$\omega_1 = n\omega$$

$$\Rightarrow g = n^2 \omega^2 r \cos^2 \lambda$$

At equator $\lambda = 0$

$$\Rightarrow g = n^2 \omega^2 r$$

Currently $9.8 = g - \omega^2 r$

$$\Rightarrow 9.8 + \omega^2 r = n^2 \omega^2 r$$

$$\Rightarrow n^2 - 1 = \frac{9.8}{\omega^2 r}$$

$$= \frac{9.8}{\left[\frac{2\pi}{24 \times 60 \times 60} \right]^2 \times 6400 \times 10^3}$$

$$= 289.83$$

$$n = 17$$

Q5) <Q2>

Q6) <Q1>

Q7) <Q2>

Q8a) <Q1> , <Q1> Right bank are more eroded in NH.

Q1)

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kepler's laws of planetary motion :

① Planets revolve around the sun in elliptical paths. with sun as one of the foci

Radial vector
from the sun
to a planet
→ sweeps equal

areas in
equal time i.e.
area velocity
is constant

② Areal velocity of the planetary motion is const. i.e. Area of the orbit swept by the planet per unit time remains same.

③ Time period of rotation (T) is proportional to radius of motion as (R)

$$T^2 \propto R^3$$

Now, $F = \frac{GMm}{R^n}$

For circular orbital motion,

$$F = \frac{GMm}{R^n} = \frac{mv^2}{R}$$

$$\Rightarrow v = \sqrt{\frac{GM}{R^{n-1}}}$$

$$\Rightarrow \text{Time Period} = \frac{2\pi R}{v} = \frac{2\pi R \sqrt{R^{n-1}}}{\sqrt{GM}}$$

Now this should be of form $\propto R^{3/2}$

$$\Rightarrow 1 + \frac{n-1}{2} = \frac{3}{2}$$

$$\Rightarrow \boxed{n = 2}$$

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The square of the planets time Period of complete revolution around the sun is proportional to cube of its semi-major axis

(Q2)

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For a satellite motion,

$$\frac{GMm}{R^2} = \frac{mv^2}{R}$$

$$\Rightarrow v = \sqrt{\frac{GM}{R}}$$

$$\Rightarrow \text{Time Period} = \frac{2\pi R}{\sqrt{\frac{GM}{R}}} = \frac{2\pi}{\sqrt{GM}} R^{3/2}$$

$$\text{Also } \frac{GMm}{R_e^2} = mg$$

$$\Rightarrow GM = g R_e^2$$

$$\Rightarrow \text{Time Period} = \frac{2\pi}{\sqrt{g R_e^2}} R^{3/2}$$

$$\Rightarrow T = \frac{2 \times 3.14 \times (3.884 \times 10^8)^{3/2}}{\sqrt{9.8 \times (6400 \times 1000)^2}}$$

$$= \underline{\underline{27.76 \text{ days}}}$$

For geostationary satellite, $T = 1 \text{ day}$

$$\Rightarrow \frac{T_1}{T_2} = \left(\frac{R_1}{R_2}\right)^{3/2}$$

$$\Rightarrow R_2 = R_1 \cdot \left(\frac{T_2}{T_1}\right)^{2/3} = 3.884 \times 10^5 \left(\frac{1}{27.76}\right)^{2/3}$$

$$= \underline{\underline{4.23 \times 10^4 \text{ km}}}$$

$$\approx \underline{\underline{42,300 \text{ km}}}$$

$$\text{Height} = \underline{\underline{35,900 \text{ km}}}$$

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Q3)

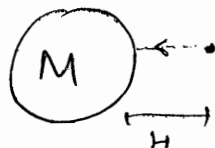
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We know on moon gravitation is $\frac{1}{6}$ th
that of Earth.

$$\Rightarrow \frac{GMm}{R^2} = mg \Rightarrow g = \frac{GM_{\text{Earth}}}{R_{\text{Earth}}^2}$$

$$\frac{g}{G} = \frac{GM_{\text{moon}}}{R_{\text{Moon}}^2}$$

$$\frac{M_{\text{moon}}}{R_{\text{Moon}}^2} = \frac{1}{6} \cdot \frac{M_{\text{Earth}}}{R_{\text{Earth}}^2}$$



From Newton's law,

$$\frac{GM_{\text{moon}} m}{(R_{\text{moon}} + H)^2} - \frac{GM_{\text{Earth}} m}{(d - R_{\text{moon}} - H)^2} = m a$$

~~if $d \gg R_{\text{moon}}$~~

Ignoring earth's effect as $d \gg R_{\text{Moon}}$
we have

$$\frac{GM_{\text{Moon}}}{(R_{\text{moon}} + H)^2} = a = \frac{dv}{dt}$$

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$$\text{Let } R_{\text{moon}} + H = r$$

$$\Rightarrow \boxed{-\frac{GM}{r^2} = \ddot{r}}$$



Multiply both sides by $2\dot{r}$



$$\Rightarrow -\frac{GM}{r^2} \cdot 2\dot{r} = 2\dot{r} \ddot{r}$$

Write acceleration
as $\left(\frac{dv}{dr}\right)$

$$\Rightarrow 2GM \cdot d\left(\frac{1}{r}\right) = d(\dot{r}^2)$$

$$a = \frac{dv}{dt} = v \frac{dv}{dr}$$

$$\Rightarrow 2GM \int d\left(\frac{1}{r}\right) = \int d(\dot{r}^2)$$

$$\Rightarrow \frac{2GM}{r} + C = \dot{r}^2$$

$$\Rightarrow \frac{GM}{r^2} = v \frac{dv}{dr}$$

But at $r \rightarrow \infty$, $\dot{r} = 0$

$$\Rightarrow \int \frac{GM}{r^2} dr = \int v dv$$

$$\Rightarrow C = 0$$

$$\Rightarrow GM \left(\frac{1}{r} - \frac{1}{R}\right) = \frac{v^2}{2}$$

$$\Rightarrow \frac{2GM}{r} = \dot{r}^2$$

$$\Rightarrow \boxed{v \approx \sqrt{\frac{2GM}{R}}}$$

$$\Rightarrow \dot{r} = -\sqrt{\frac{2GM}{r}}$$

[$\dot{r} < 0$ as
 r reduces with
time]

$$\Rightarrow v = \dot{r} = \sqrt{\frac{2GM}{R}} : \text{Speed when falls}$$

$$\boxed{\text{Radius of moon} = 1720 \text{ km}}$$

$$\Rightarrow v = \sqrt{\frac{2 \cdot \frac{g}{6} \cdot \frac{R_m^2}{R_m}}{\frac{R_m}{3}}} = \sqrt{\frac{g R_m}{3}} = \underline{2370 \text{ m/s}}$$

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Q4)

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At new moon



$$\frac{GM_S m_m}{R_{SM}^2} - \frac{GM_E m_m}{R_{EM}^2} = m_m a_{moon}$$

$$\Rightarrow \vec{a}_{moon} = \left[\frac{GM_S}{R_{SM}^2} - \frac{M_E}{R_{EM}^2} \right]$$

Q5)

For $v < v_0$: fall back to Earth $v_0 < v < v_E$: Orbit around Earth $v > v_E$: satellite can go out of Earth's attractive force.

Using the Conservation of energy,

$$k \cdot E_i + P \cdot E_i = k \cdot E_f + P \cdot E_f$$

$$\underline{k \cdot E_i + P \cdot E_i = k \cdot E_f + 0}$$

For v_E minimum, $k \cdot E_f$ should just be 0

$$\Rightarrow \frac{1}{2} m v_i^2 - \frac{GMm}{R_E} = 0$$

$$\Rightarrow v_i = v_{esc} = \sqrt{\frac{2GM}{R_E}}$$

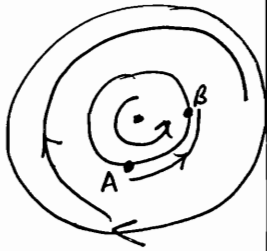
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~~सौर~~ को
new moon
होना है !!

$$G = 6.67 \times 10^{-11}$$

Q4

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$$\text{Time Period of rotation} = \frac{2\pi}{\sqrt{g R_e^2}} \times R^{3/2}$$

$$= \frac{2 \times 3.14 \times (20 \times 10^6)^{3/2}}{\sqrt{9.8 \times (6.4 \times 10^6)^2}}$$

$$= 0.32 \text{ days} = 7.8 \text{ hours}$$

$$\omega_{\text{satellite}} = \frac{2\pi}{T} = \frac{2\pi}{7.8} \text{ (hrs}^{-1}\text{)}$$

$$\omega_{\text{Point}} = \frac{2\pi}{24} \text{ (hrs}^{-1}\text{)}$$

East to west
होना चाहिए!!

$$\omega_{\text{satellite, P}} = 2\pi \left[\frac{1}{7.8} + \frac{1}{24} \right] \text{ hrs}^{-1}$$

$$\Rightarrow T_{\text{satellite, P}} = 5.9 \text{ hrs} \approx 6 \text{ hours}$$

Point on
Earth is
also
revolving

Q6

$$\frac{mv^2}{R} = \frac{GMm}{R^2} \Rightarrow v = \sqrt{\frac{GM}{R}}$$

$$= \sqrt{\frac{6.67 \times 10^{-11} \times 1.9 \times 10^{27}}{3000000}}$$

$$= 42.54 \text{ km/sec}$$

Q8

<Q1>

$$\text{Areal velocity} = \frac{d}{dt} \left(\frac{1}{2} R \times \theta \right)$$

$$= \frac{1}{2} R \times \frac{d\theta}{dt} = \frac{1}{2} R v$$

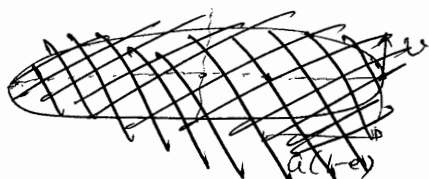
$$= \left(\frac{mv^2}{2m} \right) = \text{const.} = \left(\frac{J}{2m} \right)$$

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Also

$$T = \frac{2\pi ab}{v_r} =$$

$$\frac{2\pi a m b}{\sqrt{\frac{k\mu b^2}{a}}} = \frac{2\pi\sqrt{m} a^{3/2}}{\sqrt{k}} = \propto a^{3/2}$$



$$\frac{J^2}{k\mu} = \frac{b^2}{a} \Rightarrow$$

$$J = m v_r = \sqrt{\frac{k\mu b^2}{a}}$$

Remember
this!!

Q9)

Additional velocity = $v_{esc} - v_{orb}$

Now for v_{esc} , we have total $E = 0$

$$\Rightarrow \frac{1}{2} m v_{esc}^2 - \frac{GMm}{R_e} = 0$$

$$\Rightarrow v_{esc} = \sqrt{\frac{2GM}{R_e}}$$

For v_{orb} near the surface of earth,

$$\frac{GMm}{R^2} = \frac{m v_{orb}^2}{R}$$

$$\Rightarrow v_{orb} = \sqrt{\frac{GM}{R_e}}$$

Additional velocity required = $(\sqrt{2} - 1) \sqrt{\frac{GM}{R_e}}$

$$\text{Also } mg = \frac{GMm}{R^2} \Rightarrow GM = g R_e^2$$

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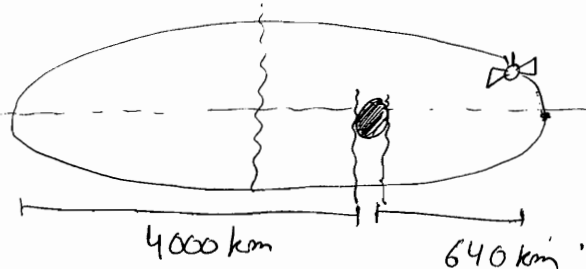
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$$\Rightarrow \text{Additional velocity} = (\sqrt{2}-1) \sqrt{gR_e}$$

$$= \underline{\underline{327 \text{ km/s}}}$$

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(10)



e and a ?

Now

$$a(1-e) = 640 \times 10^3 + 6400 = 7040 \text{ km}$$

$$a(1+e) = 4000 \times 10^3 + 6400 = 10,400 \text{ km}$$

$$\Rightarrow \frac{1+e}{1-e} = \frac{10400}{7040}$$

$$\Rightarrow \frac{2e}{2} = \frac{336}{1744}$$

$$\Rightarrow \frac{2e}{2} = \frac{336}{1744} = 0.192$$

$$\Rightarrow a = \frac{640 \times 10^3}{0.807} = \underline{\underline{8720 \text{ km}}}$$

Do not do
the same
mistake
again &
again.....
It's given
from surface
of earth...
not centre
of earth!!!!

(11)

Let a particle be moving under a
central force F .

Since force is central $\vec{F} = F(r) \hat{r}$

not $F\hat{\theta}$ is 0

$$\Rightarrow \vec{a}_\theta = 0$$

$$\Rightarrow \frac{1}{r} \frac{d}{dt} (r^2 \dot{\theta}) = 0$$

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$$\Rightarrow r^2 \dot{\theta} = \text{const.} \quad \text{--- (1)}$$

Also since Torque is 0, we have
angular momentum as constant

$$\Rightarrow m \dot{\theta} r^2 = \text{const.} \quad \text{--- (2)}$$

From (1) and (2), let

$$\underline{m \dot{\theta} r^2 = J \text{ (const)}}$$

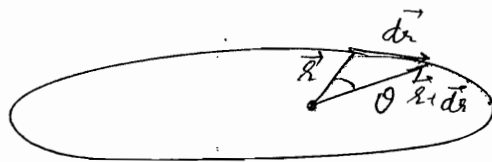
~~In order to find trajectory~~
Equation of motion is

$$\vec{F}(r) = m a_r \hat{r}$$

$$\Rightarrow F(r) = m [\ddot{r} - \dot{\theta}^2 r]$$

$$\Rightarrow \frac{d^2 r}{dt^2} - \left(\frac{d\theta}{dt}\right)^2 r - \frac{F(r)}{m} = 0$$

The areal velocity is the area swept
by the radial vector per unit time,



$$d\vec{A} = \frac{1}{2} \vec{r} \times d\vec{r}$$

$$\frac{d\vec{A}}{dt} = \frac{1}{2} \vec{r} \times \frac{d\vec{r}}{dt} = \frac{1}{2} \frac{\vec{r} \times m \vec{v}}{m} = \underline{\underline{\left(\frac{\vec{J}}{2m}\right)}}$$

$$\Rightarrow \left|\frac{d\vec{A}}{dt}\right| = \underline{\underline{\left(\frac{J}{2m}\right) = \text{const.}}}$$

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(12)

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For an elliptical Orbit,

Time Period T is given by

$$T = \frac{\pi ab 2\mu}{J}$$

Also semilatus Rectum $l = \frac{b^2}{a} = \frac{J^2}{k\mu}$

$$\Rightarrow T = \frac{\pi ab 2\mu}{\sqrt{\frac{b^2}{a} k\mu}}$$

Now $k = GMm$

$$\Rightarrow T = \frac{\pi ab \cdot 2\mu}{\sqrt{\frac{b^2}{a} \cdot GM\mu^2}} = \frac{2\pi ab \sqrt{a}}{b \sqrt{GM}}$$

$$= \frac{2\pi a^{3/2}}{\sqrt{GM}}$$

$$\frac{75 \times 365}{24 \times 60 \times 60} = \frac{2\pi \times a^{3/2}}{\sqrt{6.67 \times 10^{-11} \times 2 \times 10^{36}}}$$

$$a = 2.664 \times 10^{12} \text{ m}$$

$$= 2.664 \times 10^9 \text{ km}$$

Now, ~~ignoring~~

$$a(1-e) = 8.9 \times 10^7$$

$$a(1+e) = 5.239 \times 10^9 \text{ km}$$

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Q13)

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Let the mass of the particle be m .
Since Force is central

$$\vec{a}_0 = 0$$

$$\vec{a}_r = (\ddot{r} - \dot{\theta}^2 r) \hat{r}$$

$$F(r) = m(\ddot{r} - \dot{\theta}^2 r) \quad \text{--- (1)}$$

Also since torque = 0, $\Rightarrow$ angular
momentum remains conserved

$$\Rightarrow m \dot{\theta}^2 r^2 = J \quad (\text{const})$$

$$\Rightarrow \dot{\theta} = \sqrt{\frac{J}{mr^2}}$$

Now, Put $r = \frac{1}{u} \Rightarrow \frac{dr}{du} = -\frac{1}{u^2}$

$$\frac{dr}{dt} = \frac{dr}{d\theta} \cdot \frac{d\theta}{dt} = \sqrt{\frac{J}{mr^2}} \cdot \left(\frac{dr}{du}\right) \cdot \left(\frac{du}{d\theta}\right)$$

$$= -\sqrt{\frac{J}{m}} u^2 \cdot \frac{1}{u^2} \left(\frac{du}{d\theta}\right)$$

$$\boxed{\frac{dr}{dt} = -\frac{J}{m} \left(\frac{du}{d\theta}\right)}$$

$$\frac{d^2 r}{dt^2} = -\frac{J}{m} \cdot \frac{d^2 u}{d\theta^2} \cdot \frac{d\theta}{dt} = -\frac{J^2}{m^2} u^2 \left(\frac{d^2 u}{d\theta^2}\right)$$

From (1),

$$\frac{F}{m} = -\frac{J^2}{m^2} u^2 \left(\frac{d^2 u}{d\theta^2}\right) - \frac{J^2}{m^2} u^3$$

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$$\frac{d^2u}{d\theta^2} + u = -\frac{m \cdot F}{J^2 u^2} \quad \text{--- (2)}$$

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Now, for gravitation like central
force, let $F = -\frac{k}{r^2} = -ku^2$

$$\Rightarrow \frac{d^2u}{d\theta^2} + u = +\frac{m \cdot k}{J^2}$$

$$\Rightarrow u = \frac{mk}{J^2} + A \cos(\theta - \theta_0) \quad \text{Where } A : \text{const.}$$

$$\Rightarrow \frac{(J^2/mk)}{r} = 1 + \left(\frac{AJ^2}{mk}\right) \cos(\theta - \theta_0) \quad \text{--- (3)}$$

Comparing with conic section,

$$\frac{1}{r} = 1 + e \cos(\theta - \theta_0), \text{ we have}$$

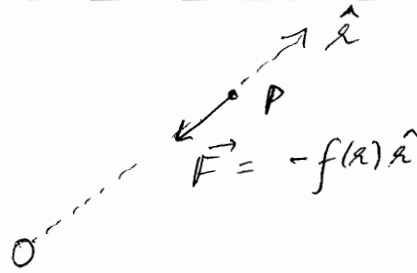
③ As a plane motion equation, we
semi latus rectum = $\left(\frac{J^2}{mk}\right)$ and

$$e = \left(\frac{AJ^2}{mk}\right)$$

Constant A can be found out using
total Energy E.

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2nd solution



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Let the force experienced by the particle be $-f(r) \hat{r}$

Now,

$$\begin{aligned}\vec{L} &= m\vec{r} \times \vec{v} \\ &= m\vec{r} \times (\dot{r}\hat{r} + r\dot{\theta}\hat{\theta}) \\ &= m\vec{r} \times r\dot{\theta}\hat{\theta} \\ &= mr^2\dot{\theta}(\hat{r} \times \hat{\theta})\end{aligned}$$

Now $\vec{L} = \frac{d\vec{L}}{dt}$

Also $\vec{L} \neq \vec{F} \times \vec{r} = -f(r) \hat{r} \times r\hat{r} = 0$

$\Rightarrow \vec{L} = 0$

$\Rightarrow \vec{L} = \text{const.}$

Also as $\vec{L} = mr^2\dot{\theta}(\hat{r} \times \hat{\theta})$

$\therefore \vec{L}_0$ is perpendicular to both $\hat{r}$ and $\hat{\theta}$. As $\vec{L}$ is constant, $\therefore \hat{r}$ and $\hat{\theta}$ always lie in the plane perpendicular to $\vec{L}_0$, which implies that the path of the particle is a planar curve.

Q14)

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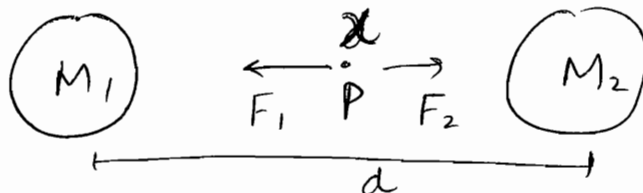
From Kepler 3rd law, $T^2 \propto R^3$

$$\Rightarrow \frac{R_N}{R_E} = \left(\frac{T_N}{T_E} \right)^{2/3} = \left(\frac{165}{1} \right)^{2/3}$$

$$\Rightarrow \boxed{R_N = 30 R_E}$$

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Q15)



At P, $F_1 = F_2$

$$\Rightarrow \sqrt{\frac{M_1}{x^2}} = \sqrt{\frac{M_2}{(d-x)^2}}$$

$$\Rightarrow \frac{d-x}{x} = \sqrt{\frac{M_2}{M_1}}$$

$$\Rightarrow \frac{d}{x} = \frac{\sqrt{M_2} + \sqrt{M_1}}{\sqrt{M_1}}$$

$$\Rightarrow \boxed{x = \frac{\sqrt{M_1} d}{\sqrt{M_2} + \sqrt{M_1}}}$$

$$d-x = \frac{\sqrt{M_2} d}{\sqrt{M_1} + \sqrt{M_2}}$$

$$V = \left(-\frac{GM_1}{x} \right) + \left(-\frac{GM_2}{d-x} \right)$$

~~$$= -\frac{GM_1}{x} - \frac{GM_2}{d-x}$$~~

$$= -\frac{GM_1 (\sqrt{M_1} + \sqrt{M_2})}{\sqrt{M_1} d} - \frac{GM_2 (\sqrt{M_1} + \sqrt{M_2})}{\sqrt{M_2} d}$$

$$= -\frac{G(\sqrt{M_1} + \sqrt{M_2})}{d} [\sqrt{M_1} + \sqrt{M_2}]$$

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$$V = -\frac{G}{d} [M_1 + M_2 + 2\sqrt{M_1 M_2}]$$

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Tut Sheet 4 : Rigid Body Dynamics

Q1)

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For a top, let us consider it as ~~spherically~~ symmetrical body.

i.e. $I_{xx} = I_{yy}$ and cross terms are 0

$$\Rightarrow \boxed{\vec{J} = I \vec{\omega}}$$

Phenomenon of rotation of axis of rotation when a torque acts perpendicular to the axis of rotation is called Precession.....

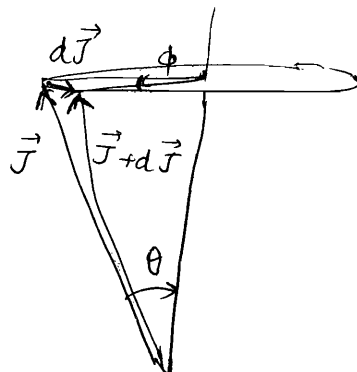


Fig 1

For such a spinning top, let torque τ be applied perpendicular to Angular Momentum. Such a torque will not change the magnitude of Angular Momentum but only its direction. Let such a torque change the angular momentum $\vec{J}$ by $d\vec{J}$, as shown in Fig 1.

$$\Rightarrow \tau = \frac{d\vec{J}}{dt}$$

$$\text{Also } |d\vec{J}| = J \sin \theta \cdot d\phi$$

(from Fig 1).

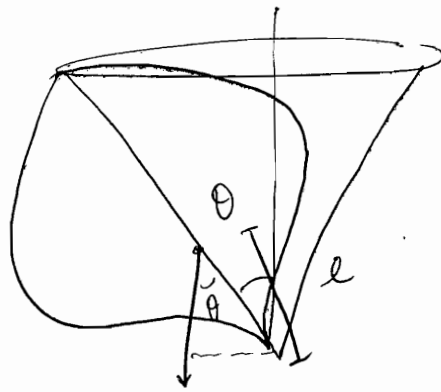
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$$\Rightarrow \tau = \frac{I \sin \theta \, d\phi}{dt} = I \sin \theta \left(\frac{d\phi}{dt} \right)$$

$\left(\frac{d\phi}{dt} \right)$ is defined as precessional
velocity. ω_p .

$$\Rightarrow \omega_p = \frac{\tau}{I \sin \theta}$$



For the shown top,
 $\tau = mgl \sin \theta$

$$\Rightarrow \omega_p = \frac{mgl}{I} = \left(\frac{mgl}{I\omega} \right)$$

Now for solid sphere

$$I = \frac{2}{5} m r^2$$

$$\begin{aligned} \Rightarrow \omega_p &= \frac{mgl}{\frac{2}{5} m r^2 \omega} = \frac{g (5 + 20) \times 10^{-3}}{\frac{2}{5} \times 4 \times 10^{-4} \times 20 \times 2 \times \pi} \\ &= \underline{\underline{12.19 \text{ rad/s}}} \end{aligned}$$

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Q2)

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Q1)

$$\omega_p = \frac{mgl}{I\omega} = \frac{mgl}{mk^2\omega} = \left(\frac{gl}{k^2\omega} \right)$$

$$= \frac{9.8 \times \pi \times 100}{25 \times 20 \times 2\pi} = \underline{\underline{0.98 \text{ rad/s}}}$$

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Q3)

For a single particle, angular momentum $\vec{L}$ is defined as moment of the linear momentum $\vec{p}$

$$\text{i.e. } \vec{L} = \vec{r} \times \vec{p} = \vec{r} \times m\vec{v}$$

For a rotating rigid body, let it be composed of N particles.

$$\Rightarrow \vec{L} = \sum_i \vec{r}_i \times m_i \vec{v}_i$$

$$\text{Also } \vec{v}_i = \vec{\omega} \times \vec{r}_i$$

where $\vec{\omega}$ is the angular velocity of the rotating body

$$\Rightarrow \vec{L} = \sum_i \vec{r}_i \times m_i (\vec{\omega} \times \vec{r}_i)$$

$$= \sum_i m_i [\vec{\omega} (r_i^2) - \vec{r}_i (\vec{r}_i \cdot \vec{\omega})]$$

$$= \sum_i m_i [(\omega_x \hat{i} + \omega_y \hat{j} + \omega_z \hat{k}) (x_i^2 + y_i^2 + z_i^2) - (x_i \hat{i} + y_i \hat{j} + z_i \hat{k}) (\omega_x x_i + \omega_y y_i + \omega_z z_i)]$$

$$\Rightarrow \vec{L}_x = \sum_i m_i [\omega_x (x_i^2 + y_i^2 + z_i^2) - x_i (\omega_x x_i + \omega_y y_i + \omega_z z_i)]$$

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$$\Rightarrow L_x = \sum_i m_i (y_i^2 + z_i^2) \omega_x - \sum_i m_i (x_i y_i \omega_y + x_i z_i \omega_z)$$

$$= I_{xx} \omega_x + I_{xy} \omega_y + I_{xz} \omega_z$$

where

$$I_{xx} = \sum_i m_i (y_i^2 + z_i^2)$$

$$I_{xy} = - \sum_i m_i (x_i y_i)$$

$$I_{xz} = - \sum_i m_i (x_i z_i)$$

Similarly

$$\vec{L} = \vec{I} \vec{\omega}$$

where

$$\vec{I} = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{bmatrix}$$

Time derivative of Angular Momentum is
~~given by~~ given by

$$\text{i.e. } \frac{d\vec{L}}{dt}$$

$$= \frac{d}{dt} \sum (\vec{r}_i \times m_i \vec{v}_i)$$

$$= \sum m_i \left[\frac{d\vec{r}_i}{dt} \times \vec{v}_i \right] + \sum \vec{r}_i \times m_i \left(\frac{d\vec{v}_i}{dt} \right)$$

$$\vec{\tau} = \sum (\vec{r}_i \times \vec{F}_i) = \text{nothing but Torque acting on the body}$$

$$\text{Hence } \vec{\tau} = \frac{d\vec{L}}{dt}$$

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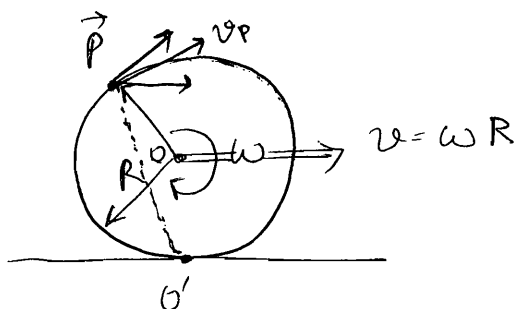
If $\tau = 0$

$$\Rightarrow \frac{d\vec{L}}{dt} = 0 \Rightarrow \vec{L} = \text{const.}$$

i.e. Law of Conservation of Angular Momentum says that in absence of external torque, Angular Momentum of a body remains conserved. It finds application in motion of all bodies under central force where torque = 0

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Q4)



For a rolling body, $v = \omega R$

$$\begin{aligned} \vec{v}_P &= \vec{v}_{P, CM} + \vec{v}_{CM} \\ &= \vec{\omega} \times \vec{r}_{OP} + \vec{\omega} \times \vec{R}_{CM} \\ &= \vec{\omega} \times (R_{CM} + \vec{r}) \\ &= \omega \times \vec{r}_{O'P} \end{aligned}$$

Hence $\vec{v}_P$ is $\perp$ to $\vec{r}_{O'P}$ i.e. equivalent to rotation about point of contact.

kinetic Energy =

$$\begin{aligned} &= \sum \frac{1}{2} m_i v_i^2 \\ &= \sum \frac{1}{2} m_i (\vec{\omega} \times \vec{r}_{O'P})^2 \\ &= \sum \frac{1}{2} m_i \omega^2 d_i^2 \quad [d_i = O'P'] \end{aligned}$$

Prove for
1 scalar
k.E.
and
1 vector
velocity

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$$= \left(\sum \frac{1}{2} m_i d_i^2 \right) \omega^2$$

$$T = \frac{1}{2} I_P \omega^2$$

$$I_P = \sum m_i d_i^2$$

= moment of inertia
about point of contact.

For a circular plate or type (not
sphere)

$$I_P = I_C + m r^2$$

$$\Rightarrow T = \frac{1}{2} (I_C + m r^2) \omega^2$$

$$= \frac{1}{2} I_C \omega^2 + \frac{1}{2} m (\omega r)^2$$

$$T = \frac{1}{2} I_C \omega^2 + \frac{1}{2} m v_C^2$$

First term represents k.E. due to
rotation about axis through COM, and
other term represents translation of COM

Q5) For a body with principal axes x, y, z ,
we have

$$\vec{J} = \omega_x I_x \hat{i} + \omega_y I_y \hat{j} + \omega_z I_z \hat{k}$$

$$(3) \text{ Now } \tau = \left(\frac{dJ}{dt} \right) = \left(\frac{dJ}{dt} \right)_R + (\vec{\omega} \times \vec{J})$$

$$\vec{\omega} \times \vec{J} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \omega_x & \omega_y & \omega_z \\ I_x & I_y & I_z \end{vmatrix} = \sum \hat{i} [\omega_y I_z - \omega_z I_y]$$

$$\Rightarrow N_x = I_x \dot{\omega}_x - (I_y \omega_y \omega_z - I_z \omega_y \omega_z)$$

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①
★ 1st line will
be

$$\vec{N}_{ext,0} = \frac{d\vec{L}_0}{dt}$$

where 0 is
the fixed point
on the body,

$\vec{N}_{ext,0}$ and $\vec{L}_0$
are torques &
angular
momentum of
the body about
the fixed point

★ Now this equation holds for all inertial frames.

Considering an inertial frame S . let S' be the
frame attached to the body $\Rightarrow \left(\frac{d\vec{L}_0}{dt} \right)_S = \left(\frac{d\vec{L}_0}{dt} \right)_{S'} + \vec{\omega} \times \vec{L}_0$

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$$\Rightarrow N_x = I_x \dot{\omega}_x - (I_y - I_z) \omega_y \omega_z$$

$$N_y = I_y \dot{\omega}_y - (I_z - I_x) \omega_z \omega_x$$

$$N_z = I_z \dot{\omega}_z - (I_x - I_y) \omega_x \omega_y$$

These are called Euler's Equations of a rigid body.

For a spherical top, $I_x = I_y = I_z = I$

$$\Rightarrow \vec{N} = I \vec{\omega}$$

Since the reaction at the point of contact exerts no torque, we have

$$\vec{J} = \text{const} = I \vec{\omega}$$

~~For a spherical top, $I_x = I_y = I_z = I$~~

For a symmetrical top, $I_x = I_y \neq I_z$

$$\Rightarrow \text{~~For a spherical top~~}$$

$$I_x \dot{\omega}_x = (I_y - I_z) \omega_y \omega_z$$

$$\dot{\omega}_x = \left(1 - \frac{I_z}{I_x}\right) \omega_z \omega_y$$

define $\left(1 - \frac{I_z}{I_x}\right) \omega_z = \Omega$

$$\Rightarrow \dot{\omega}_x = \left(1 - \frac{I_z}{I_x}\right) \omega_z \omega_y = \Omega \omega_y \quad \text{--- (1)}$$

Also $I_y \dot{\omega}_y = (I_z - I_x) \omega_z \omega_x$

$$\Rightarrow \dot{\omega}_y = \left(\frac{I_z}{I_y} - 1\right) \omega_z \omega_x$$

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$$\Rightarrow \dot{\omega}_y = -\Omega \omega_x \quad - (2)$$

From (1), $\ddot{\omega}_x = \Omega \dot{\omega}_y$
 $= \underline{\underline{-\Omega^2 \omega_x}} \quad (\text{from (2)})$

Similarly $\ddot{\omega}_y = -\Omega \dot{\omega}_x$
 $= \underline{\underline{-\Omega^2 \omega_y}} \quad - (4)$

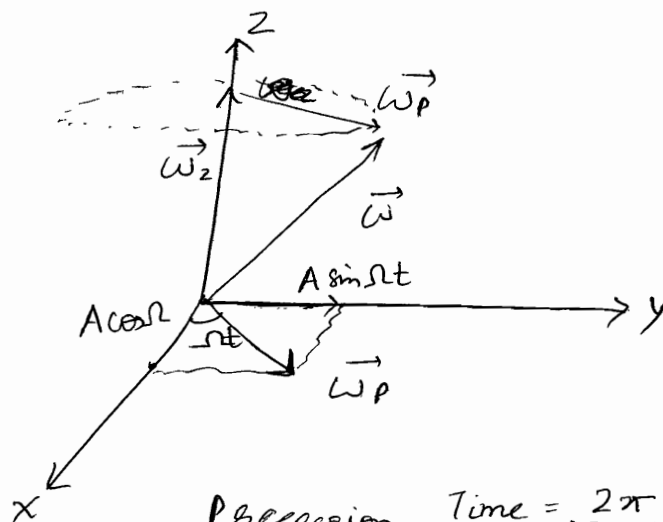
Also $I_z \dot{\omega}_z = (I_x - I_y) \omega_x \omega_y$
 $= 0$

$\Rightarrow \underline{\underline{\omega_z = \text{const}}} \quad - (5)$

$$\vec{\omega} = A \cos \Omega t \hat{i} + A \sin \Omega t \hat{j} + \omega_z \hat{k}$$

$$= \vec{\omega}_p + \omega_z \hat{k}$$

$$|\vec{\omega}_p| = A = \text{const}$$



Precession Time = $\frac{2\pi}{\Omega} = \frac{2\pi}{\omega_z \left(1 - \frac{I_z}{I_x}\right)}$

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— (3)

Q6)

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Elements

$$I_{xx} = \sum (y_i^2 + z_i^2) m_i$$

$$I_{yy} = \sum (z_i^2 + x_i^2) m_i$$

$$I_{zz} = \sum (x_i^2 + y_i^2) m_i$$

are called Moments of Inertia
while the elements

$$I_{xy} = -\sum m_i x_i y_i$$

$$I_{yz} = -\sum m_i y_i z_i$$

$$I_{zx} = -\sum m_i z_i x_i$$

are called Products of Inertia.

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Q7) <Q57

Q8)

<Q37

For a spherical top, $I_{xx} = I_{yy} = I_{zz} = I$

$$\& I_{xy} = I_{yz} = I_{zx} = 0$$

$$\therefore L_x = I\omega_x$$

$$L_y = I\omega_y$$

$$L_z = I\omega_z$$

$$\Rightarrow \vec{L} = I \vec{\omega}$$

But for free rotation $\Rightarrow$ NO Torque

$$\frac{d\vec{L}}{dt} = 0$$

$$\Rightarrow \vec{L} = \text{constant}$$

$$\Rightarrow \underline{\underline{\vec{\omega} = \text{const}}}$$

$\therefore$ The most general free rotation of a spherical top is a uniform rotation about an axis fixed in space & whose direction is given by direction of $\vec{\omega}$.

Q9)

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For principle axes,

$$\vec{J} = I_x \omega_x \hat{i} + I_y \omega_y \hat{j} + I_z \omega_z \hat{k}$$

$$K.E. = \frac{1}{2} \vec{J} \cdot \vec{\omega}$$

$$K = \frac{1}{2} [I_x \omega_x^2 + I_y \omega_y^2 + I_z \omega_z^2]$$

For no external Torque, kinetic energy
will remain same,

$$\Rightarrow \frac{\omega_x^2}{\left(\frac{2K}{I_x}\right)} + \frac{\omega_y^2}{\left(\frac{2K}{I_y}\right)} + \frac{\omega_z^2}{\left(\frac{2K}{I_z}\right)} = 1$$

This is called Ellipsoid of Inertia
or Poincaré Equation of Motion.

Principle axes are those axis of the
body about which the cross terms of
moment of inertia vanish.

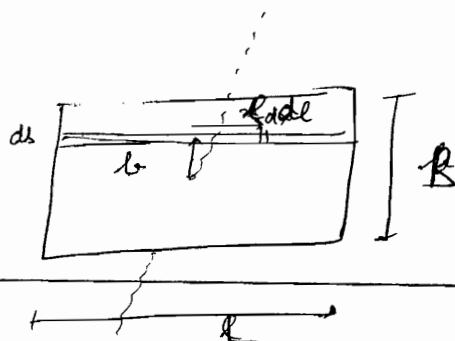
$$\text{i.e. } I_{xy} = I_{yz} = I_{zx} = 0$$

The moments about the principle axes
are called Principal moments viz.

$$I_{xx}, I_{yy}, I_{zz}$$

For the strip shown

$$\begin{aligned} dI &= dm (b^2 + l^2) \\ &= \sigma db dl (b^2 + l^2) \end{aligned}$$



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$$I_{db} = \sigma db \int_{-L/2}^{L/2} (b^2 + l^2) dl$$

$$= \sigma db \cdot \left[b^2 L + \frac{L^3}{12} \right]$$

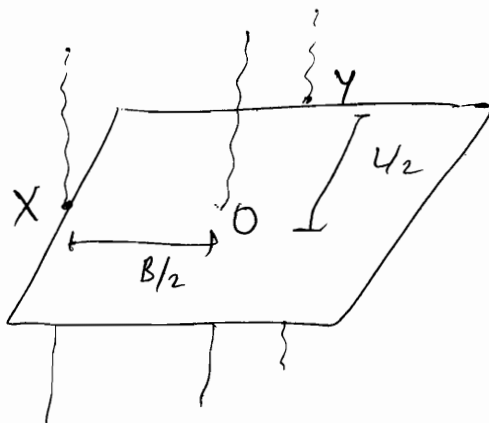
$$I = \int I_{db} = \int_{-B/2}^{B/2} \sigma b^2 L db + \frac{\sigma L^3}{12} \int_{-B/2}^{B/2} db$$

$$= \sigma \cdot L \cdot \frac{B^3}{12} + \frac{\sigma L^3}{12} \cdot B$$

$$= \sigma \cdot \frac{BL}{12} [B^2 + L^2]$$

$$= \frac{M}{BL} \cdot \frac{BL}{12} [B^2 + L^2]$$

$$I = \frac{M}{12} [B^2 + L^2]$$



$$I_x = I_o + \frac{m B^2}{4}$$

$$= \frac{M}{12} [B^2 + L^2] + \frac{M B^2}{4}$$

$$= \frac{M}{12} [4B^2 + L^2]$$

$$I_y = I_o + \frac{m L^2}{4}$$

$$= \frac{M}{12} [4L^2 + B^2]$$

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Q10)

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$$\text{No. of independent variables} = 3 \times 3 = 9$$

(No. of particles * No. of coordinates to describe them)

$$\text{No. of constraint equation} = 3$$

(d_{12}, d_{23}, d_{13} are fixed)

$$\Rightarrow \text{No. of independent variable} = \text{Degree of freedom} = 9 - 3 = 6$$

The Euler Angles ϕ, θ, ψ are defined as:

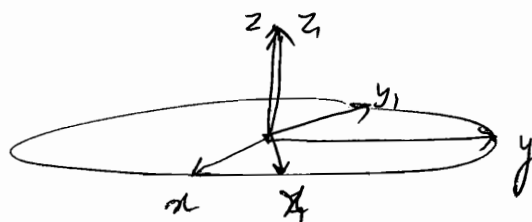
Let x, y, z be the three space axes &

x', y', z' be the three body axes. The

set (x, y, z) is transformed to (x', y', z') by three successive rotations. These angles of the rotations, namely, Euler Angles define the orientation of (x', y', z') w.r.t.

(x, y, z) .

First rotation is about z -axis by angle ϕ , to give x_1, y_1, z_1

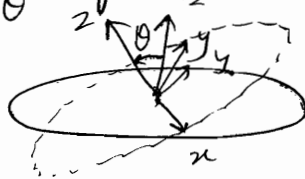


Second rotation is about either of the x or

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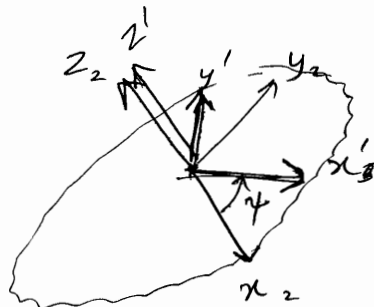
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y_1 axis. say, z_1 axis. to give (x, y, z)
by angle θ



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The third rotation is about z_2 axis
to give (x', y', z') by angle ψ



Q1 Tut 27

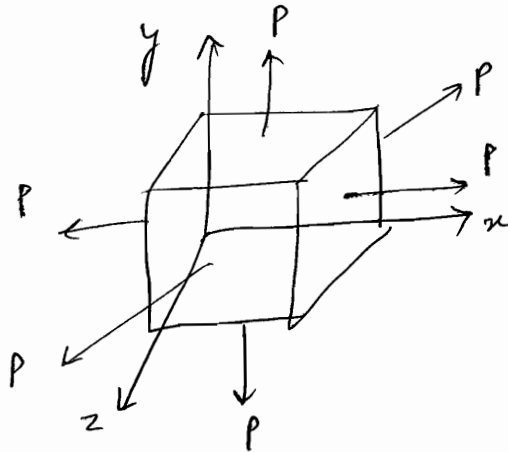
Tut 5 : Mechanics of Continuous Media

Q1)

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(30 marks)

Consider a cube of length l_0 , on which a Force is applied equally on all faces such that a stress P is created.



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Now Young's modulus $Y = \frac{P}{\Delta l/l}$

and Poisson ratio $\sigma = \frac{(\Delta l_{\perp}/l_{\perp})}{(\Delta l_{\parallel}/l_{\parallel})}$

$\Rightarrow$ due to stress along x direction

$$l_x = l_0 + l_0 \frac{P}{Y}$$

$$l_y = l_0 - l_0 \frac{P \sigma}{Y}$$

$$l_z = l_0 - \frac{l_0 P \sigma}{Y}$$

similarly considering stress along all directions

$$l_x = l_0 + \frac{l_0 P}{Y} - \frac{2\sigma l_0 P}{Y}$$

$$l_y = l_0 + \frac{l_0 P}{Y} - \frac{2\sigma l_0 P}{Y}$$

$$l_z = l_0 + \frac{l_0 P}{Y} - \frac{2\sigma l_0 P}{Y}$$

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$$\begin{aligned}\Rightarrow V' &= l_1 l_2 \\ &= \left[l_0 + \frac{dp}{\gamma} (1-2\sigma) \right]^3 \\ &= l_0^3 \left[1 + \frac{3p}{\gamma} (1-2\sigma) \right] \quad (\text{from Binomial Approximation})\end{aligned}$$

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$$\Rightarrow \frac{\Delta V}{V} = \frac{3p}{\gamma} (1-2\sigma)$$

$$\text{Also, } K = \frac{p}{\left(\frac{dV}{V}\right)}$$

$$\Rightarrow \frac{1}{K} = \frac{3}{\gamma} (1-2\sigma)$$

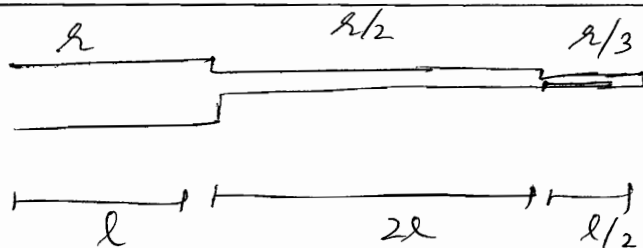
$$\Rightarrow \boxed{\gamma = 3K (1-2\sigma)}$$

Q2)

The flow of the fluid is said to be streamlined or laminar flow, if the velocity at every point in the fluid remains constant, in magnitude as well as direction.

Fluid flow remains streamlined till velocity does not exceed a limiting value called the critical velocity, beyond which flow becomes turbulent.

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Now the flow of liquid through a capillary tube is defined by Poisslieu's formula as

$$Q = \left(\frac{P \pi a^4}{8 \eta l} \right)$$

where a : radius of capillary
 P : Pressure across capillary
 η : viscosity of fluid
 l : length of capillary.

$$\Rightarrow P = \left(\frac{8 Q \eta l}{\pi a^4} \right)$$

Now since the tubes are in series, the flow in each will be same.

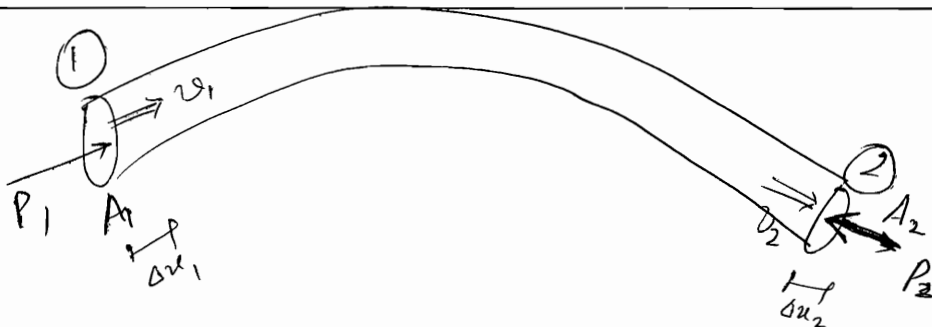
$$\Rightarrow P \propto \left(\frac{l}{a^4} \right)$$

$$\begin{aligned} \Rightarrow P_1 : P_2 : P_3 &= \frac{l_1}{a_1^4} : \frac{l_2}{a_2^4} : \frac{l_3}{a_3^4} \\ &= \frac{l}{r^4} : \frac{2l}{\frac{r^4}{16}} : \frac{l}{\frac{2 \cdot r^4}{81}} \\ &= 1 : 32 : \left(\frac{81}{2} \right) \end{aligned}$$

$$\Rightarrow \underline{P_2 = 32P}, \quad \underline{P_3 = \frac{81}{2}P}$$

3

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Let us consider two cross sections
across the flow of a ~~compressible~~ fluid.
Let the area, velocity of flow, pressure
and their heights be (A_1, v_1, P_1, h_1)
and (A_2, v_2, P_2, h_2)
Let in Δt time, fluid flow be Δx_1
across 1st cross section and Δx_2 across
2nd cross section,

$$\begin{aligned} \text{Work Done} &= F_1 \Delta x_1 - F_2 \Delta x_2 \\ \text{in movement} \\ \text{of liquid} \\ &= P_1 A_1 \Delta x_1 - P_2 A_2 \Delta x_2 \\ &= P_1 A_1 v_1 \Delta t - P_2 A_2 v_2 \Delta t \end{aligned}$$

Since fluid is incompressible

$$\begin{aligned} \Delta m &= A_1 \Delta x_1 \rho = A_2 \Delta x_2 \rho \\ \Rightarrow A_1 v_1 &= A_2 v_2 \end{aligned}$$

This is called Continuity equation.

$$\Rightarrow \Delta W = P_1 \frac{\Delta m}{\rho} - P_2 \frac{\Delta m}{\rho}$$

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Now work done ΔW results in
change in energy

$$\Rightarrow \Delta W = \frac{1}{2} \Delta m (v_2^2 - v_1^2) + \Delta m g (h_2 - h_1)$$

$$\Rightarrow \frac{P_1 \Delta m}{\rho} - \frac{P_2 \Delta m}{\rho} = \frac{1}{2} \Delta m (v_2^2 - v_1^2) + \Delta m g (h_2 - h_1)$$

$$\Rightarrow P_1 + \frac{1}{2} \rho v_1^2 + \rho g h_1 = P_2 + \frac{1}{2} \rho v_2^2 + \rho g h_2$$

= const.

Hence Total energy i.e. sum of
Pressure energy, kinetic Energy and
Potential Energy remains constant
throughout the flow.

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Tut Sheet 6 LTR

Q1)

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For a frames' moving in +x direction
wrt. a frame S, the Lorentz transformation
give out,

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v}{c^2}x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Since the two event are simultaneous in
S $\Rightarrow t_1 = t_2$

$$\Rightarrow t_1' = \frac{t_1 - \frac{v}{c^2}x_1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$t_2' = \frac{t_2 - \frac{v}{c^2}x_2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow t_1' - t_2' = \frac{v}{c^2} \frac{(x_2 - x_1)}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Since events do not take place at a
constant point $\Rightarrow x_1 \neq x_2$

$$\Rightarrow t_1' - t_2' \neq 0 \Rightarrow t_1' \neq t_2' \Rightarrow \text{Not simultaneous.}$$

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Q2)

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$$1 \text{ BeV} = 1 \text{ billion eV}$$

$$= 1 \text{ GeV}$$

$$\text{Now } E = m_0 c^2 + T = \sqrt{p^2 c^2 + (m_0 c^2)^2}$$

$$\text{For proton } m_0 c^2 = 938 \text{ MeV}$$

$$= 0.938 \text{ BeV}$$

$$\Rightarrow 0.938 + 1 = \sqrt{p^2 c^2 + (0.938)^2}$$

$$\sqrt{2.876} = 1.695 \text{ BeV} = p^2 c^2$$

$$\Rightarrow p = 1.695 \text{ BeV}/c$$

$$= \underline{\underline{9.04 \times 10^{-19} \text{ kg m/s}}}$$

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Q3)

From Lorentz transform, for a
frame S' travelling in $+x$ direction with
velocity v wrt. frame S , we have

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v}{c^2} x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

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$$\Rightarrow x'^2 + y'^2 + z'^2 - c^2 t'^2$$

$$= \frac{(x - vt)^2 - c^2 \left(t - \frac{v}{c^2} x\right)^2}{\left[1 - \frac{v^2}{c^2}\right]} + y^2 + z^2$$

$$= \frac{x^2 + v^2 t^2 - 2xvt - c^2 t^2 - \frac{v^2}{c^2} x^2 + 2tvx}{1 - \frac{v^2}{c^2}} + y^2 + z^2$$

$$= \frac{x^2 \left[1 - \frac{v^2}{c^2}\right] - c^2 t^2 \left[1 - \frac{v^2}{c^2}\right] + y^2 + z^2}{1 - \frac{v^2}{c^2}}$$

$$= x^2 + y^2 + z^2 - c^2 t^2$$

$\therefore x^2 + y^2 + z^2 - c^2 t^2$ is Lorentz invariant

Q4) From inverse Lorentz transform

$$x = \frac{x' + vt'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$t = \frac{t' + \frac{v}{c^2} x'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y = y'$$

$$z = z'$$

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$$\Rightarrow x = \frac{60 + \frac{3c}{5} \cdot 8 \times 10^{-8}}{\sqrt{1 - \left(\frac{3}{5}\right)^2}}$$

$$= \frac{60 + \frac{3}{5} \times 3 \times 8}{\left(\frac{4}{5}\right)} = \underline{\underline{93 \text{ m}}}$$

$$t = \frac{8 \times 10^{-8} + \frac{3c}{5c^2} \times 60}{\sqrt{1 - \left(\frac{3}{5}\right)^2}} \Delta$$

$$= \frac{8 \times 10^{-8} + 12 \times 10^{-8}}{\left(\frac{4}{5}\right)} \Delta$$

$$= \frac{20}{4} \times 5 \times 10^{-8} \Delta$$

$$= \underline{\underline{25 \times 10^{-8} \Delta}}$$

Q5) From Lorentz transforms, for a frame.....

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v}{c^2} x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

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Taking differentials

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$$\Rightarrow dx' = \frac{dx - v dt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$dy' = dy$$

$$dz' = dz$$

$$dt' = \frac{dt - \frac{v}{c^2} dx}{\sqrt{1 - \frac{v^2}{c^2}}}$$

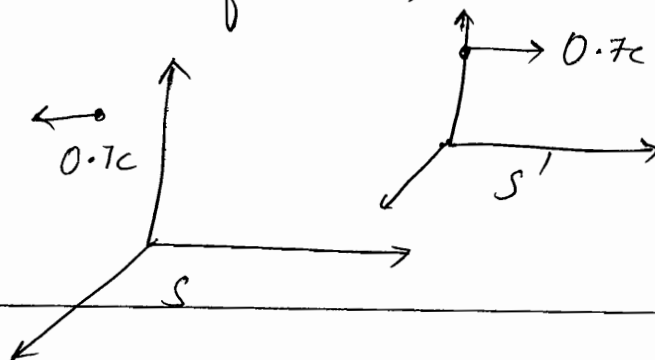
From these equations

$$u_x' = \left(\frac{dx'}{dt'} \right) = \frac{dx - v dt}{dt - \frac{v}{c^2} dx} = \left(\frac{u_x - v}{1 - \frac{v u_x}{c^2}} \right)$$

$$u_y' = \frac{dy'}{dt'} = \frac{dy \sqrt{1 - \frac{v^2}{c^2}}}{dt - \frac{v}{c^2} dx} = \left(\frac{u_y \sqrt{1 - \frac{v^2}{c^2}}}{1 - \frac{v u_x}{c^2}} \right)$$

$$u_z' = \frac{dz'}{dt'} = \left(\frac{u_z \sqrt{1 - \frac{v^2}{c^2}}}{1 - \frac{v u_x}{c^2}} \right)$$

These are referred to as addition
theorems of velocity.



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Consider two frames S and S' where
 S' moves at $v = 0.7c$ wrt frame S .
Let us consider a particle moving
in frame S with $u_x = -0.7c$,
 u_x' of the particle wrt frame S' will
be,

$$u_x' = \frac{-0.7c - 0.7c}{1 + \frac{0.7c \cdot 0.7c}{c^2}}$$

$$= \frac{-1.4c}{1 + 0.49}$$

$$= \underline{\underline{-0.939c}}$$

$\Rightarrow$ By an analogy, $v_{A,B} = -0.939c$
ie. they will see each other as
moving away with $v = 0.939c$

Q6)

The half life of μ meson in its own
frame of reference (proper time) is
 $2 \times 10^{-8} s$

When observed by the person on a
stationary ~~frame~~ frame, like on
earth, time dilation will be observed.

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$$\Rightarrow T_{\text{earth}} = \frac{2 \times 10^{-8}}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{2 \times 10^{-8}}{\sqrt{1 - \frac{1}{4}}}$$

$$= \underline{\underline{2.31 \times 10^{-8} \text{ s}}}$$

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Q7)

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

If $m = 2m_0$

$$\Rightarrow 2 = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \Rightarrow 1 - \frac{v^2}{c^2} = \frac{1}{4}$$

$$\Rightarrow \frac{v}{c} = \frac{\sqrt{3}}{2} \Rightarrow \boxed{v = \frac{\sqrt{3}}{2} c}$$

Q8) For a frames' moving at velocity v wrt. frame S at a speed of v in $+x$ direction, we have by Lorentz equations

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}} ; y' = y ; z' = z ;$$

$$t' = \frac{\left(t - \frac{v}{c^2} x\right)}{\sqrt{1 - \frac{v^2}{c^2}}}$$

For $\frac{v}{c} \ll 1$ $1 - \frac{v^2}{c^2} \approx 1$

$$\Rightarrow \left. \begin{aligned} x' &= x - vt \\ t' &= t \end{aligned} \right\} \begin{array}{l} \text{these are} \\ \text{Galilean} \\ \text{transformations} \end{array}$$

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The physical significance of Lorentz Transformations are :

- ① They preserve speed of light in addition of velocities.
- ② Using Lorentz transformation, length contraction, time dilation & relativity of simultaneity can be proved.

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(Q9) For a body ^(cube) moving at velocity v wrt frame S in body frame of reference S' , we have

$$\text{density} = \frac{m_0}{l_0^3}$$

In frame of S , we have

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$l_x = l_0 \sqrt{1 - \frac{v^2}{c^2}}$$

$$l_y = l_0$$

$$l_z = l_0$$

$$\begin{aligned} \Rightarrow \text{density} &= \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}} \cdot l_0^3 \sqrt{1 - \frac{v^2}{c^2}}} \\ &= \frac{d_0}{\left(1 - \frac{v^2}{c^2}\right)} \end{aligned}$$

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Q10)

Hence density increases for a moving body.

$$\frac{19.3 \times 10^3 \text{ kg/m}^3}{(\sqrt{1-(0.9)^2})^2} = \underline{\underline{101.57 \times 10^3 \text{ kgm}^{-3}}}$$

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Q11)

Let the ~~relative~~ increase in velocity be dv . The ~~relative~~ increase in Energy

$$dE = d\left(\frac{m_0 c^2}{\sqrt{1-\frac{v^2}{c^2}}}\right)$$

$$= \frac{1}{\cancel{2}} \frac{m_0 c^2}{\left(1-\frac{v^2}{c^2}\right)^{3/2}} \cdot \frac{-1}{c^2} \cdot 2v dv$$

$$= \frac{E}{\left(1-\frac{v^2}{c^2}\right)} \cdot \frac{v}{c^2} dv$$

$$dE = \frac{E v}{(c^2 - v^2)} dv$$

~~$$\frac{dE}{dv} = \left(\frac{E v}{c^2 - v^2}\right)$$~~

$$\boxed{\left(\frac{dE}{E}\right) = \left(\frac{v^2}{c^2 - v^2}\right) \left(\frac{dv}{v}\right)}$$

$\therefore$ relative increase in Energy $\left(\frac{dE}{E}\right)$ is $\left(\frac{v^2}{c^2 - v^2}\right)$ times the relative increase in velocity.

⊛ Change
= Δx

relative change

$$= \left(\frac{\Delta x}{x}\right)$$

Q12)

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$$0.51 + 1 = \sqrt{p^2 c^2 + (0.51)^2}$$

~~1.51 MeV = mc^2~~

~~1.51 MeV = mc^2~~

~~1.51 MeV = mc^2~~

⊛

Always prefer

$$E = mc^2$$

$$E = \sqrt{p^2 c^2 + (mc^2)^2}$$

$$\Rightarrow pc = 1.42 \text{ MeV}$$

$$\Rightarrow p = (1.42 \text{ MeV}/c)$$

$$\text{Also } 1.51 \text{ MeV} = mc^2$$

$$\Rightarrow m = \frac{1.51 \text{ MeV}}{c^2}$$

$$\Rightarrow \text{speed} = \frac{p}{m} = \frac{1.42 \text{ MeV}/c}{1.51 \text{ MeV}/c^2} = 0.94c = 2.82 \times 10^8$$

$$1.51 \text{ MeV} = \frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow v = 0.94c$$

Q13)

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{25 - 0.6 \cdot 5 \times 10^{-8}}{\sqrt{1 - (0.6)^2}} = \frac{16}{0.8} \text{ m} = 20 \text{ m}$$

$$t' = \frac{t - \frac{v}{c^2} x}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{5 \times 10^{-8} - \frac{0.6}{3 \times 10^8} \cdot 25}{0.8} = 0$$

$$S'(20, 0, 0, 0)$$

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Q14)

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Q15)

Q-97 931 MeV

From frame 1 bound to earth

$$t = \frac{6 \times 10^3}{0.99 \times 3 \times 10^8} = \underline{\underline{2.02 \times 10^{-5} \text{ s}}}$$

From frame bound to meson (proper time)

$$t_0 = 2.02 \times 10^{-5} \sqrt{1 - (0.99)^2}$$

$$= \underline{\underline{2.84 \times 10^{-6} \text{ s}}}$$

Q16)

1 eV = $1.6 \times 10^{-19} \text{ J}$ $\otimes$ 1 eV is the energy
acquired by an e- when it is
accelerated through 1V
potential.

$$1.3 \text{ MeV} = mc^2$$

$$\Rightarrow m = 1.3 \left(\frac{\text{MeV}}{c^2} \right)$$

$$\text{Also } 1.3 \text{ MeV} = \sqrt{p^2 c^2 + (0.5)^2} \text{ MeV}$$

$$\Rightarrow 1.195 \text{ MeV} = pc$$

$$\Rightarrow p = 1.195 \left(\frac{\text{MeV}}{c} \right)$$

$$\Rightarrow v = \frac{p}{m} = \frac{1.195 \frac{\text{MeV}}{c}}{1.3 \frac{\text{MeV}}{c^2}}$$

$$= \underline{\underline{2.75 \times 10^8 \text{ m/s}}}$$

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$$\frac{p}{m} = \frac{\frac{1.195 \text{ MeV}}{c}}{\frac{1.3 \text{ MeV}}{c^2}} = v$$

$$v = 2.75 \times 10^8 \text{ m/s}$$

① H₂
② Michelson Morley

③ Davison Germer (wave particle duality)
④ Stern Gerlach ($s = 1/2$)
⑤ Yang & Thir Ex. (Parity Violation)

Q17)

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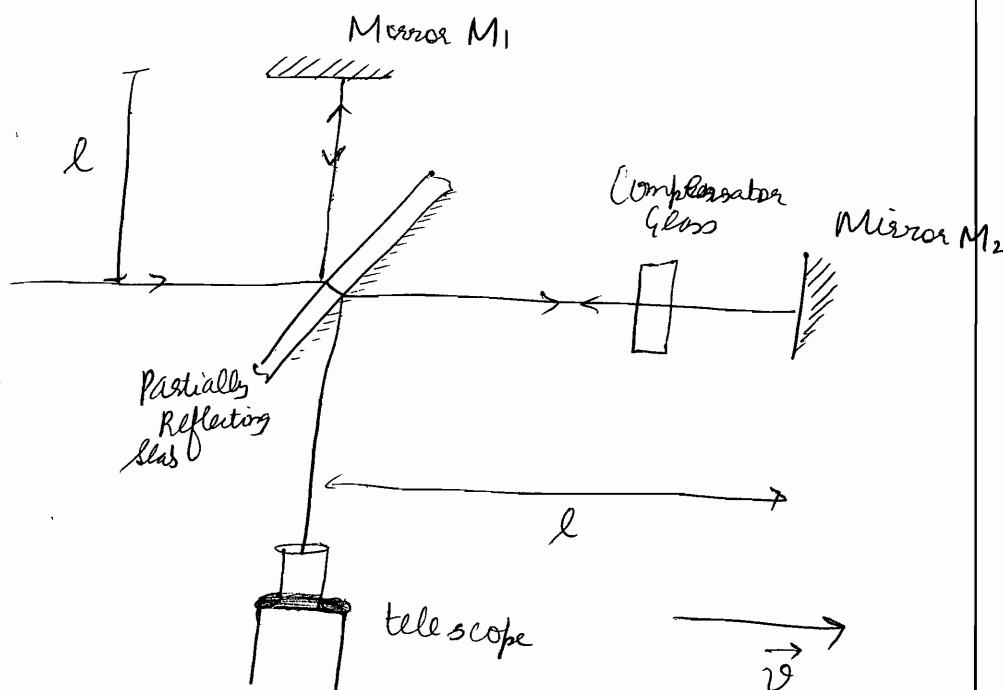
Q18)

Q-3

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Null Result of Michelson Morley Experiment

Michelson Morley Experiment negated the existence of ether and is explicit proof of constancy of speed of light. The set up of the experiment uses Michelson Interferometer. (as shown below)



Michelson interferometer is based upon the principle that when path difference b/w waves travelling in perpendicular directions exist, a fringe pattern is formed. When the instrument is rotated through 90° , the path difference reverses and fringe shift in the pattern should be observed. But

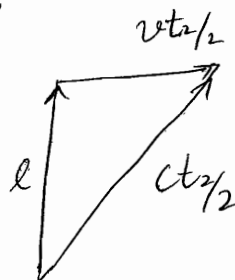
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no fringe shift was observed. The
result of 0 fringe shift is called Null
Result. Einstein's STR successfully
explained the results.

The path difference can be calculated
as. For Horizontal wave,

$$\begin{aligned}
 t_1 &= \frac{l}{c+v} + \frac{l}{c-v} = \left(\frac{2cl}{c^2-v^2} \right) \\
 &= \frac{2l}{c} \left[1 - \frac{v^2}{c^2} \right]^{-1} \\
 &= \underline{\underline{\frac{2l}{c} \left[1 + \frac{v^2}{c^2} \right]}}
 \end{aligned}$$

For vertical ray,



$$\left(\frac{ct_2}{2} \right)^2 = \left(\frac{vt_2}{2} \right)^2 + l^2$$

$$\Rightarrow t_2^2 [c^2 - v^2] = 4l^2$$

$$\Rightarrow t_2 = \frac{2l}{\sqrt{c^2 - v^2}} = \underline{\underline{\frac{2l}{c} \left[1 + \frac{v^2}{2c^2} \right]}}$$

$$\Delta x = c(t_1 - t_2)$$

$$= 2l \left[\frac{v^2}{2c^2} \right] = \underline{\underline{\left(\frac{lv^2}{c^2} \right)}}$$

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when the experiment is rotated,

$$\Delta x = \frac{lv^2}{c^2}$$

Hence we should get fringe shift
corresponding to $\left(\frac{2lv^2}{c^2}\right)$

But Here, no shift was observed and the
explanation was constancy of velocity of
light, as given by Einstein.

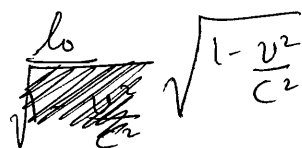
$$t_1 = t_2 = \left(\frac{2l}{c}\right)$$

Q19) If A length as seen in the frame at rest
wrt. the body is l_0 . If the same
length is viewed from a frame moving
at velocity v wrt. body, the length will
appear to be smaller by a factor of
 $\sqrt{1 - \frac{v^2}{c^2}}$

i.e.
$$l = l_0 \sqrt{1 - \frac{v^2}{c^2}}$$

This is called Lorentz Contraction!!

For an observer, its length will be

 $l_0 \sqrt{1 - \frac{v^2}{c^2}}$

$$\therefore \text{time to pass} = \frac{l_0}{v} \sqrt{1 - \frac{v^2}{c^2}} = 4 \mu s$$

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$$\Rightarrow \frac{100}{v} \sqrt{1 - \frac{v^2}{c^2}} = 4 \times 10^{-6}$$

$$\Rightarrow 100 \sqrt{\frac{1}{v^2} - \frac{1}{c^2}} = 4 \times 10^{-6}$$

$$\Rightarrow \sqrt{\frac{1}{v^2} - \frac{1}{c^2}} = 4 \times 10^{-8}$$

$$\frac{1}{v^2} - \frac{1}{c^2} = 1.6 \times 10^{-15}$$

$$\frac{1}{v^2} = 1.11 \times 10^{-17} + 1.6 \times 10^{-15}$$

$$\frac{1}{v^2} = 1.611 \times 10^{-15}$$

$$v = 0.249 \times 10^8 \text{ m/s}$$

Q20)

For an object moving with velocity v wrt. frame S , the mass m is given as

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

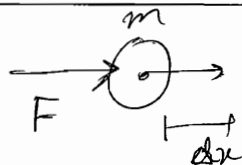
Where m_0 is the mass of particle at rest

$$\begin{aligned} \Rightarrow dm &= \frac{m_0}{\left(1 - \frac{v^2}{c^2}\right)^{3/2}} \cdot \frac{-1}{2} \cdot \frac{-2v}{c^2} \cdot dv \\ &= \frac{m}{1 - \frac{v^2}{c^2}} \cdot \frac{v}{c^2} dv = \frac{m v dv}{c^2 - v^2} \end{aligned}$$

Now consider the object being applied upon by a Force F s.t. Work Done

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$$\begin{aligned}
 dW &= F \cdot dx \\
 &= \frac{d(p)}{dt} dx \\
 &= \frac{dm}{dt} \cdot v dx + m \cdot \frac{dv}{dt} dx \\
 &= v^2 dm + m \cdot v dv \\
 &= v^2 dm + dm (c^2 - v^2) \\
 &= c^2 dm
 \end{aligned}$$

Now $dW = \Delta h \cdot E$

$$\Rightarrow \int dT = \int_{m_0}^{m} c^2 dm$$

$$\Rightarrow T = c^2(m - m_0)$$

$$\begin{aligned}
 \Rightarrow E &= T + \text{rest mass energy} \\
 &= c^2(m - m_0) + (m_0 c^2)
 \end{aligned}$$

$$\Rightarrow \boxed{E = mc^2}$$

Q21)

$$M_{\mu_0} = 207 Me$$

$$\text{Proper Mean life time} = 2.2 \times 10^{-6} s$$

$$\text{Time measured in lab} = \frac{T_0}{\sqrt{1 - \frac{v^2}{c^2}}} = 6.9 \times 10^{-6}$$

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$$\frac{2.2 \times 10^{-6}}{6.9 \times 10^{-6}} = \sqrt{1 - \frac{v^2}{c^2}}$$

$$v = 0.948c$$

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$$\text{Speed of muons in lab} = 2.84 \times 10^8 \text{ m/s}$$

Q22)



The observation on the ship will
observe simple rotation of rod.

The observed on the earth will [जटिल]
also experience an increase in the
length of the rod by $\frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$ factor

$$l_{\text{initial}} = 0.5 \sqrt{1 - \frac{v^2}{c^2}} \text{ m}$$

$$l_{\text{final}} = 0.5 \text{ m}$$

Q23)

<Q-57

$$\frac{u_x + v}{1 + \frac{u_x v}{c^2}} = \frac{1.8c}{1 + 0.81} = \underline{\underline{0.99c}}$$

Q24)

Let B be moving at speed +v wrt. A

$$\Rightarrow x_B = \left[\frac{x_A - v_B t_A}{\sqrt{1 - \left(\frac{v_B}{c}\right)^2}} \right]; y_B = y_A; z_B = z_A;$$

$$t_B = \left[\frac{t_A - \frac{v_B x_A}{c^2}}{\sqrt{1 - \frac{v_B^2}{c^2}}} \right]$$

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$$t_1 - t_2 = 10^{-6} \text{ s}$$

$$t_1' - t_2' = \frac{(t_1 - t_2) + 0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow 10^{-6} \times 3 = \frac{10^{-6}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow \left[1 - \frac{v^2}{c^2}\right] = \frac{1}{9} \Rightarrow v = \sqrt{\frac{8}{9}} c$$

$$x_1 - x_2 = 0$$

$$x_1' - x_2' = \frac{(x_1 - x_2) - (v)(t_1 - t_2)}{\sqrt{1 - \frac{v^2}{c^2}}} \checkmark$$

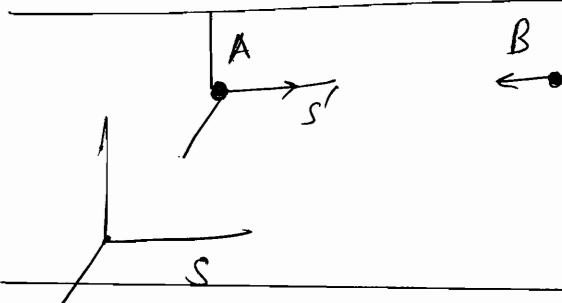
$$= - \frac{\sqrt{\frac{8}{9}} c \cdot 10^{-6}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$= \frac{\sqrt{\frac{8}{9}} c \times 10^{-6}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$= \frac{2545}{3} \text{ m}$$

$$= \underline{\underline{848.52 \text{ m}}}$$

Q 25)



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Let us attach frame S' to Body A. This
implies velocity of frame $S' \Rightarrow v = \left(\frac{4c}{5}\right)$

Writing Lorentz equation

$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$t' = \frac{t - \frac{v}{c^2} x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow u_{x'} = \frac{dx'}{dt'} = \frac{dx - v dt}{dt - \frac{v}{c^2} dx}$$

$$= \left(\frac{u_x - v}{1 - \frac{v u_x}{c^2}} \right)$$

Now let us consider motion of particle
B.

wrt S , $u_B = -\left(\frac{4c}{5}\right)$

$$\Rightarrow u_{B,A} = u_{B,S'} = u_{x'} = \frac{-\frac{4c}{5} - \frac{4c}{5}}{1 + \frac{\frac{4c}{5} \cdot \frac{4c}{5}}{c^2}}$$

$$\Rightarrow u_{x'} = -\frac{8}{5} \frac{c}{1 + \frac{16}{25}} = -\frac{40c}{41} = -0.97c$$

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In Frame S

$$\begin{aligned}
 h \cdot E \cdot B &= m_0 c^2 - m_0 c^2 \\
 &= \frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}} - m_0 c^2 = m_0 c^2 \left[\frac{1}{\sqrt{1 - \left(\frac{4}{5}\right)^2}} - 1 \right] \\
 &= m_0 c^2 \left[\frac{5}{3} - 1 \right] = \frac{2}{3} m_0 c^2
 \end{aligned}$$

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$$\begin{aligned}
 h \cdot E \cdot B' &= m_0 c^2 \left[\frac{1}{\sqrt{1 - \left(\frac{40}{41}\right)^2}} - 1 \right] \\
 &= m_0 c^2 \left[\frac{41}{9} - 1 \right] \\
 &= m_0 c^2 \cdot \left(\frac{32}{9} \right) \\
 &= \frac{32}{9} m_0 c^2
 \end{aligned}$$

Q26) In Non relative physics, velocity is not const. but relative

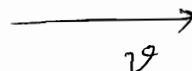
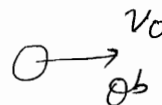
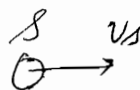
⇒ Source and observer will encounter different velocities of wave
hence doppler effect is given as

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$$\gamma = \gamma_0 \sqrt{\frac{v-v_0}{v-v_s}}$$

⊛ Here it matters if source
moves wrt. observer or observer
moves wrt. source.

⊛ Asymmetric



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In relativistic phenomenon, velocity of
light is same for both source and observer.

$$\text{here } \gamma = \gamma_0 \sqrt{\frac{1 - \frac{v}{c}}{1 + \frac{v}{c}}}$$

⊛ Here it does not matter whether
source moves relative to observer
or vice versa as there is no way
out.

⊛ Symmetric

⊛ Even transverse doppler effect is
observed which has no classical
analogue.

Longitudinal Doppler Effect

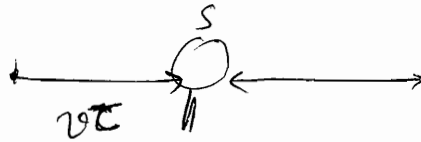
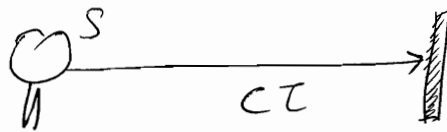
$$\propto T' = T + \frac{\Delta x}{v}$$

$$\propto T' = T + \frac{\Delta x + v \Delta t'}{v}$$

$$\propto T' \left[1 - \frac{v}{c} \right] = T$$

$$\gamma' = \gamma \sqrt{\frac{1 - \frac{v}{c}}{1 + \frac{v}{c}}}$$

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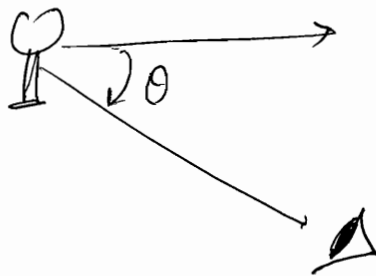


$$\lambda = (c - v) T = (c - v) \alpha T_0$$

$$\Rightarrow \frac{c}{\lambda} = c \left(1 - \frac{v}{c}\right) \propto \frac{1}{\lambda_0}$$

$$\Rightarrow \lambda = \lambda_0 / \left(1 - \frac{v}{c}\right) \propto \lambda_0 \sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}}$$

Transverse Doppler



$$\lambda = (c - v \cos \theta) T$$

$$\frac{c}{\lambda} = c \left[1 - \frac{v}{c} \cos \theta\right] \propto \frac{1}{\lambda_0}$$

$$\Rightarrow \lambda = \lambda_0 / \left[1 - \frac{v}{c} \cos \theta\right]$$

$$= \lambda_0 \sqrt{1 - \frac{v^2}{c^2}} \left(1 - \frac{v}{c} \cos \theta\right)$$

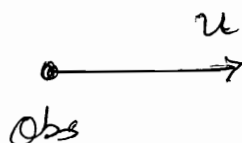
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⇒ For $\theta = 90^\circ$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

In the given problem



In observer frame,



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γ



From conservation of momentum,
 $p_\mu = p_\gamma$ — (1)

~~$\Rightarrow p_\mu = (E_\gamma/c)$ — (1)~~

~~Also $E_\gamma + m_\mu c^2 = m_\pi c^2$ — (2)~~

~~$E_\gamma^2 + m_\mu^2 c^4 = m_\pi^2 c^4$~~

~~$m_\mu^2 c^4 + m_\mu^2 c^4 = m_\pi^2 c^4$~~

~~$\Rightarrow \pi^2 c^4 + E_\gamma^2 = m_\pi^2 c^4$~~

$E_\mu + E_\gamma = \pi c^2$ — (2)

$E_\mu^2 = p_\mu^2 c^2 + m_\mu^2 c^4$

$E_\gamma^2 = p_\gamma^2 c^2$

$\Rightarrow E_\mu^2 - E_\gamma^2 = m_\mu^2 c^4$

$\Rightarrow E_\mu - E_\gamma = \frac{m_\mu^2 c^4}{\pi c^2} = \frac{m_\mu^2 c^2}{\pi}$

$\Rightarrow E_\mu = \frac{1}{2} \left[\pi c^2 + \frac{m_\mu^2 c^2}{\pi} \right]$

Note this
standard
way to solve

$E_a + E_b$
 $E_a^2 - E_b^2$
 $E_a - E_b$

$\Rightarrow E_a$
 E_b

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$$\Rightarrow h \cdot E_{\mu} = E_{\mu} - \mu c^2$$

$$= \frac{\pi c^2}{2} + \frac{\mu^2 c^2}{2\pi} - \mu c^2$$

$$= \frac{\pi^2 c^2 + \mu^2 c^2 - 2\pi \mu c^2}{2\pi}$$

$$h \cdot E_{\mu} = \frac{c^2 (\pi - \mu)^2}{2\pi}$$

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Q28)

Q20

Q30)

Let a consider a rod of length L_0 travelling at velocity v relative to a frame S . Let us attach frame S' to the moving body.

From Lorentz transformation,

$$x' = \frac{x - vt}{\sqrt{1 - v^2/c^2}}$$

In order to measure the length correctly, the measurements of two end points must be simultaneously from frame S

$$\text{i.e. } t_1 = t_2$$

$$\Rightarrow x'_1 = \frac{x_1 - vt_1}{\sqrt{1 - v^2/c^2}}$$

$$x'_2 = \frac{x_2 - vt_2}{\sqrt{1 - v^2/c^2}}$$

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$$\Rightarrow (x_1' - x_2') = \frac{(x_1 - x_2)}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Since measurement made from frame S'
that is at rest w.r. rod,
 $x_1' - x_2'$ is the proper length L_0 .

$$\Rightarrow L_0 = \frac{L_{obs}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

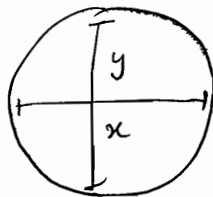
$$\Rightarrow \boxed{L = L_0 \left(1 - \frac{v^2}{c^2}\right)^{1/2}}$$

Also note that there is no affect on
length L to motion of particle as,

$$dy' = dy$$

$$dz' = dz$$

Hence for spherical ball



$x = y = d$ in frame of S'

$$x \rightarrow x \sqrt{1 - \frac{v^2}{c^2}} = d \sqrt{1 - \frac{v^2}{c^2}}$$

$$y \rightarrow y = d$$

$$z \rightarrow z = d$$

Hence ball will appear as spheroid

~~spheroid~~ or prolate or ellipsoid.

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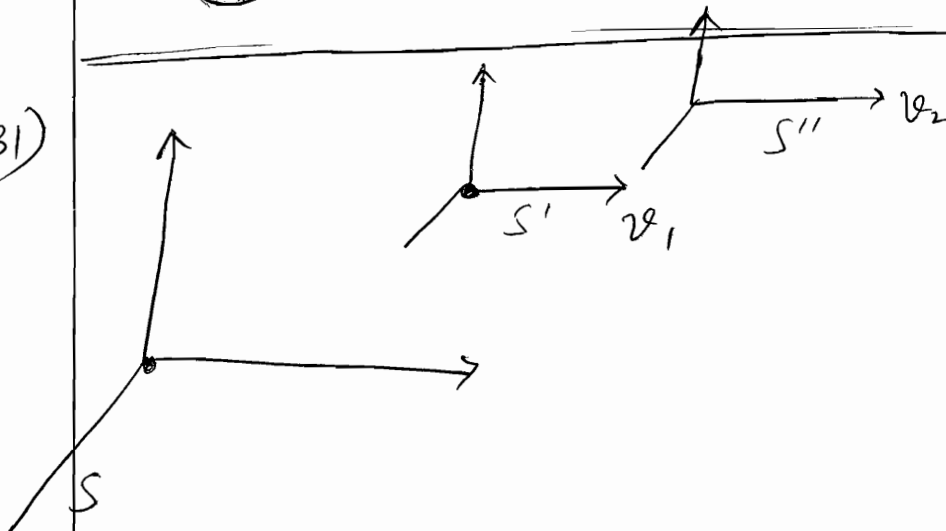
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Prolate Spheroid

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Q31)



Let frame S' moves at velocity v_1
wrt. S in $\oplus$ ve x direction.

$$\Rightarrow x' = \frac{x - vt}{\sqrt{1 - v^2/c^2}}$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v}{c^2}x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\sqrt{1 - \frac{v^2}{c^2}}$$

Now let us suppose S'' frame
moves wrt. S' at velocity v_2
in $\oplus$ ve x direction

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$$\Rightarrow x'' = \frac{x' - v_2 t'}{\sqrt{1 - \frac{v_2^2}{c^2}}}$$

$$y'' = y'$$

$$z'' = z'$$

$$t'' = \frac{t' - \frac{v_2}{c^2} x'}{\sqrt{1 - \frac{v_2^2}{c^2}}}$$

$$\text{Now, } x'' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}} - \frac{v_2 \cdot \left(t - \frac{v}{c^2} x\right)}{\sqrt{1 - \frac{v_2^2}{c^2}}}$$

$$= \frac{x \left[1 + \frac{vv_2}{c^2} \right] - t \left[\frac{v + v_2}{\sqrt{1 - \frac{v^2}{c^2}}} \right]}{\sqrt{1 - \frac{v_2^2}{c^2}}}$$

$$= \frac{x - t \left(\frac{v + v_2}{1 + \frac{vv_2}{c^2}} \right)}{\frac{\sqrt{1 - \frac{v^2}{c^2}} \sqrt{1 - \frac{v_2^2}{c^2}}}{\left(1 + \frac{vv_2}{c^2} \right)}}$$

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$$U = \left(\frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}} \right) \leftarrow \text{We have to reach here...} \quad U_x = \left(\frac{u_x' + v}{1 + \frac{v u_x'}{c^2}} \right)$$

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Such double transformations
should be done by matrix method and
fortunately we have MINKOWSKY MATRIX
which we can initial write as equivalent
of Lorentz transform

Aim is to
find u_3 s.t.

$$\alpha_3 = \frac{1}{\sqrt{1 - \frac{v_3^2}{c^2}}}$$

$$\beta_3 = \left(\frac{v_3}{c} \right)$$

$$\begin{bmatrix} \alpha & 0 & 0 & i\alpha\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha\beta & 0 & 0 & \alpha \end{bmatrix} = \begin{bmatrix} \alpha_2 & 0 & 0 & i\alpha_2\beta_2 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha_2\beta_2 & 0 & 0 & \alpha_2 \end{bmatrix} \begin{bmatrix} \alpha_1 & 0 & 0 & i\alpha_1\beta_1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha_1\beta_1 & 0 & 0 & \alpha_1 \end{bmatrix}$$

$$\text{Now } \alpha_1 = \frac{1}{\sqrt{1 - \frac{v_1^2}{c^2}}} ; \alpha_2 = \frac{1}{\sqrt{1 - \frac{v_2^2}{c^2}}} ; \beta_1 = \frac{v_1}{c} ; \beta_2 = \frac{v_2}{c}$$

$$= \begin{bmatrix} \alpha_1\alpha_2 + i\alpha_1\alpha_2\beta_1\beta_2 & 0 & 0 & i\alpha_1\alpha_2(\beta_1 + \beta_2) \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha_1\alpha_2(\beta_1 + \beta_2) & 0 & 0 & \alpha_1\alpha_2 + i\alpha_1\alpha_2\beta_1\beta_2 \end{bmatrix}$$

$$\text{Now } \alpha_1\alpha_2 [1 + \beta_1\beta_2]$$

$$= \frac{1}{\sqrt{1 - \frac{v_1^2}{c^2}} \sqrt{1 - \frac{v_2^2}{c^2}}} \left[1 + \frac{v_1 v_2}{c^2} \right]$$

$$= \frac{c^2 + v_1 v_2}{\sqrt{(c^2 - v_1^2)(c^2 - v_2^2)}} = \frac{1}{\sqrt{\frac{c^4 + v_1^2 v_2^2 - (v_1^2 + v_2^2)c^2}{(c^2 + v_1 v_2)^2}}}$$

$$= \frac{1}{\sqrt{\frac{(c^2 + v_1 v_2)^2 - 2c^2 v_1 v_2 - v_1^2 c^2 - v_2^2 c^2}{(c^2 + v_1 v_2)^2}}}$$

$$= \frac{1}{\sqrt{1 - \frac{(v_1 + v_2)^2 c^2}{c^4 \left(1 + \frac{v_1 v_2}{c^2}\right)^2}}} = \frac{1}{\sqrt{1 - \left(\frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}}\right)^2 \frac{1}{c^2}}}$$

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Also,

$$\begin{aligned} & \alpha, \alpha_2 [\beta_1 + \beta_2] \\ &= \frac{\alpha}{1 + \beta_1 \beta_2} [\beta_1 + \beta_2] \\ &= \alpha \cdot \frac{\left(\frac{v_1}{c} + \frac{v_2}{c}\right)}{\left(1 + \frac{v_1 v_2}{c^2}\right)} = \alpha \frac{\left(\frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}}\right)}{c} \end{aligned}$$

$$\text{let } \frac{\left(\frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}}\right)}{c} = \beta$$

$\Rightarrow$ Minkowsky Matrix, is

$$\begin{bmatrix} \alpha & 0 & 0 & i\alpha\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha\beta & 0 & 0 & \alpha \end{bmatrix}$$

$$\text{where } \alpha = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\text{where } v = \left(\frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}}\right)$$

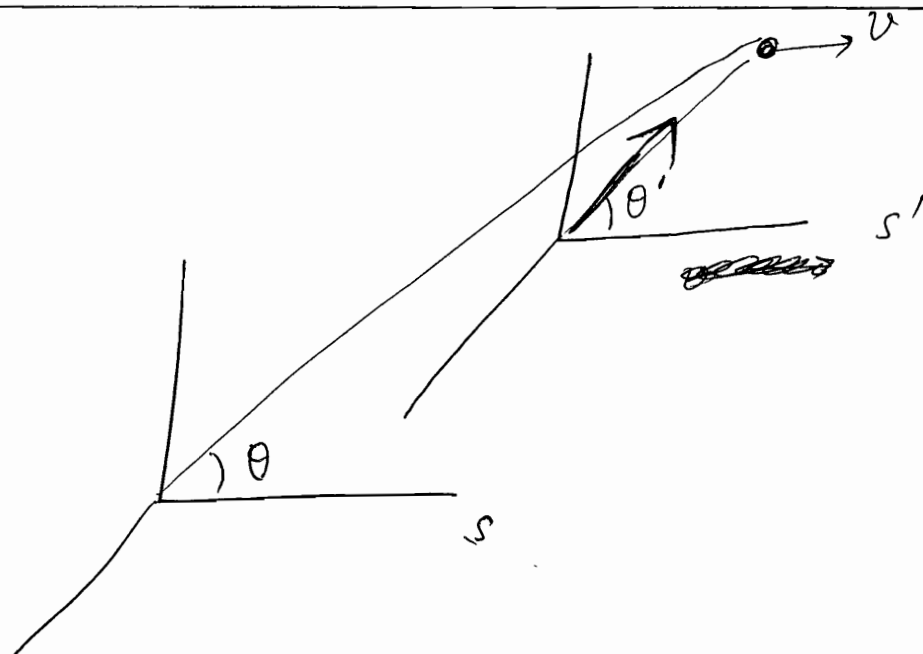
$$\beta = \left(\frac{v}{c}\right)$$

Hence 2 transformation equal to 1 transformation
with relativistic composition of velocities.

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Q 32)

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To prove :
$$\tan \theta = \frac{\sin \theta' \sqrt{1 - \frac{v^2}{c^2}}}{\cos \theta' + \left(\frac{v}{c}\right)}$$

For the set up as shown above,

$$u_{x'} = + c \cos \theta' \quad \text{--- (1)}$$

$$u_{y'} = + c \sin \theta' \quad \text{--- (2)}$$

Now from inverse Lorentz transform,

$$x = \frac{x' + vt'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$t = \frac{t' + \frac{v}{c^2} x'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y = y'$$

$$z = z'$$

$$\Rightarrow u_x = \frac{dx}{dt} = \frac{dx' + v dt'}{dt' + \frac{v}{c^2} dx'} = \left(\frac{u_{x'} + v}{1 + \frac{v u_{x'}}{c^2}} \right) \quad \text{--- (3)}$$

$$u_y = \frac{dy}{dt} = \frac{dy' \sqrt{1 - \frac{v^2}{c^2}}}{dt' + \frac{v}{c^2} dx'} = \frac{u_{y'}}{1 + \frac{v u_{x'}}{c^2}} \sqrt{1 - \frac{v^2}{c^2}} \quad \text{--- (4)}$$

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Using ① and ② in ③ and ④

$$u_x = \left(\frac{+C \cos \theta' + v}{1 + \frac{v}{c} \cos \theta} \right) = +C \cos \theta$$

$$u_y = \frac{+C \sin \theta'}{1 + \frac{v}{c} \cos \theta} \sqrt{1 - \frac{v^2}{c^2}} = +C \sin \theta$$

$$\Rightarrow \tan \theta = \frac{\sin \theta'}{\cos \theta' + \left(\frac{v}{c}\right)} \sqrt{1 - \frac{v^2}{c^2}}$$

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Q33)

$$v = 0.6c$$

~~For~~

For body moving with
 $v = 0.6c$, mass m

is

$$m = \frac{m_0}{\sqrt{1 - (0.6)^2}} = \frac{m_0}{0.8} = \left(\frac{5m_0}{4} \right)$$

Now $v_x = 0$
 $v_y = 0.6c$
 $v_z = 0$

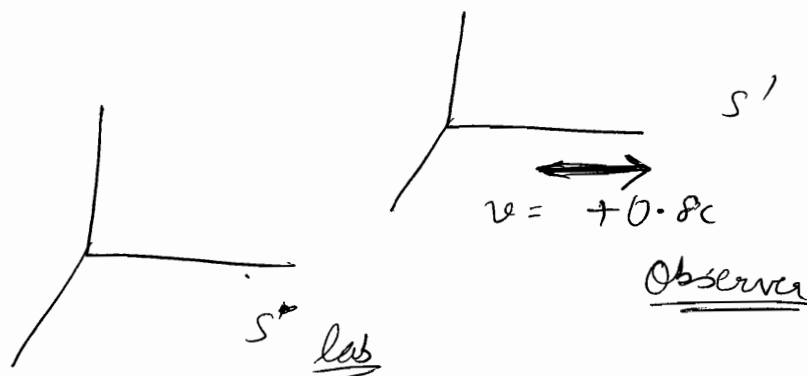
$$E = mc^2 = \frac{5m_0}{4} c^2$$

$$p_\mu = \left[0, \frac{3}{4} m_0 c, 0, i \frac{5m_0}{4} c \right]$$

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Now the Momentum 4 vector obeys
Minkowsky Transformation,
For frame ~~of~~ Observer ~~by~~ S'

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$$p_\mu' = \alpha_{\mu\nu} p_\nu$$

$$\alpha_{\mu\nu} = \begin{bmatrix} \alpha & 0 & 0 & i\alpha\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha\beta & 0 & 0 & \alpha \end{bmatrix}$$

$$\Rightarrow \text{Now } \alpha = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{1}{\sqrt{1 - (0.8)^2}} = \frac{1}{0.6} = \left(\frac{5}{3}\right)$$

$$\beta = \frac{v}{c} = 0.8 = \left(\frac{4}{5}\right)$$

$$\Rightarrow \begin{bmatrix} p_x' \\ p_y' \\ p_z' \\ i\frac{E'}{c} \end{bmatrix} = \begin{bmatrix} \alpha & 0 & 0 & i\alpha\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\alpha\beta & 0 & 0 & \alpha \end{bmatrix} \begin{bmatrix} 0 \\ \frac{3}{4}m_0c \\ 0 \\ i\frac{5}{4}m_0c \end{bmatrix}$$

$$= \begin{bmatrix} -\frac{5}{4}\alpha\beta m_0c \\ \frac{3}{4}m_0c \\ 0 \\ i\frac{5}{4}\alpha m_0c \end{bmatrix} = \begin{bmatrix} -\frac{5}{3}m_0c \\ \frac{3}{4}m_0c \\ 0 \\ i\frac{25}{12}m_0c \end{bmatrix}$$

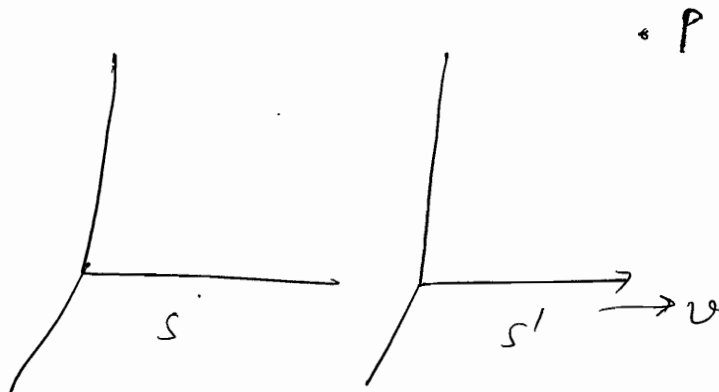
Q34) ① All Basic laws of Physics take the same physical form in all inertial frames.

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② velocity of light is constant in all inertial frames.

Let us consider two frames S and S' , latter moving at velocity v w.r.t. former in positive x -direction. Let the two frame coincide at $x=0$.



An event occurs at point P which is noticed by S and S' at t and t' since velocity of light is constant for both frames,

$$x^2 + y^2 + z^2 - c^2 t^2 = x'^2 + y'^2 + z'^2 - c^2 t'^2 \quad \text{--- ①}$$

$$\text{Let } x' = \alpha(x - vt)$$

$$\text{and } t' = \beta t + \gamma x$$

Putting in ①, we get

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$$\begin{aligned} & \alpha^2 x^2 + \alpha^2 v^2 t^2 - 2\alpha^2 v x t + y'^2 + z'^2 \\ & - c^2 \beta^2 t^2 - c^2 \gamma^2 x^2 - 2c^2 \beta \gamma x t \\ & = x^2 + y^2 + z^2 - c^2 t^2 \end{aligned}$$

Comparing coefficients,

$$\alpha^2 - c^2 = 1 \quad \text{--- (2)}$$

$$\alpha^2 v^2 - c^2 \beta^2 = -c^2 \quad - (3)$$

$$-2\alpha^2 v - 2C^2 \beta r = 0 \quad - (4)$$

$$\alpha^2 = 1 + c^2 \cancel{\beta^2}$$

~~XXXXXXXXXXXXXXXXXXXX~~

$$v^2 + v^2 c^2 \gamma^2 + c^2 = c^2 \beta^2$$

$$\Rightarrow \beta^2 = \left(1 + v^2 \gamma^2 + \frac{v^2}{c^2}\right)$$

$$\alpha^4 v^2 = c^4 \beta^2 r^2$$

[illegible]

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$$\alpha = \sqrt{1 + \frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2}\right)}$$

$$= \sqrt{1 + \frac{v^2}{c^2 - v^2}} = \sqrt{\frac{c^2}{c^2 - v^2}}$$

$$\alpha = \sqrt{\frac{1}{1 - \left(\frac{v^2}{c^2}\right)}}$$

$$\beta = \sqrt{1 + \frac{v^4}{c^4} \left(1 - \frac{v^2}{c^2}\right)} + \frac{v^2}{c^2}$$

$$= \sqrt{1 + \frac{v^4}{c^4 - v^2 c^2}} + \frac{v^2}{c^2}$$

$$= \sqrt{\frac{c^4 - \cancel{v^2 c^2} + \cancel{v^4} + \cancel{c^2 v^2} - \cancel{v^4}}{c^4 - v^2 c^2}}$$

$$\beta = \sqrt{\frac{1}{1 - \frac{v^2}{c^2}}}$$

Hence, $x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v}{c^2} x}{\sqrt{1 - \frac{v^2}{c^2}}}$$

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As there is no preferential frame, we
can have S travelling with velocity
(-v) w.r.t. S'. Writing Lorentz
inverse transform

$$x = \frac{x' + vt'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y = y'$$

$$z = z'$$

$$t = \frac{t' + \frac{v}{c^2} x'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Now let a clock be kept at x'

let $t'_1 - t'_2 = T_0$ (Proper time as the
clock is stationary)

From inverse transforms,

$$t_1 - t_2 = \frac{\frac{v}{c^2} (0) + T_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow \Delta T_{\text{obs}} = \frac{\Delta T_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

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$$\text{Time} = \frac{5 \mu s}{\sqrt{1 - (0.8)^2}} = \frac{5}{0.6} \mu s = \underline{\underline{8.33 \mu s}}$$

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Q36)

Let a rod be moving with a velocity v w.r.t. frame S along the $+x$ direction. Also let S' be frame attached to the rod. Let P be fixed mark in S frame & A, B be two end points of rod. Consider the two events

- (I) Front End A coincides with P
- (II) Back End B coincides with P

In frame S'

length of rod = l_0 (measured when it is at rest)

velocity of P is $-v$ in S' frame

$$\Rightarrow \Delta t' = \left(\frac{l_0}{v} \right)$$

As observed in frame S

$$\Delta t = \frac{L}{v}$$

$$\Delta t = \frac{l_0}{v}$$

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Now Δt is proper time as for both events occur at same position P

$$\Rightarrow \Delta t' = \frac{\Delta t}{\sqrt{1 - \frac{v^2}{c^2}}} \quad (\Delta x = 0)$$

$$\Rightarrow \frac{L_0}{v} = \frac{L_{obs}}{v \sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow L_{obs} = L_0 \sqrt{1 - \frac{v^2}{c^2}}$$

Hence length contraction.....

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Q37)

Q38)
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Q - 27

Let S and S' be two frames.

dx be length measured in ~~rest~~ stationary frame and dx' be length measured in the moving frame. i.e. when the body is at rest w.r.t. frames dx' : Proper length

$$dx = \sqrt{1 - \frac{v^2}{c^2}} \cdot dx' \quad (\text{length contraction})$$

Let time dt' be measured in moving frame with x' const. where clock is kept at rest w.r.t. frame $\Rightarrow dt'$: proper time

$$dt = \frac{dt'}{\sqrt{1 - \frac{v^2}{c^2}}} \quad (\text{time dilation})$$

$$\Rightarrow dx \, dy \, dz \, dt = dx' \, dy' \, dz' \, dt'$$

Q39)

Q 32

Q40)

Q18)

No it does not as it was inferred that Earths drags ether along with it and relative velocity = 0

Einstein's hypothesis finally settled the debate

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OPTICS (Selected Questions)

Tut Sheet 7A Damped & Forced Oscillations

Q 1)

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For the weak damping system, or the underdamped system, the response is given by:

$$x = A e^{-\lambda t} \sin(\omega_d t - \theta)$$

$$\text{where } \omega_d = \sqrt{\omega_0^2 - c^2} \approx \omega_0 \quad (\text{given})$$

$$\text{We are given } \omega_0 \tau \gg 1$$

$$\text{i.e. } \omega_0 \cdot \frac{1}{2c} \gg 1 \quad (\tau: \text{relaxation time})$$

$$\Rightarrow \omega_0 \gg 2c$$

$$\Rightarrow \sqrt{\omega_0^2 - c^2} \approx \omega_0$$

$$\Rightarrow x = A e^{-\lambda t} \sin(\omega_0 t - \theta) \quad \checkmark$$

Rate of energy dissipation or Power of dissipation, is given as:

$$\left(\frac{dE}{dt}\right) = -b \cdot v \cdot v = -2c \dot{x}^2 \quad \checkmark$$

$$= -2c A^2 e^{-2\lambda t} \sin^2(\omega_0 t - \theta)$$

$$= -2c A^2 \sin^2(\omega_0 t - \theta) e^{-2\lambda t} \quad (-2c)$$

$$= [4c^2 A^2 \sin^2(\omega_0 t - \theta) \cos^2(\omega_0 t - \theta)] e^{-2\lambda t}$$

Q 2)

Lecture 2

Q 3)

Lecture 3

Q 4)

Lecture 3

Q 5)

Lecture 3

Q 6)

For a damped oscillator, the equation

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of motion is

$$\ddot{x} - b\dot{x} - \cancel{4k}x = 0$$

Solution for the underdamped case is

$$x = Ae^{-\cancel{1}ct} \sin(\omega dt + \theta)$$

We are given $\omega d = 2\pi \times 3000 \text{ s}^{-1}$

$$\Rightarrow \text{Time for 3000 oscillations} = 3000 T.$$

$$= 3000 \cdot \frac{2\pi}{2\pi \times 3000} = 10 \text{ s}$$

Let initial Amplitude be A

$$\Rightarrow \text{Amplitude after 10 seconds} = Ae^{-10c}$$

We are given,

$$Ae^{-10c} = \frac{A}{10} \Rightarrow \boxed{c = 0.23 \text{ s}^{-1}}$$

$$\text{Damping constant} = b = 2c = \boxed{0.46 \text{ s}^{-1}}$$

Energy variation fall off twice as
amplitude

$$\text{i.e. } E = E_0 e^{-2ct} \sin(\omega dt - \theta)$$

$\therefore$ Time for energy to reduce to $1/10^{\text{th}}$
of the value is given as

$$\Rightarrow \frac{E_0}{10} = E_0 e^{-2ct}$$

$$\Rightarrow \boxed{t = 5 \text{ s}}$$

Q.7

lecture 2

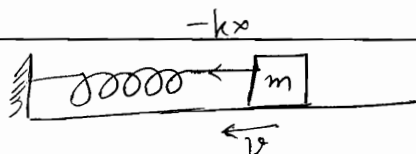
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damping
constant = ~~2c~~
= b

$$b = 2mc ??$$

Q8)

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For the spring system, we have,

$$-kx - b\dot{x} = m\ddot{x}$$

$$\Rightarrow \ddot{x} + 2c\dot{x} + \omega_0^2 x = 0 \quad \text{where } \frac{k}{m} = \omega_0^2$$

<derivation>

For a light damping, we have. $\frac{b}{m} = 2c$

Under damped solution

$$\text{i.e. } x = A e^{-ct} \sin(\omega_d t - \theta)$$

$$\omega_d = \sqrt{\omega_0^2 - c^2}$$

$$\begin{aligned} \text{Frequency of Oscillation} &= \frac{1}{2\pi} \sqrt{2\omega_0^2 - c^2} \\ &= \frac{1}{2\pi} \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}} \end{aligned}$$

In the absence of damping i.e. $b=0$,
we have

$$f = \frac{1}{2\pi} \sqrt{\frac{k}{m}}$$

We are given,

$$\frac{1}{2\pi} \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}} = \frac{1}{\sqrt{2}} \cdot \frac{1}{2\pi} \sqrt{\frac{k}{m}}$$

$$\Rightarrow \frac{2k}{m} - \frac{b^2}{2m^2} = \frac{k}{m}$$

$$\begin{aligned} \Rightarrow \frac{k}{m} &= \frac{b^2}{2m^2} \Rightarrow b = \sqrt{2km} \\ &= \sqrt{2 \times 10 \times 0.1} \\ &= \sqrt{2} \text{ Ns/m} \end{aligned}$$

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$\frac{N \cdot s}{m}$

Q9)

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Q10)

Lecture 2

The equation for oscillation under external harmonic force is :

$$\ddot{x} + 2c\dot{x} + \omega_0^2 x = F_0 \sin(pt) \quad \text{where } F_0 = \frac{F}{m}$$

$$\text{Let } x = A \sin(pt - \theta)$$

$$\text{where } \omega_0^2 = \frac{k}{m} \\ 2c = 2b = \beta$$

$$\Rightarrow -Ap^2 \sin(pt - \theta) + 2cAp \cos(pt - \theta) + \omega_0^2 A \sin(pt - \theta) = F_0 \sin(pt - \theta + 0)$$

$$\Rightarrow (\omega_0^2 A - Ap^2) \sin(pt - \theta) + 2cAp \cos(pt - \theta) = F_0 \cos \theta \sin(pt - \theta) + F_0 \sin \theta \cos(pt - \theta)$$

Comparing coefficients, we have

$$\omega_0^2 A - Ap^2 = F_0 \cos \theta$$

$$2cAp = F_0 \sin \theta$$

$$\Rightarrow \tan \theta = \left(\frac{2cp}{\omega_0^2 - p^2} \right) \\ A = \frac{F_0}{\sqrt{(\omega_0^2 - p^2)^2 + (2cp)^2}}$$

For maximum Amplitude,

$$\frac{dA}{dp} = 0 \quad \& \quad \frac{d^2A}{dp^2} < 0$$

$$\Rightarrow \frac{-F_0}{[(\omega_0^2 - p^2)^2 + (2cp)^2]^{3/2}} \cdot [2(\omega_0^2 - p^2)(-2p) + 4(2cp)c] = 0$$

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$$\Rightarrow 8c^2p = 4p(\omega_0^2 - p^2)$$

$\Rightarrow$ Cancelling the trivial solution $p=0$,

$$2c^2 - \omega_0^2 = -p^2$$

$$\Rightarrow \boxed{p = \sqrt{\omega_0^2 - 2c^2}}$$

$$\text{At } p = \sqrt{\omega_0^2 - 2c^2}$$

$$A = \frac{F}{m \sqrt{(2c^2)^2 + (2c)^2 (\omega_0^2 - 2c^2)}}$$

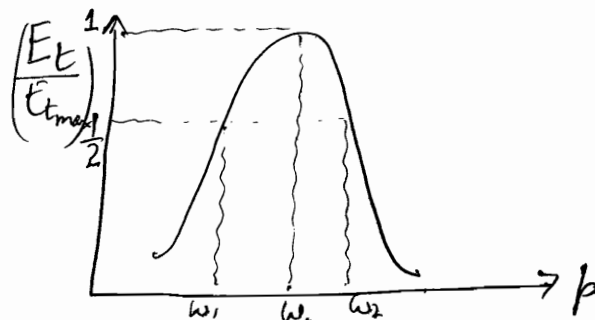
$$= \frac{F}{m \sqrt{4c^4 + 4c^2 \omega_0^2 - 8c^4}}$$

$$= \frac{F}{m \sqrt{4c^2 \omega_0^2 - 4c^4}}$$

$$= \frac{F}{2cm \sqrt{\omega_0^2 - c^2}}$$

$$\boxed{A = \frac{f}{2c \sqrt{\omega_0^2 - c^2}}}$$

The energy transferred by the external system varies as



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The sharpness of resonance is the
measure of the sharpness of this curve.
Let the resonance occur at ω_0 and
let half power max be transferred at
 ω_1 and ω_2 .

Mathematically, sharpness is defined
in terms of Quality factor Q ,

$$Q = \left(\frac{\omega_0}{\omega_2 - \omega_1} \right)$$

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Test Sheet : 7B → Beats, Stationary Wave

Q1)

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Mathematically, group velocity v_g is given
as,

$$v_g = v_p - \lambda \left(\frac{dv_p}{d\lambda} \right) \quad \text{--- ①}$$

$$\text{Given } v_p = \left[\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right]^{1/2}$$

$$\Rightarrow \frac{dv_p}{d\lambda} = \frac{1}{2} \left[\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right]^{-1/2} \left[\frac{-2\pi T}{\lambda_p^2} + \frac{g}{2\pi} \right]$$

From ①

$$\Rightarrow v_g = \left[\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right]^{1/2} + \frac{\left[\frac{\pi T}{\lambda_p} - \frac{g\lambda}{4\pi} \right]}{\left(\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right)^{1/2}}$$

$$= \frac{\left[\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} + \frac{\pi T}{\lambda_p} - \frac{g\lambda}{4\pi} \right]}{\left(\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right)^{1/2}}$$

$$= \frac{\left[\frac{3\pi T}{\lambda_p} + \frac{g\lambda}{4\pi} \right]}{\left(\frac{2\pi T}{\lambda_p} + \frac{g\lambda}{2\pi} \right)^{1/2}}$$

$$= \frac{2\pi T}{\lambda_p v_p} + \frac{1}{2} v_p$$

$$v_g = \frac{v_p}{2} + \frac{2\pi T}{\lambda_p v_p}$$

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For phase velocity to be minimum,

$$\frac{dv_p}{d\lambda} = 0$$

$$\Rightarrow \frac{2\pi T}{\rho \lambda^2} = \frac{g}{2\pi} \Rightarrow$$

$$\lambda = 2\pi \sqrt{\frac{T}{\rho g}}$$

$$\Delta \left(\frac{d^2 v_p}{d\lambda^2} \right) > 0$$

for $\lambda = 2\pi \sqrt{\frac{T}{\rho g}}$, double derivative
is ~~negative~~ ^{positive}, hence a minima.

Q2) Lecture 1

Q3) Given λ n

5090	1.647
5340	1.640
5890	1.630

$$\Rightarrow \frac{\Delta n}{\Delta \lambda} = \frac{-2.8 \times 10^{-5} - 1.8 \times 10^{-5}}{2} \\ = -2.3 \times 10^{-5}$$

$$\text{Phase velocity} = v = \frac{c}{n} \\ = \frac{3 \times 10^8}{1.640} \text{ m/s} \\ = 1.82 \times 10^8 \text{ m/s} \checkmark$$

Group velocity is given as

$$v_g = v_p - \lambda \left(\frac{dv_p}{d\lambda} \right)$$

$$\text{Now } v_p = \frac{c}{n}$$

$$\Rightarrow \frac{dv_p}{d\lambda} = -\frac{c}{n^2}$$

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$$\Rightarrow v_g = v_p - \lambda \frac{dv_p}{dn} \cdot \frac{dn}{d\lambda}$$

$$= v_p + \lambda \cdot \frac{c}{n^2} \left(\frac{dn}{d\lambda} \right)$$

$$\approx v_p + \lambda \cdot \frac{c}{n^2} \left(\frac{\Delta n}{\Delta \lambda} \right)$$

$$= 1.82 \times 10^8 + \frac{5340 \times 10^{10} \times 3 \times 10^8}{(1.64)^2} \cdot 2.3 \times 10^{-5}$$

$$= (1.82 + 0.137) \times 10^8$$

$$= \underline{\underline{1.683 \times 10^8 \text{ m/s}}} \checkmark$$

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Q4) Lecture 1

Q5) For a string the velocity of a travelling wave v is given as

$$v = \sqrt{\frac{T}{\lambda}}$$

T : Tension

λ : mass per unit length

$$= \sqrt{\frac{25}{0.25}} = \underline{\underline{10 \text{ m/s}}}$$

velocity
on
string

$$v_{\text{wave}} = 5 \text{ Hz}$$

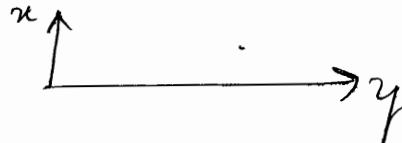
$$\lambda_{\text{wave}} = 0.01 \text{ m}$$

The equation of wave is given as

$$x = A \sin(\omega t - ky)$$

$$k = \left(\frac{2\pi}{\lambda} \right) \quad \omega = (2\pi \nu)$$

$$\lambda = \frac{v}{\nu} = \frac{10}{5} = \underline{\underline{2 \text{ m}}}$$



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Q6)

$$\Rightarrow \text{wave equation : } x = 0.01 \sin(10\pi t - \pi y) \text{ m}$$

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$$v_p = \sqrt{\frac{g}{k}} = \sqrt{\frac{g\lambda}{2\pi}}$$

$$v_g = v_p - \lambda \left(\frac{dv_p}{d\lambda} \right)$$

$$= \sqrt{\frac{g\lambda}{2\pi}} - \lambda \cdot \sqrt{\frac{g}{2\pi}} \cdot \frac{1}{2\sqrt{\lambda}}$$

$$= \frac{1}{2} \cdot \sqrt{\frac{g\lambda}{2\pi}} = \frac{1}{2} \sqrt{\frac{g}{k}} = \left(\frac{v_p}{2} \right)$$

Q7)

$$\psi = 10 \cos(5x + 25t)$$

Comparing with $\psi = A \cos(kx + \omega t)$, we have

$$\lambda = \frac{2\pi}{k} = \frac{2\pi}{5} = \underline{\underline{1.256 \text{ m}}}$$

$$v = \lambda \nu = \frac{\omega}{k} = \frac{25}{5} = \underline{\underline{5 \text{ m/s}}}$$

$$\nu = \frac{\omega}{2\pi} = \frac{25}{2\pi} \approx \underline{\underline{4 \text{ Hz}}}$$

direction of propagation : $(-x)$ axis

$$\psi_1 = 10 \cos(5x + 25t)$$

$$\psi_2 = 20 \cos(5x + 25t + \pi/3)$$

$$\psi = \psi_1 + \psi_2$$

$$= \cancel{10 \cos(5x + 25t)} + 10 \cos(5x + 25t)$$

$$+ 20 \cos(5x + 25t) \cdot \frac{1}{2}$$

$$- 20 \sin(5x + 25t) \cdot \frac{\sqrt{3}}{2}$$

$$= 30 \cos(5x + 25t) - 10\sqrt{3} \sin(5x + 25t)$$

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$$= -\sqrt{700} \sin \left[5x + 25t - \tan^{-1} \left(\frac{2}{\sqrt{3}} \right) \right]$$

$$= \sqrt{700} \cos \left[5x + 25t + \left[\frac{\pi}{2} - \tan^{-1} \left(\frac{2}{\sqrt{3}} \right) \right] \right]$$

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$$\Rightarrow \theta = \frac{\pi}{2} - \tan^{-1} \left(\frac{2}{\sqrt{3}} \right)$$

$$= \tan^{-1} \left(\frac{\sqrt{3}}{2} \right)$$

$$\approx \underline{41^\circ}$$

Remember
to convert
the form
to initial
type....

(Q8)

(Q9)

Q1

$$k = \frac{2\pi}{1} = 2\pi \text{ m}^{-1}$$

$$v = \frac{\omega}{k} \Rightarrow \omega = vk = 2\pi \times 3 = \underline{6\pi \text{ s}^{-1}}$$



For wave 1

$$y = A \sin (6\pi t - 2\pi x)$$

For wave 2

$$y = A \sin (6\pi t + 2\pi x)$$

$$\Rightarrow y_{\text{total}} = y_1 + y_2$$

$$= \underline{2A \sin (6\pi t) \cos (2\pi x)}$$

Nodes

$$y = 0 \Rightarrow 2\pi x = (2n+1)\pi/2$$

$$\Rightarrow \boxed{x = \frac{(2n+1)}{4} \text{ m}} = \frac{1}{4} \text{ m}, \frac{3}{4} \text{ m}, \dots$$

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Antinodes

$$y = y_{\max}$$

$$\Rightarrow 2\pi x = n\pi$$

$$\Rightarrow x = \frac{n\pi}{2} = 0, \frac{1}{2}\pi, \frac{1}{2}\pi, \frac{3}{2}\pi, \dots$$

if strings are fixed at both ends, then at $x=0$,
it will be node ALWAYS. hence choose the
wave wisely i.e. amplitude of resultant wave
should have $\sin kx$ term

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★ Very
important
to take
right type of
wave at
starting in
standing waves

let the two waves be

$$y_1 = A \sin(kx - \omega t)$$

$$y_2 = A \sin(kx + \omega t)$$

$$k = 2\pi$$

$$\omega = 6\pi$$

$$\Rightarrow y = A \sin kx \cos \omega t$$

$$\Rightarrow \text{At } x=0, y=0 : \text{satisfied}$$

Node

$$kx = n\pi$$

$$\Rightarrow x = \frac{n\pi}{k} = \frac{(n)}{2} \pi$$

Antinode

$$kx = (2n+1)\frac{\pi}{2}$$

$$\Rightarrow x = \frac{(2n+1)}{4} \pi$$

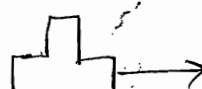
Q10)

No underroot
in
sound
doppler !!

For doppler effect in sound,



R



$$v = 10$$

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$$\begin{aligned} \nu_{\text{direct}} &= \nu_0 \frac{v - 0}{v + v_s} = 900 \frac{330}{340} \\ &= \underline{\underline{873 \text{ Hz}}} \end{aligned}$$

$$\begin{aligned} \nu_{\text{indirect}} &= \nu_{\text{wall}} \\ &= \nu_0 \frac{v - 0}{v - v_s} = 900 \frac{330}{320} \\ &= \underline{\underline{928.9 \text{ Hz}}} \end{aligned}$$

$$\begin{aligned} \text{No. of beats} &= \nu_1 - \nu_2 \\ &= 55 \end{aligned}$$

But person's ear can't decipher more than 10 beats - Hence No beats heard.

Q11)



Let the wave be described as
 $y_1 = A \sin(kx + \omega t)$ (travelling in $-x$ direction)

Since the point O is fixed, it will be reflected off; let the reflected wave be:

$$y_2 = A \sin(kx - \omega t)$$

$\therefore$ the total standing wave will be superposition of y_1 and y_2

$$\Rightarrow y_{\text{total}} = y_1 + y_2 = 2A \sin kx \cos \omega t$$

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Now for nodes

$$y(x) = 0$$

$$\Rightarrow \sin kx = 0$$

$$\Rightarrow kx = n\pi$$

$$\Rightarrow x = \frac{n\pi}{k} = \frac{n\pi}{\frac{2\pi}{\lambda}} \cdot \lambda = \left(\frac{n\lambda}{2}\right)$$

$\therefore$ Nodes at $\left[0, \frac{\lambda}{2}, \lambda, \frac{3\lambda}{2}, \dots\right]$

At antinodes,

$$\sin kx = 1$$

$$\Rightarrow y = 2A \cos \omega t$$

Now the kinetic energy for a small
length dl around the antinode is
given by,

$$k.E. = \frac{1}{2} m \dot{y}^2$$

$$= \frac{1}{2} \cdot dl \cdot \rho \cdot (2A\omega)^2 \cdot \cos^2 \omega t$$

$$= dl \cdot 2A^2 \omega^2 \rho \cos^2 \omega t$$

$$\Rightarrow k.E \text{ per unit length} = \underline{2A^2 \omega^2 \rho \cos^2 \omega t}$$

Q12)

Lecture 1

Q13)

The equation for a standing wave
fixed at both the ends is given as

$$y = 2A \sin kx \cos \omega t$$

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$$\begin{aligned} \text{At } x=0 \quad y=0 \\ \text{At } x=L \quad y=0 \Rightarrow kL = n\pi \\ \Rightarrow k = \left(\frac{n\pi}{L}\right) \end{aligned}$$

$$\Rightarrow y = 2A \sin\left(\frac{n\pi x}{L}\right) \cos \omega t$$

Let a consider a different portion dx
around the point x

$$\Rightarrow y|_x = 2A \sin\left(\frac{n\pi x}{L}\right) \cos \omega t$$

$$\dot{y}|_x = -2A\omega \sin\left(\frac{n\pi x}{L}\right) \cdot \sin \omega t$$

The energy density is given as

$$\frac{E}{dl} = \frac{1}{2} \frac{dm}{dl} v^2$$

$$= \frac{1}{2} \rho \frac{dl}{dl} \cdot 4A^2 \omega^2 \sin^2\left(\frac{n\pi x}{L}\right) \sin^2 \omega t$$

$$\begin{aligned} \Rightarrow \text{Energy per unit length} &= \text{Energy density } (E_x) \\ &= 2\rho A^2 \omega^2 \sin^2\left(\frac{n\pi x}{L}\right) \sin^2 \omega t \end{aligned}$$

$$\text{In a time period } \langle \sin^2 \omega t \rangle = \frac{1}{2}$$

$$\Rightarrow \boxed{\langle E_x \rangle = \rho A^2 \omega^2 \sin^2\left(\frac{n\pi x}{L}\right)}$$

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Now E_x will be maximum where

$$\sin^2\left(\frac{n\pi x}{L}\right) = 1$$

i.e. $\left|\sin\left(\frac{n\pi x}{L}\right)\right| = 1$ i.e. antinodes

& minimum

at $\left|\sin\left(\frac{n\pi x}{L}\right)\right| = 0$ i.e. nodes.

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Q14)

Let a Consider a tube of cross sectional area A . Let PQ and RS be two transverse sections of the tube located at x and $x + \Delta x$ from fixed point O. Let the longitudinal displacement be denoted by $\xi(x)$.

$\therefore$ the displacement of the two planes be $\xi(x)$ and $\xi(x + \Delta x)$

Now the difference between the two planes is

final separation
between the two
planes

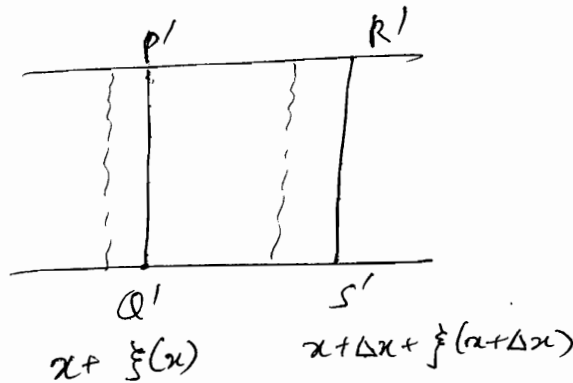
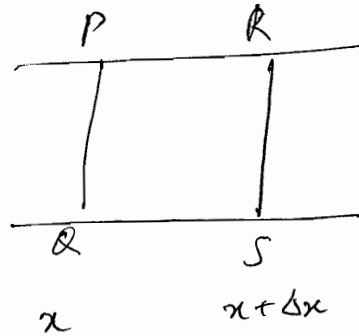
$$\begin{aligned} & [(x + \Delta x) + \xi(x + \Delta x)] - [x + \xi(x)] \\ &= \Delta x + \Delta x \left(\frac{d\xi}{dx} \right) \end{aligned}$$

$$\begin{aligned} \text{Change in volume} &= A \times \text{Change in length} \\ &= A \Delta x \left(\frac{d\xi}{dx} \right) \end{aligned}$$

$$\text{Original volume} = A \Delta x$$

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$$\Rightarrow \left(\frac{dV}{V} \right) = \text{strain} = \left(\frac{d\xi}{dx} \right)$$



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judge
साहब
के derivation
note की
जर्न !!

Now let the change the volume be dV
and the final volume be V .

Since mass remains same,

$$\Rightarrow M_1 = M_2$$

$$\Rightarrow \rho_0 [V - dV] = \rho V$$

$$\Rightarrow \rho = \rho_0 \left[\frac{V - dV}{V} \right] = \rho_0 \left[1 - \frac{dV}{V} \right]$$

$$\boxed{\rho = \rho_0 \left[1 - \frac{d\xi}{dx} \right]}$$

Or take
initial
volume as
 V and
binomially
we can
have some
result....

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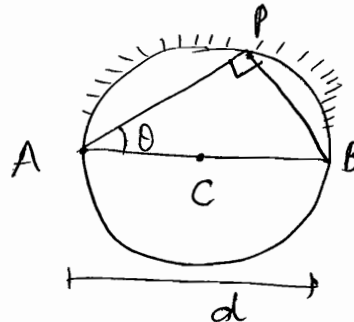
Tut Sheet 8 : Geometrical Optics.

Q1)

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Now we know for a point that joins
2 diametrically opposite points, an angle
of 90° is subtended.

Now the Path length
is $d \cos \theta + d \sin \theta$



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For extremely,

$$ds = 0$$

$$\Rightarrow -d \sin \theta d\theta + d \cos \theta d\theta = 0$$

$$\Rightarrow \sin \theta = \cos \theta \Rightarrow \boxed{\theta = 45^\circ}$$

$$\begin{aligned} \Rightarrow \overline{APB} &= 2d \cdot \cos 45^\circ + 2d \sin 45^\circ \\ &= \underline{\underline{2\sqrt{2}d}} = \underline{\underline{\sqrt{2}d}} \end{aligned}$$

$$\frac{d^2s}{d\theta^2} = -d \cos \theta \Big|_{\theta=45^\circ} \neq d \sin \theta \Big|_{\theta=45^\circ}$$

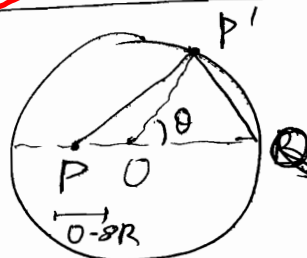
$$< 0$$

Hence the path is maximum

Q2)

To prove: $\cos\left(\frac{\theta}{2}\right) = \frac{3}{4}$

$$PP' = \sqrt{OP^2 + OP'^2 - 2OP \cdot OP' \cos(\pi - \theta)}$$



$$= \sqrt{(0.8R)^2 + R^2 + 2 \cdot 0.8R^2 \cos \theta}$$

$$= \sqrt{1.64 + 1.6 \cos \theta} R$$

$$P'O = \sqrt{R^2 + R^2 - 2R^2 \cos \theta}$$

$$= R \sqrt{2 - 2 \cos \theta} = R \cdot 2 \sin\left(\frac{\theta}{2}\right)$$

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$$\frac{ds}{d\theta} = 0 \quad \text{where } s = R\sqrt{1.64 + 1.6 \cos \theta} + 2R \sin\left(\frac{\theta}{2}\right)$$

$$\Rightarrow \frac{-1R}{2\sqrt{1.64 + 1.6 \cos \theta}} \cdot 1.6 \sin \theta + 2R \cos\left(\frac{\theta}{2}\right) \cdot \frac{1}{2} = 0$$

$$\Rightarrow \cos\left(\frac{\theta}{2}\right) = \frac{0.8 \sin \theta}{\sqrt{1.64 + 1.6 \cos \theta}}$$

$$\Rightarrow 1.64 + 1.6 \cos \theta = \left[1.6 \sin\left(\frac{\theta}{2}\right)\right]^2$$

$$\Rightarrow 1.64 + 1.6 \left[2 \cos \frac{\theta}{2} - 1\right] = 2.56 \left[1 - \cos \frac{\theta}{2}\right]$$

$$\Rightarrow \cos^2\left(\frac{\theta}{2}\right) [3.2 + 2.56] = 2.56 - 1.64$$

$$\Rightarrow \cos \frac{\theta}{2} = \frac{2.52}{5.76}$$

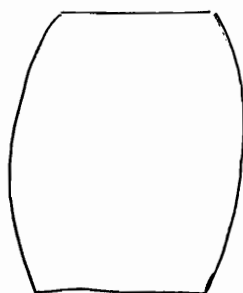
$$\cos\left(\frac{\theta}{2}\right) = 0.4375$$

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देना से करो
आ जायगा !!

Q3)

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$$t = 1 \text{ cm}$$

$$R_1 = +4 \text{ cm} \quad R_2 = -4 \text{ cm}$$

Let the system matrix be $\begin{bmatrix} b & -a \\ -d & c \end{bmatrix}$

It is composition of 3 operations:

$$\begin{bmatrix} 1 & -P_2 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ \frac{d}{n_2} & 1 \end{bmatrix} \begin{bmatrix} 1 & -P_1 \\ 0 & 1 \end{bmatrix}$$

$$P_1 = \frac{n_2 - n_1}{R_1} = \frac{1.5 - 1}{4} = \frac{0.5}{4} = \frac{1}{8}$$

$$P_2 = \frac{n_2 - n_1}{R_2} = \frac{1 - 1.5}{-4} = \frac{1}{8}$$

$$\Rightarrow \begin{bmatrix} b & -a \\ -d & c \end{bmatrix} = \begin{bmatrix} 1 & -\frac{1}{8} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ \frac{2}{3} & 1 \end{bmatrix} \begin{bmatrix} 1 & -\frac{1}{8} \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1 - \frac{1}{12} & -\frac{1}{8} \\ \frac{2}{3} & 1 \end{bmatrix} \begin{bmatrix} 1 & -\frac{1}{8} \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} \frac{11}{12} & -\frac{1}{8} \\ \frac{2}{3} & 1 \end{bmatrix} \begin{bmatrix} 1 & -\frac{1}{8} \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} \frac{11}{12} & -\left(\frac{11}{12} + 1\right) \cdot \frac{1}{8} \\ \frac{2}{3} & -\frac{1}{12} + 1 \end{bmatrix} = \begin{bmatrix} \frac{11}{12} & -\frac{23}{96} \\ \frac{2}{3} & \frac{11}{12} \end{bmatrix}$$

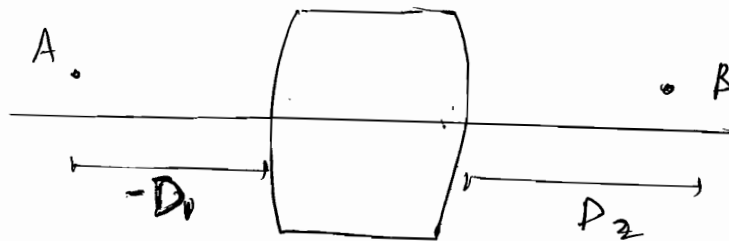
Determinant
should be
1, not 0

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Now Consider 2 points D_1 and D_2
distance away

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~~$$\begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ D_2 & 1 \end{bmatrix} \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -D_1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix}$$~~

~~$$\Rightarrow \begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} a & b \\ D_2 a + c & D_2 b + d \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -D_1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix}$$

$$= \begin{bmatrix} a - bD_1 & b \\ D_2 a + c - D_1 D_2 b - dD_1 & D_2 b + d \end{bmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix}$$~~

~~For $\frac{a}{D_1} + \frac{d}{D_2} = \frac{b - c}{D_1 D_2}$~~

~~$$M = D_2 b + d = \frac{D_2 a + c}{D_1} = 1$$~~

~~$$\Rightarrow D_2 = \frac{1 - d}{b}$$

$D_1 =$~~

$$\begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ D_2 & 1 \end{bmatrix} \begin{bmatrix} b & -a \\ -d & c \end{bmatrix} \begin{bmatrix} 1 & 0 \\ D_1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix}$$

$$= \begin{bmatrix} b & -a \\ D_2 b - d & c - D_2 a \end{bmatrix} \begin{bmatrix} 1 & 0 \\ D_1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix}$$

$$= \begin{bmatrix} b - aD_1 & -a \\ D_2 b - d + cD_1 & -aD_1 D_2 + c - D_2 a \end{bmatrix}$$

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$$\frac{b}{D_1} + \frac{c}{D_2} = \frac{d}{D_1 D_2} + a$$

$$M = C - D_2 a$$

$$= \cancel{b - a D_1}$$

Now focal length = $-\frac{1}{a} = \underline{\underline{\left(\frac{96}{23}\right)}}$

this
comes from
Det = 1

Nodal
~~and~~ points at

$$D_1 = \frac{b-1}{a} = \frac{\frac{11}{12} - 1}{\left(\frac{23}{96}\right)}$$

$$D_2 = \frac{c-1}{a}$$

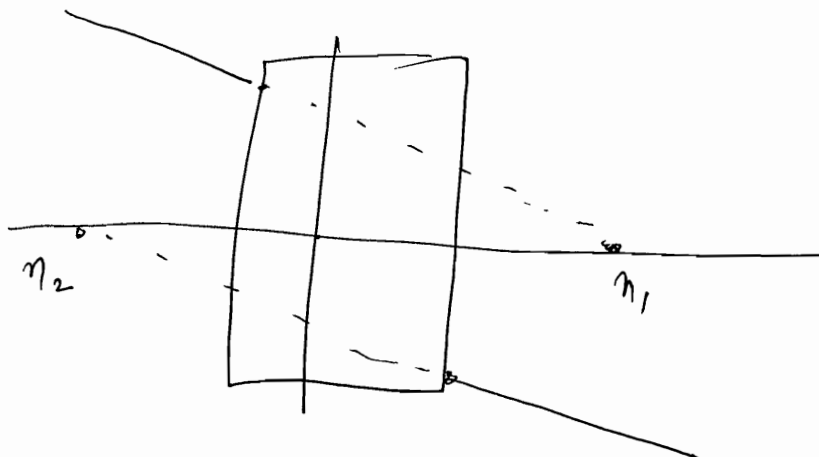
$$= \frac{\frac{11}{12} - 1}{\left(\frac{23}{96}\right)} = \underline{\underline{-\left(\frac{8}{23}\right)}}$$

$$= -\frac{1}{12} \times \frac{96}{23}$$

$$= \underline{\underline{-\left(\frac{8}{23}\right)}}$$

both negative? one

unit points? same ✓



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Q 4)

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$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2}$$

$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2}$$

Considering the matrix method

$$\begin{bmatrix} 1 & 0 \\ v & 1 \end{bmatrix} \begin{bmatrix} 1 & -\frac{1}{2} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ t & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

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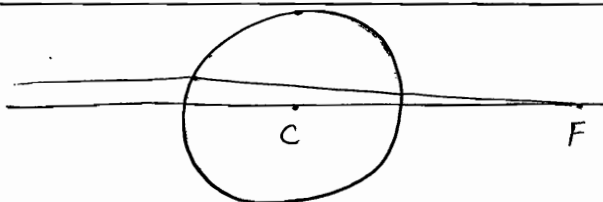
अरे भय्या focal lengths तो दे दो !!

Q5

कृपया इस स्थान में कुछ न लिखें।
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$$P_1 = \left(\frac{\mu-1}{R} \right) = P$$

$$P_2 = + \left(\frac{\mu-1}{R} \right) = P$$



Combined matrix

$$= \begin{bmatrix} 1 & -P \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ \frac{2R}{\mu} & 1 \end{bmatrix} \begin{bmatrix} 1 & -P \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1 & -\frac{2PR}{\mu} & -P \\ 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & -P \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1 & -\frac{2PR}{\mu} & -P + \frac{P^2 R}{\mu} \\ 0 & 1 & -\frac{PR}{\mu} + 1 \end{bmatrix}$$

Step 2
multiply
front !!

$$\begin{bmatrix} \lambda \\ y \end{bmatrix} = \begin{bmatrix} 1 - \frac{PR}{\mu} & -P + \frac{P^2 R}{\mu} \\ \frac{R}{\mu} & -\frac{PR}{\mu} + 1 \end{bmatrix} \begin{bmatrix} \lambda_i \\ y_i \end{bmatrix}$$

$$\frac{PR}{\mu} = 0.375$$

$$R = 20 \text{ cm} = 0.2 \text{ m}$$

$$\mu = 1.6$$

$$P = 3 \text{ m}^{-1}$$

$$\begin{bmatrix} 0.625 & -4.8 \\ 0.125 & 0.625 \end{bmatrix}$$

~~Calculating~~ Calculating nodal points

$$\begin{bmatrix} 1 & 0 \\ v & 1 \end{bmatrix} \begin{bmatrix} a & 1 \\ c & d \end{bmatrix} \begin{bmatrix} 1 & 0 \\ u & 1 \end{bmatrix} = \begin{bmatrix} a & b \\ av+c & bv+d \end{bmatrix} \begin{bmatrix} 1 & 0 \\ u & 1 \end{bmatrix} = \begin{bmatrix} a+bu & b \\ av+c+bu+ud & bv+d \end{bmatrix}$$

$$v = -\frac{a+d}{b}$$

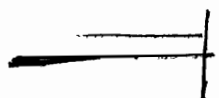
$$u = -\frac{a+d}{b}$$

$$v = -0.28$$

$$0.077$$

$$u = -0.28$$

$$0.077$$



$$\begin{bmatrix} \lambda \\ y_0 \end{bmatrix} = \begin{bmatrix} 1 & -4.875 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ y_0 \end{bmatrix}$$

$$\Rightarrow \lambda = -4.875 y_0$$

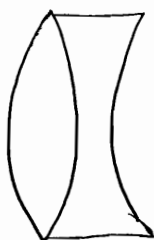
$$-\tan \theta = -4.875 y_0$$

$$f = \frac{1}{4.875} = 0.205$$

$$\Rightarrow \text{Position of Paraxial focus} = 0.282 \text{ m}$$

Q6)

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1: crown
2: flint

कृपया इस स्थान
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$$\frac{1}{f_1} + \frac{1}{f_2} = \frac{1}{20} \quad \text{--- (1)}$$

$$\text{Also } -\frac{f_1}{f_2} = \frac{w_1}{w_2} \approx \frac{(\mu_b - \mu_r)/\mu_{av1}}{(\mu_b' - \mu_r')/\mu_{av2}}$$

~~$\Rightarrow \frac{1}{f_1} = \frac{2.1}{f_2}$~~

~~$\Rightarrow \frac{2.1}{f_2} + \frac{1}{f_2} = \frac{1}{20}$~~

~~$f_1 = 10.41$~~

~~$f_2 = -22$~~

$$w_1 = \frac{\mu_b - \mu_r}{\left(\frac{\mu_b + \mu_r}{2} - 1\right)} = 0.01546$$

$$w_2 = 0.02715$$

$$\frac{w_1}{w_2} = 0.5694 \Rightarrow \frac{1}{f_2} = -\frac{1}{f_1} \cdot 0.5694$$

$$\Rightarrow \frac{1}{f_1} = -\frac{1.75}{f_2}$$

$$\Rightarrow -\frac{0.75}{f_2} = \frac{1}{20} \Rightarrow f_2 = -15.1 \text{ cm}$$

$$f_1 = 8.6 \text{ cm}$$

Note the
formula

Q7)

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$$S = \frac{h^2 \phi(k, n)}{f}$$

$$\text{Now } \frac{S_2}{S_1} = \frac{h_2^2}{h_1^2} = \left(\frac{2}{6}\right)^2$$

$$\Rightarrow S_2 = \frac{1}{9} \times 1.8 \text{ cm} = 0.2 \text{ cm}$$

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Q8)

$$\begin{aligned} h_1^2 : h_2^2 : h_3^2 \\ = 25 : (144 - 100) : 144 \\ = 25 : 44 : 144 \end{aligned}$$

Q9)

Lecture 1

Q10)

System Matrix for lens is given as

$$\begin{bmatrix} 1 & -(\frac{1}{f_1}) \\ 0 & 1 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

The total matrix is, therefore

$$\begin{aligned} & \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & -10 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0.1 & 1 \end{bmatrix} \begin{bmatrix} 1 & -5 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ +0.2 & 1 \end{bmatrix} \\ & = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & -10 \\ 0.1 & 1 \end{bmatrix} \begin{bmatrix} 0 & -5 \\ 0.2 & 1 \end{bmatrix} \end{aligned}$$

In Matrix Method, we transverse from one point to another, so sign will be accordingly!!



Very Important Point: it will be +0.2
∵ we are going towards right so $\left(\frac{D}{n_1}\right) > 0$. So
sir ji were right in using +D₁ and +D₂

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$$= \begin{bmatrix} 1 & 0 \\ v & 1 \end{bmatrix} \begin{bmatrix} -2 & -10 \\ 0.2 & 0.5 \end{bmatrix}$$

$$= \begin{bmatrix} -2 & -10 \\ -2v + 0.2 & -10v + 0.5 \end{bmatrix}$$

Now, $-2v + 0.2 = 0$

$\Rightarrow v = 0.1$ ✓

$\frac{y_2}{y_1} = 0.5 - 1 = -0.5 = -\left(\frac{1}{2}\right)$

Magnification of $\left(\frac{1}{2}\right)$ and inverted image will be formed. $y_2 = -0.005$

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long way

① find u_o, v_o
for unit
planes

$u_o = -0.1$

$v_o = -0.05$

② find $f = 0.1$

$\Rightarrow u_{eff} = -0.3$
 $f = 0.1$

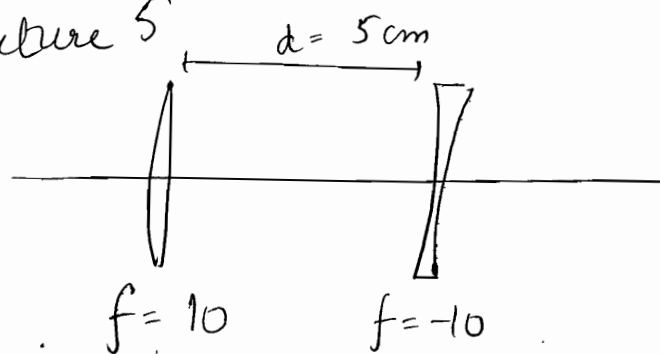
$\Rightarrow v = \left(\frac{3}{20}\right)$

$v_{eff} = 0.1$

Q11)

Q12)

lecture 5



$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2} - \frac{t}{f_1 f_2}$$

$$= \frac{1}{10} + \frac{1}{10} + \frac{5}{10 \cdot 10} = \left(\frac{1}{20}\right)$$

$f = 20 \text{ cm}$

Power $= \frac{1}{f} = 0.05$
 $= 5 \text{ diopter}$

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To find principal points :

Unit point = principal point

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Q13)

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Q14)

lecture 6

Fermat's Method : lecture 5

$$\gamma = \gamma_0 = \frac{c}{\lambda} = \frac{3 \times 10^8}{600 \times 10^{-9}} = \underline{\underline{5 \times 10^{14} \text{ Hz}}}$$

$$v = \frac{c}{n} = \frac{3 \times 10^8}{3} \times 2 = \underline{\underline{2 \times 10^8 \text{ m/s}}}$$

$$\lambda = \frac{v}{\gamma} = \frac{2 \times 10^8}{5 \times 10^{14}} = 4 \times 10^{-7} = \underline{\underline{400 \text{ nm}}}$$

Q15)

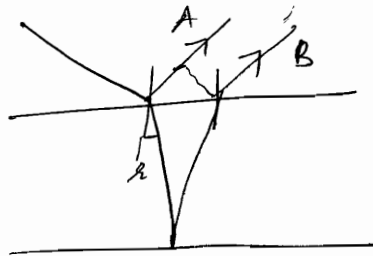
lecture 5

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Tut Sheet 9 : Interference

Q1)

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Path difference between ray A and ray B
is $2\mu t \cos r = \frac{\lambda}{2}$

Now from snell's law

$$\mu_1 \sin i = \mu_2 \sin r$$

$$\Rightarrow \sin 45^\circ = \mu_2 \sin r = 1.33 \sin r$$

$$\Rightarrow \sin r = 0.53$$

$$\Rightarrow \underline{\underline{\cos r = 0.84}}$$

For a destructive interference

$$\Delta = (2n-1) \frac{\lambda}{2}$$

$$\Rightarrow 2\mu t \cos r - \frac{\lambda}{2} = (2n-1) \frac{\lambda}{2}$$

$$\Rightarrow 2\mu t \cos r = n\lambda$$

$$\Rightarrow t = \left(\frac{n\lambda}{2\mu \cos r} \right)$$

For minimum thickness, $n=1$

$$\Rightarrow t = \frac{\lambda}{2\mu \cos r} = \frac{5890 \text{ \AA}}{2 \times 1.33 \times 0.84} = 0.2636 \mu\text{m}$$

Q2)

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For interference, intensity is given by

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \phi$$

where I_1 and I_2 are intensity due to first & second wave while ϕ is the phase difference.

$$\text{Now } I_{av} = I_1 + I_2 + 2\sqrt{I_1 I_2} \langle \cos 0 \rangle$$

$$= [I_1 + I_2] \checkmark$$
~~$$I_{av} = I_1 + I_2 + 2\sqrt{I_1 I_2} \langle \cos 0 \rangle$$~~

$$I_{max} = I_1 + I_2 + 2\sqrt{I_1 I_2}$$

$$= (\sqrt{I_1} + \sqrt{I_2})^2 = 1.1 (I_1 + I_2)$$

$$I_{min} = I_1 + I_2 - 2\sqrt{I_1 I_2}$$

$$= (\sqrt{I_1} - \sqrt{I_2})^2 = 0.9 (I_1 + I_2)$$

1st Method

$$\text{Let } \frac{I_1}{I_2} = x > 1$$

$$\Rightarrow$$
~~$$I_1 + I_2 + 2\sqrt{I_1 I_2} = 1.1(I_1 + I_2)$$~~

$$x + 1 + 2\sqrt{x} = 1.1(x + 1) \quad \text{--- (1)}$$

$$x + 1 - 2\sqrt{x} = 0.9(x + 1)$$

$$4\sqrt{x} = 0.2(x + 1)$$

$$20\sqrt{x} = x + 1$$

$$400x = x^2 + 1 + 2x$$

~~$$x^2 - 398x + 1 = 0$$~~

$$x^2 - 398x + 1 = 0 \Rightarrow x = \frac{398 \pm \sqrt{398^2 - 4}}{2} \approx 398$$

($\because x > 1$)

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2nd Method

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$$\Rightarrow \frac{\sqrt{I_1} + \sqrt{I_2}}{\sqrt{I_1} - \sqrt{I_2}} = \sqrt{\frac{1.1}{0.9}} = 1.105$$

$$\Rightarrow \frac{\sqrt{I_1}}{\sqrt{I_2}} = \frac{2.1}{0.1}$$

$$\Rightarrow \frac{I_1}{I_2} = \frac{401.9}{1} \approx \underline{402}$$

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remember
this

Q3)

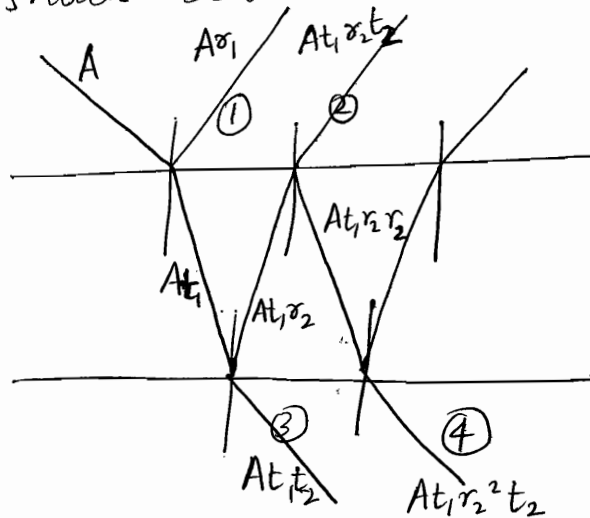
air
to
film

Let the index of the film be 1.33

$$\Rightarrow \text{Transmission Ratio} = \frac{2\mu_1}{\mu_1 + \mu_2} = \frac{2}{1 + 1.33} = 0.858 = t_1 \text{ (say)}$$

$$\text{Reflection ratio} = \frac{A_r}{A_i} = \frac{\mu_1 - \mu_2}{\mu_1 + \mu_2} = \frac{1 - 1.33}{1 + 1.33} = -0.14 = r_1 \text{ (say)}$$

Let us consider the multiple reflections as shown below



Similarly

film
to
air

$$\left[\begin{aligned} t_2 &= \frac{2\mu_1}{\mu_1 + \mu_2} = \frac{2 \times 1.33}{2.33} = 1.141 \\ r_2 &= \frac{1.33 - 1}{1.33 + 1} = 0.14 \end{aligned} \right]$$

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Consider waves ③ and ②

$$A_1 = A r_1 = \underline{0.14 A}$$

$$A_2 = A t_1 t_2 r_2 = A (0.858) (1.141) 0.14 \\ = 0.137 A \approx \underline{0.14 A}$$

Consider wave ③ and ④

$$A_3 = A t_1 t_2 = A (0.858) (1.141) \text{ complete darkness} \\ = \underline{0.97 A}$$

$$A_4 = A t_1 t_2 r_2^2 = 0.97 \times (0.14)^2 A \\ = \underline{0.02 A}$$

→ Rather give
the proof of

B.S.
"Aparna"

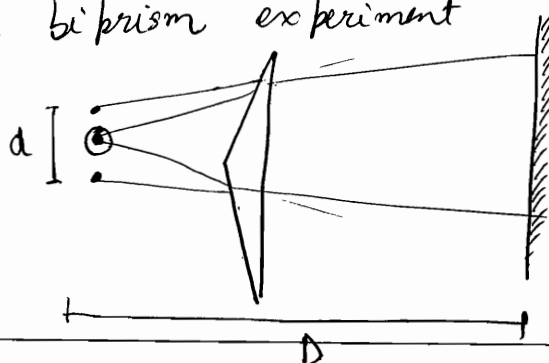
In transmitted light, the ray 2' differs from ray 1' in amplitude ratio r_2^2 . Therefore the interfering transmitted rays are of very unequal magnitude & hence the contrast of maxima and minima is poor.

Now since for reflected waves, we have almost similar intensities $\Rightarrow$ contrast in interference will be more while

For transmitted waves, difference between intensities is so much that there will be no contrast.

Q4)

For a biprism experiment



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$$d \sin \theta = n \lambda \quad (\text{for constructive interference})$$

$$\Rightarrow d \cdot \frac{y_n}{D} = n \lambda$$

$$\Rightarrow d \frac{\beta}{D} = \lambda$$

$$\Rightarrow \beta = \left(\frac{D \lambda}{d} \right)$$

When mica sheet introduced, path
difference for central fringe
 $= d \sin \theta = (\mu - 1) t$

$$\Rightarrow d \cdot \frac{y_n}{D} = (\mu - 1) t$$

$$\Rightarrow d \cdot \frac{1.89}{D} = 0.59 t$$

$$\Rightarrow \frac{\lambda}{\beta} \cdot 1.89 = 0.59 t$$

$$\Rightarrow 5900 \text{ \AA} \times \frac{1.89}{0.43} = 0.59 t$$

$$\Rightarrow t = 4.395 \mu\text{m}$$

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Q5)

Lecture 7

Q6)

Colours produced in thin films

When a thin film of oil on water, or a soap bubble, exposed to an extended to a white light ~~of different colours~~ (such as sky)

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is observed under reflected light, brilliant colours are seen in the film or the bubble. These colours arise due to the interference of the light waves reflected from the top and bottom surface of the film.

At a particular point of the film, and for a particular position of the eye (i.e. for a particular t and n), rays of only certain wavelengths will have a path difference satisfying the condition of maxima. Hence only these colours will be visible with maximum intensity. The colouration will vary with the thickness of the film. as well as with the position of eye. If different points of the film are observed with eye in the same position, a different set of colours will be observed.

In the problem, considering normal incidence,

For appearing dark in reflected light

$$\Delta = (2n+1) \frac{\lambda}{2}$$

$$\Rightarrow 2\mu t - \frac{\lambda}{2} = (2n+1) \frac{\lambda}{2}$$

$$\Rightarrow 2\mu t = n\lambda \Rightarrow t = \left(\frac{n\lambda}{2\mu} \right)$$

For minimum thickness, $n=1 \Rightarrow t = \frac{\lambda}{2\mu} = 0.2 \mu\text{m}$

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Q7)

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lecture 5

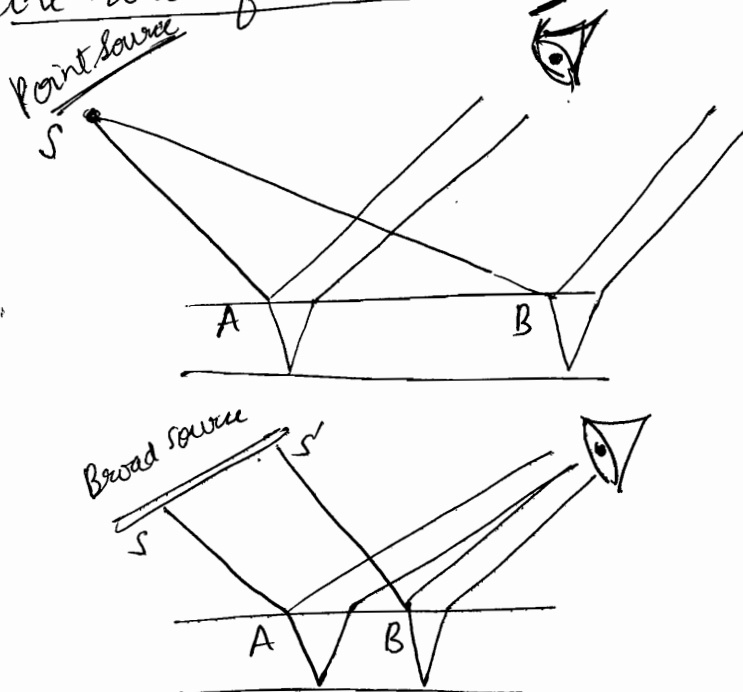
$$\frac{I_{\max}}{I_{\min}} = \frac{(\sqrt{I_1} + \sqrt{I_2})^2}{(\sqrt{I_1} - \sqrt{I_2})^2} = \left(\frac{2+1}{2-1}\right)^2 = \frac{9}{1}$$

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Q8)

REQUIREMENT OF EXTENDED SOURCE

In the case of interference exhibited by thin films (eg. soap film) an extended (broad) source of light is required so that light from each point on film may reach the eye of the observer to facilitate the whole film to be viewed at a glance.

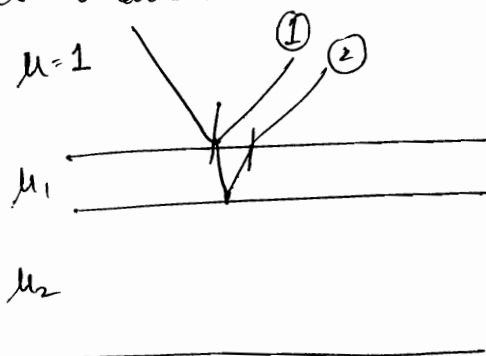


For example, Consider a narrow source. SA reaches the eye but SB does not. Thus only limited portion of the film is visible.

Now consider a broader extended source. Both SA and S'B could reach the eye. Therefore a large portion of film is visible and colours can be seen.

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For non reflecting film, a film of
lower index will be used



$$\mu_1 < \mu_2$$

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$$\Rightarrow \Delta_{12} = 2\mu t \cos \alpha \approx 2\mu t \quad (\text{Incident})$$

For non reflectivity, destructive interference
is required.

$$\Rightarrow 2\mu t = (2n-1) \frac{\lambda}{2}$$

$$\Rightarrow t = \frac{(2n-1)\lambda}{4\mu} = (2n-1) \cdot \frac{5.5 \times 10^{-5}}{\frac{5}{4} \times 4}$$

$$= (2n-1) \cdot 1.1 \times 10^{-5} \text{ cm}$$

For minimum thickness, $t = \underline{0.11 \mu\text{m}}$

Q9)

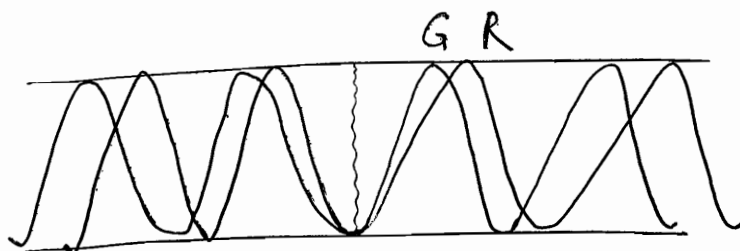
WHITE LIGHT IN MICHELSON INTERFERENCE

If a white light source (eg.
Incandescent lamp) is used, no
fringes will be seen except when the path
difference is so small that is less than
few wavelengths, only then few curved
and coloured localized fringes are
obtained.

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For a white light source, only a few fringes are observed because light contains all wavelengths between $4000 \text{ \AA}$ to $7500 \text{ \AA}$. Fringes in different colour will coincide.

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Intensity distribution for green light and red light is shown. The central fringe is uncoloured and different colours will begin to separate on either sides.

After about 10 fringes, many colours will tend to overlap at a point resulting in white.

Q10)

$$\Delta \text{central point} = (\mu - 1)t \quad \left(\text{Ignoring diffraction effects} \right)$$

$$\phi = \frac{2\pi}{\lambda} (\mu - 1)t$$

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos\left(\frac{2\pi}{\lambda} (\mu - 1)t\right)$$

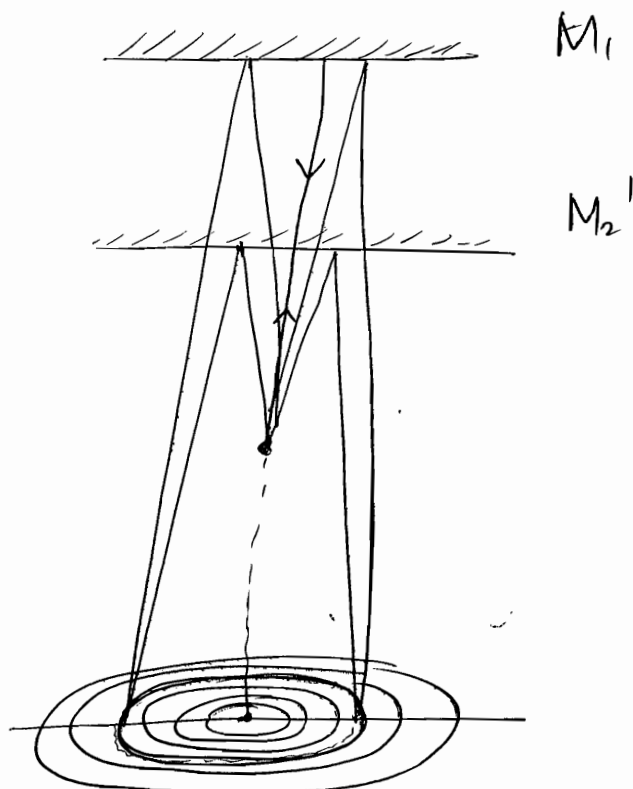
$$= 2I \left[1 + \cos\left(\frac{2\pi}{\lambda} (\mu - 1)t\right) \right]$$

$$I = 4I \cos^2\left(\frac{\pi}{\lambda} (\mu - 1)t\right)$$

[Note that even if we take diffraction into account $\left(\frac{\sin \alpha}{\alpha}\right) \rightarrow 1 \because \theta = 0$ for central point]

Q11)

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Circular Fringes

When M_1 and M_2 are exactly perpendicular
i.e. M_1 and M_2' are exactly parallel,
the film M_1, M_2' is of uniform
thickness, and we obtain circular
fringes localized at infinity.

$$\Delta = 2d \cos r + \frac{\lambda}{2}$$

For maximum intensity $\Delta = n\lambda$
Since d is constant, Haidinger
rings of constant inclination are
obtained i.e. difference rings for
different angles of refraction.

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Straight Fringes

When M_1 and M_2 are not $\perp$, we have similar configuration as wedge shaped film. $\therefore$ straight fringes are obtained.

Extended Source

In order to view the whole mirror filled with light we use extended source of light.

For change in distance d ,

$$\text{Change in } \Delta = 2d$$

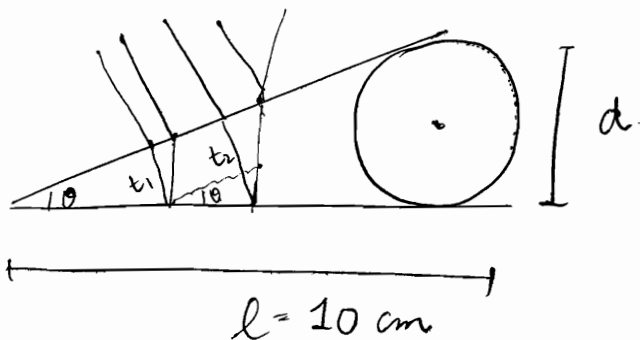
Now if $2d = N\lambda$, N fringes are shifted

Q12)

$$\lambda = 5000 \text{ \AA}$$

$$\beta = 1.25 \text{ cm}$$

$$d = ?$$



Since intervening medium is air, $\mu = 1$

$$\Rightarrow \Delta_1 = 2t_1 = n\lambda$$

$$\Delta_2 = 2t_2 = (n+1)\lambda$$

$$\Rightarrow t_2 - t_1 = \lambda/2 \Rightarrow \frac{t_2 - t_1}{\beta} = \tan \theta \approx \frac{d}{l}$$

$$\Rightarrow \beta = \left(\frac{\lambda l}{2d} \right)$$

$$\Rightarrow d \approx \frac{\lambda l}{2\beta} = 2 \mu\text{m}$$

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Q13)

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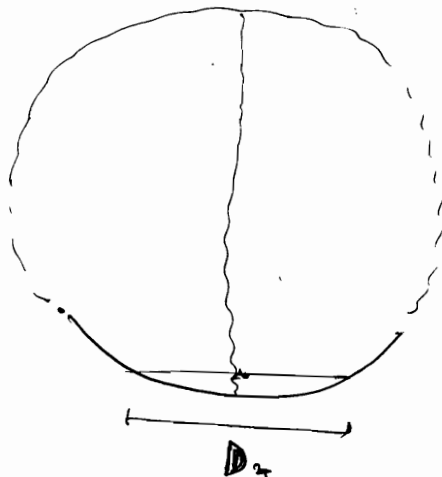
Q14)

Q-11 lecture 9

$$\left(\frac{D}{2}\right)^2 = t \cdot (2R - t)$$

$$\approx 2Rt$$

$$D^2 = 8Rt$$



Now for bright rings
 $\Delta = n\lambda$

$$\Rightarrow 2\mu t + \frac{\lambda}{2} = n\lambda$$

$$\Rightarrow t = \frac{(2n-1)\lambda}{2\mu} = (2n-1) \frac{\lambda}{4\mu}$$

$$\Rightarrow \frac{D_n^2}{8R} = (2n-1) \frac{\lambda}{4\mu}$$

$$\Rightarrow \frac{(2.303 \times 10^{-6})^2}{8R} = \frac{9\lambda}{4\mu}$$

$$\frac{(4.134 \times 10^{-6})^2}{8R} = \frac{29\lambda}{4\mu}$$

$$\mu = 1.5$$

$$\lambda = 5282 \text{ \AA}$$

$$\Rightarrow R = 0.84 \text{ m}$$

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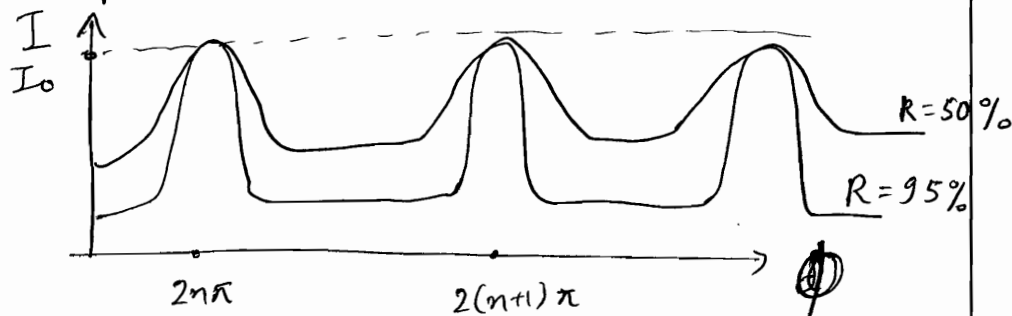
Q15)

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Lecture 9

SHARPER FRINGES IN FABRY PEROT

The maxima of the fringes in this case is very sharp & narrow as compared to Michelson.



The reason for this is found in the fact that there are several interfering beams. The intensity I_t is given by

$$I_{\text{total}} = \frac{I_0}{1 + \frac{4R}{(1-R)^2} \sin^2\left(\frac{\phi}{2}\right)}$$

where ϕ is the phase difference between successive rays and R is reflectivity.

For maxima, $\delta = 2n\pi \Rightarrow I_{\text{total}} = I_0$

When $\frac{\delta}{2} \neq n\pi$, $(1-R) \rightarrow 0$ due to

large R , and I_{total} drops sharply

When the reflectance is large & the intensity is due to superposition of several beams, the maxima are sharp & narrow.

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Q16)

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$$E_1 = 2 \cos (\vec{k}_1 \cdot \vec{r} - \omega t + \frac{\pi}{3}) \text{ kV/m}$$

$$E_2 = 5 \cos (\vec{k}_2 \cdot \vec{r} - \omega t + \frac{\pi}{4}) \text{ kV/m}$$

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$$I_1 = \frac{1}{2} \epsilon E_1^2 c = \frac{1}{2} \epsilon \cdot 4 \cdot c \times 10^6 \text{ W/m}^2$$

$$I_2 = \frac{1}{2} \epsilon E_2^2 c = \frac{1}{2} \epsilon \cdot 25 \cdot c \times 10^6 \text{ W/m}^2$$

$$I_{12} = (\sqrt{I_1} + \sqrt{I_2})^2 = (2+5)^2 \cdot \frac{1}{2} \epsilon \cdot c \times 10^6$$

$$= \frac{49}{2} \epsilon c \times 10^6 \text{ W/m}^2$$

$$I_{\min} = (2-5)^2 \frac{1}{2} \epsilon c \times 10^6 = \frac{9}{2} \epsilon c \times 10^6 \text{ W/m}^2$$

$$\Rightarrow V = \frac{49 - 9}{49 + 9} = \left(\frac{40}{58} \right)$$

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Q17)

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Q18)

Q19)

Lecture 7

Q3

$$\text{Coherence length} = 2d = \frac{\lambda^2}{\Delta\lambda}$$

$\therefore$ at this length, the pattern disappears.

~~For~~

For pattern to disappear

$$\frac{2d}{\lambda} - \frac{2d}{\lambda + \frac{\Delta\lambda}{2}} \geq \frac{1}{2}$$

$$\Rightarrow \frac{2d}{\lambda} - \frac{2d}{\lambda} \left(1 - \frac{\Delta\lambda}{2\lambda}\right) \geq \frac{1}{2}$$

$$\Rightarrow \frac{2d}{\lambda} \cdot \frac{\Delta\lambda}{2\lambda} \geq \frac{1}{2}$$

$$\Rightarrow \boxed{2d \geq \left(\frac{\lambda^2}{\Delta\lambda}\right)}$$

$$\Rightarrow l = 2d = \frac{(600)^2 \times 10^{-18}}{(0.01) \times 10^{-9}}$$

$$= \underline{\underline{0.036 \text{ m}}}$$

Q20)

$$\vec{E}_1 = E_1 \sin(kx - \omega t) \hat{i}$$

$$\vec{E}_2 = E_2 \sin(kx - \omega t + \phi) \hat{j}$$

$$(\vec{E}_1 + \vec{E}_2) = E_1 \sin(\dots) \hat{i} + E_2 \sin(\dots + \phi) \hat{j}$$

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$$I \propto E^2$$

$$\rightarrow I = \langle (\vec{E}_1 + \vec{E}_2) \cdot (\vec{E}_1 + \vec{E}_2) \rangle \quad \checkmark \quad \text{Perfect}$$

$$= \langle E_1^2 \sin^2(kx - \omega t) + E_2^2 \sin^2(kx - \omega t + \phi) \rangle$$
~~$$+ 2E_1 E_2 \sin(kx - \omega t) \sin(kx - \omega t + \phi)$$~~

$$I = E_1^2 \langle \sin^2(kx - \omega t) \rangle + E_2^2 \langle \sin^2(kx - \omega t + \phi) \rangle$$

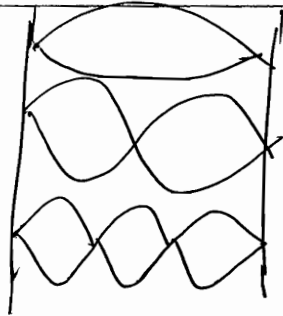
$$= \frac{(E_1^2 + E_2^2)}{2}$$

$$= \text{Independent of } \phi$$

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(Q21)

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For Fabry Perot etalon,
each mode differs by :

$$2d = n\lambda_1$$

$$2d = (n+1)\lambda_2$$

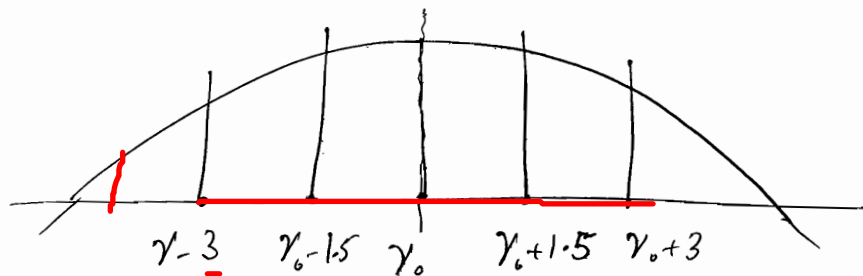
~~$\Rightarrow \frac{2d}{\lambda_1} - \frac{2d}{\lambda_2} = 1$~~

$$\frac{2d}{\lambda_2} - \frac{2d}{\lambda_1} = 1$$

$$\frac{2d(\nu_2 - \nu_1)}{c} = 1$$

$$\Rightarrow \Delta\nu = \frac{c}{2d}$$

between 2 successive
modes



$$\Delta\nu = \frac{3 \times 10^8}{2 \times 100 \times 10^{-3}} = 1.5 \text{ GHz}$$

5 modes can be supported.

Q

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Tut Sheet 10 Diffraction

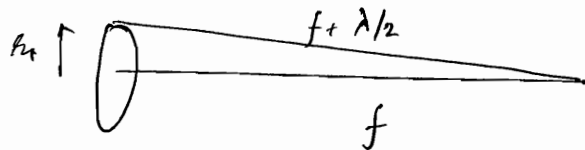
Q1)

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Considering the first image as the 1st
bright image i.e. when the Central
Zone acts as ~~central~~ 1st Half Period
Zone, we have

Focal length f is calculated as

1st way



$$f^2 + h_1^2 = \left(f + \frac{\lambda}{2}\right)^2$$

$$\Rightarrow h_1^2 = \frac{\lambda^2}{4} + 2f\frac{\lambda}{2}$$

Now $f \gg \lambda$

$$\Rightarrow \boxed{\frac{h_1^2}{\lambda} = f}$$

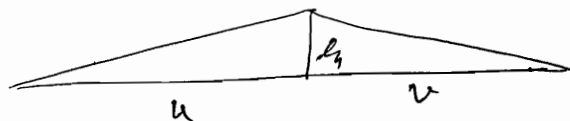
$\Rightarrow$ Using

$$\frac{1}{v} - \frac{1}{u} = \frac{1}{f} \Rightarrow \frac{1}{v} + \frac{1}{6} = \frac{5.89 \times 10^{-3}}{\left(\frac{2 \cdot 3}{2}\right)^2}$$

we have

$$\Rightarrow \boxed{v = 3.58 \text{ m}}$$

2nd way



$$\sqrt{u^2 + h_1^2} + \sqrt{v^2 + h_1^2} - u - v = \frac{\lambda}{2}$$

$$\Rightarrow \frac{u \cdot \frac{h_1^2}{2u^2}}{2u^2} + \frac{v \cdot \frac{h_1^2}{2v^2}}{2v^2} = \frac{\lambda}{2}$$

$$\Rightarrow \frac{1}{u} + \frac{1}{v} = \frac{2\lambda}{2\lambda^2} = \frac{\lambda}{\lambda^2} = \frac{1}{\left(\frac{\lambda^2}{\lambda}\right)}$$

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Q2)

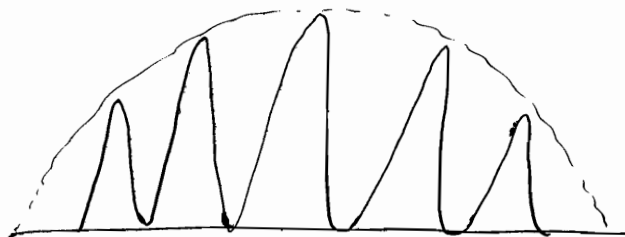
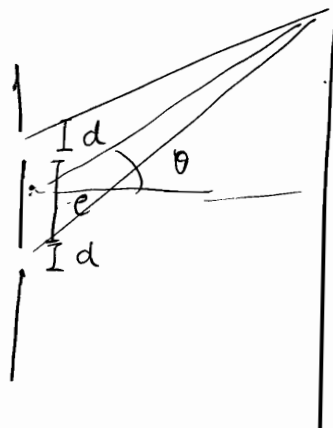
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For double slit fraunhofer diffraction

$$I = I_0 \frac{\sin^2 \alpha}{\alpha^2} \cos^2 \beta$$

$$\alpha = \frac{\pi d \sin \theta}{\lambda}$$

$$\beta = \frac{\pi (e+d) \sin \theta}{\lambda}$$



For interference maxima, we have

$$\beta = n\pi$$

$$\pi \frac{(e+d) \sin \theta}{\lambda} = n\pi$$

$$(e+d) \sin \theta = n\lambda$$

For calculating fringe width (F.W.)

$$\frac{y_{n+1}}{D} = \frac{(n+1)\lambda}{(e+d)}$$

$$\frac{y_n}{D} = \frac{n\lambda}{e+d}$$

$$\Rightarrow \frac{F.W.}{D} = \frac{\lambda}{e+d}$$

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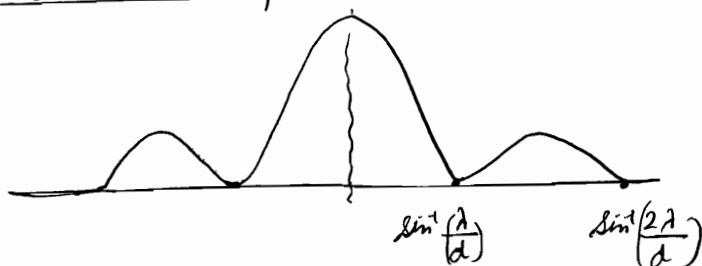
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$$\Rightarrow F.W = \left(\frac{\lambda D}{e+d} \right) = \frac{4800 \text{ \AA} \times 50 \text{ cm}}{0.1 \text{ mm} + 0.02 \text{ mm}}$$

$$= \underline{2 \text{ mm}}$$

For minima first, $\alpha = \pi \Rightarrow \theta = \left(\frac{\lambda}{d} \right) = 24 \text{ mm} \Rightarrow y = \theta \cdot D = \underline{12 \text{ mm}}$

Q3)



For diffraction pattern

$$I = I_0 \frac{\sin^2 \alpha}{\alpha^2} \quad \alpha = \frac{\pi d \sin \theta}{\lambda}$$

For 1st minima $\alpha = \pi \Rightarrow \theta = \sin^{-1} \left(\frac{\lambda}{d} \right)$

2nd minima $\Rightarrow \theta = \sin^{-1} \left(\frac{2\lambda}{d} \right)$

The position of 2nd dark minima, y_2 is

$$y_2 = D \cdot \sin^{-1} \left(\frac{2\lambda}{d} \right)$$

$$\Rightarrow 2 \cdot \sin^{-1} \left(\frac{2\lambda}{d} \right) = 1.6 \times 10^{-2}$$

$$\Rightarrow \frac{2\lambda}{d} \approx 0.8 \times 10^{-2}$$

$$\Rightarrow \lambda = \frac{0.8 \times 10^{-2} \times 0.14 \times 10^{-3}}{2} = \underline{5600 \text{ \AA}}$$

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Q4)

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For a 2-slit Fraunhofer diffraction,

$$I = I_0 \frac{\sin^2 \alpha}{\alpha^2} \cdot \cos^2 \beta$$

given : $d = 2\lambda$
 $e+d = 6\lambda$

$$\alpha = \frac{\pi d \sin \theta}{\lambda}$$

$$\beta = \frac{\pi (e+d) \sin \theta}{\lambda}$$

We have for missing order

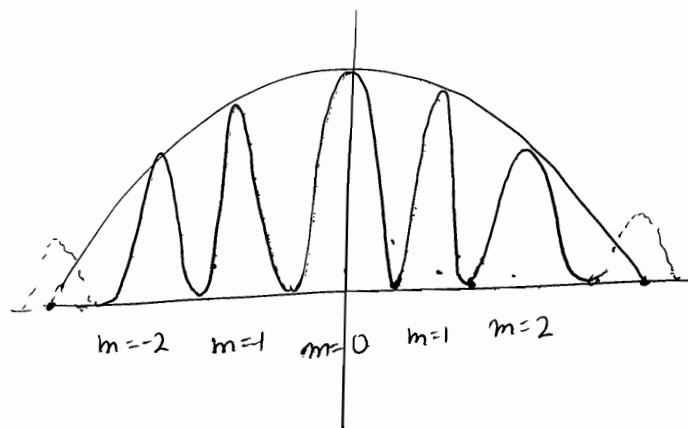
$$\cos^2 \beta = 1 \quad \text{and} \quad \alpha = n\pi$$

$$\text{i.e.} \quad \beta = m\pi \quad \text{and} \quad \alpha = n\pi$$

$$\Rightarrow \frac{\beta}{\alpha} = \frac{m}{n}$$

$$\Rightarrow \frac{e+d}{d} = \frac{m}{n}$$

$$\Rightarrow \boxed{\frac{m}{n} = 3}$$



Q5)

Lecture 12

Grating equation, for principal maxima
of interference

$$\beta = n\pi \Rightarrow \pi (e+d) \sin \theta = n\lambda$$

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$$\Rightarrow (e+d) \sin \theta = n \lambda$$

$$\Rightarrow (e+d) \cos \theta \cdot d\theta = n \cdot d\lambda$$

$$\Rightarrow \boxed{\frac{d\theta}{d\lambda} = \frac{n}{(e+d) \cos \theta}}$$

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Angular
Dispersion

or
dispersive
power

Prism Spectrum

- ① Produces only 1 spectrum.
- ② Spectrum is pure one.
- ③ This will give a bright spectrum since all light is concentrated on 1 spectrum.
- ④ Dispersive power depends on λ .
 $\Rightarrow$ IRRATIONAL SPECTRUM

Grating Spectrum

Gives a no. of spectrum on either side of zero order

This is impure at higher orders due to overlapping of spectral lines

Most of the energy goes into zero order where no spectrum is formed. Rest of the energy is distributed among the spectra of various orders

Dispersive power depends on order of spectrum n but independent of λ . $\Rightarrow$

RATIONAL SPECTRUM

Q6) lecture 12

Q7) ~~separation between~~ given $e = 4a$
 $d = a$

$$\text{For } \frac{e+d}{d} = \frac{m}{n} \Rightarrow \frac{5a}{a} = \frac{m}{n}$$

$$\Rightarrow m = 5n = \underline{5} \quad (n=1)$$

$$\text{Fringes} = 4 + 1 + 4 = 9 \text{ fringes}$$

Q8)

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Q9)

Q10)

Q11)

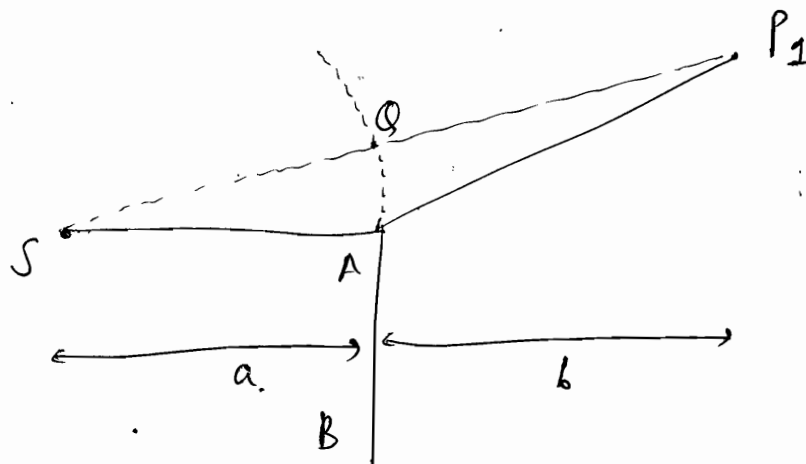
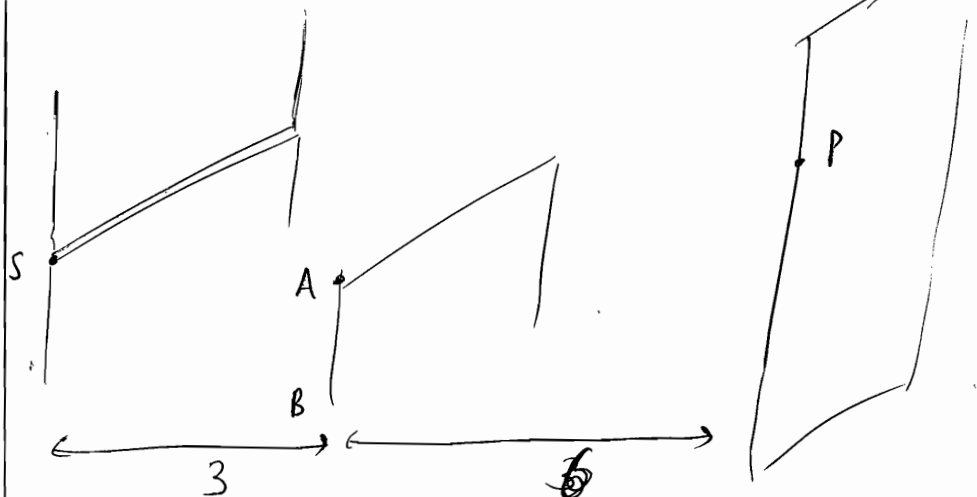
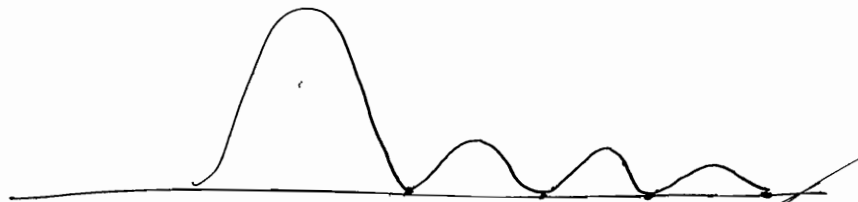
Q12)

lecture 14

lecture 14

X

X (NO Disturbance....) अरे भय्या, ... टिमाग नहीं !!



Point P will be dark, when even no. of HPZ

$$AS + AP - PS = (2n+1) \frac{\lambda}{2}$$

$$AP - QP = (2n+1) \frac{\lambda}{2}$$

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do not
write the
condition
for normal
interference

remember
this step

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~~For 1st dark Band~~
 ~~$AS + AP_1 = SP_1$~~
~~For 2nd dark Band~~
 ~~$AS + AP_2 = SP_2$~~

~~$AP_2 = ?$~~

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$$AP^2 = b^2 + y_1^2 \Rightarrow AP_1 = \sqrt{b^2 + y_1^2} = b \left(1 + \frac{y_1^2}{2b^2} \right)$$

$$OP_1 = SP - SQ = \sqrt{(a+b)^2 + y_1^2} - a$$

$$= (a+b) \left[1 + \frac{y_1^2}{2(a+b)^2} - a \right]$$

$$= b + \frac{y^2}{2(a+b)}$$

Simple
pythagoras
+ binomial

$$\Rightarrow (2n+1) \frac{\lambda}{2} = b + \frac{y_n^2}{2b} - b - \frac{y_n^2}{2(a+b)}$$

$$= \frac{y_n^2}{2} \left[\frac{1}{b} - \frac{1}{a+b} \right] = \frac{y_n^2}{2b(a+b)} a$$

$$\Rightarrow y_n = \sqrt{\frac{b(a+b) \lambda (2n+1)}{a}}$$

$$\text{Now } y_1 = \sqrt{\frac{b(a+b) \lambda \sqrt{2}}{a}}$$

$$y_4 = \sqrt{\frac{b(a+b) \lambda \sqrt{8}}{a}}$$

$$\Rightarrow \Delta y = \sqrt{\frac{6 \times 9 \times 6400 \times 10^{-10}}{3}} (\sqrt{8} - \sqrt{2})$$

$$= \underline{\underline{4.8 \text{ mm}}}$$

(Q13)

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(Q14)

Lecture 14

(Q1)

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Tut Sheet 11 resolving power

Q1)

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Resolving Power is primary of instruments
with circular aperture

eg. Telescope
Microscope.

So if nothing is given assume
circular aperture.....

For a circular aperture,

$$I = I_0 \left[\frac{J(\alpha)}{\alpha} \right]^2 \quad \alpha = \frac{\pi D \sin \theta}{\lambda}$$

Since $J(\alpha)$ is a damping sine
function, for the first ring, assuming

$J(\alpha) \Rightarrow \sin \alpha$, we have

$$I = I_0 \frac{\sin^2 \alpha}{\alpha^2} \quad \checkmark$$

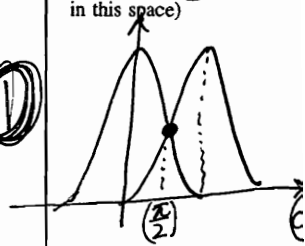
$$I_{\max} = I_0 \quad [\alpha = 0]$$

$$I_{\text{middle}} = I_0 \frac{\sin^2 \left(\frac{\pi}{4} \right)}{\left(\frac{\pi}{4} \right)^2} = \frac{8 I_0}{\pi^2}$$

$$\Rightarrow \frac{I_{\text{middle}}}{I_0} = \frac{8}{\pi^2}$$

कृपया इस स्थान
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$$I_{\text{middle}} = 2 I_0 \frac{\sin^2 \left(\frac{\pi}{4} \right)}{\frac{\pi^2}{4}}$$

$$= \frac{8 I_0}{\pi^2}$$

$$= \frac{8}{\pi^2} I_{\max}$$

Q 2)

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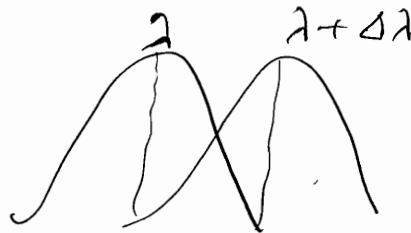
For diffraction grating, for principal
maxima

$$\beta = n\pi$$

$$\frac{n(e+d) \sin \theta}{\lambda} = n\pi$$

$$(e+d) \sin \theta = n\lambda$$

for maxima



For just resolved,

$$\theta_{\text{Maxima}(\lambda + \Delta\lambda)} = \theta_{\text{Minima} \lambda}$$

$$\Rightarrow \frac{n(\lambda + \Delta\lambda)}{e+d} = \frac{n\lambda}{e+d} + \frac{\lambda}{N(e+d)}$$

$$\Rightarrow \frac{n\Delta\lambda}{e+d} = \frac{\lambda}{N(e+d)}$$

$$\Rightarrow \frac{\lambda}{\Delta\lambda} = nN$$

$$\Rightarrow \frac{\lambda}{\Delta\lambda} = nN$$

Now for $N = 3 \times 10^4$, $n = 2$, $\lambda = 6000 \text{ \AA}$

$$\Delta\lambda = \frac{6000}{2 \times 3 \times 10^4} = \underline{\underline{0.1 \text{ \AA}}}$$

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Q3)

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$$\text{Resolving Power} : \frac{\lambda}{\Delta \lambda} = Nn$$

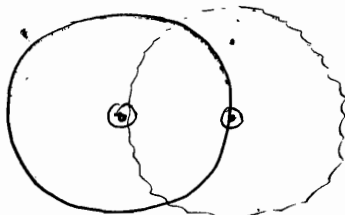
$$\text{Dispersive Power} = \left(\frac{d\theta}{d\lambda} \right) = \frac{n}{(e+d) \cos \theta}$$

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Q4)

For resolution in circular aperture,

$$1 \text{ sec} = \frac{1}{60 \times 60} \text{ degree}$$



$$1 \text{ degree} = \frac{\pi}{180} \text{ radian}$$

$$\alpha = 1.22 \pi$$

$$\Rightarrow \frac{\pi d \sin \theta}{\lambda} = 1.22 \pi$$

$$\Rightarrow \lambda = \left(\frac{d \sin \theta}{1.22} \right) = \frac{d \theta}{1.22} \Rightarrow d = \frac{1.22 \lambda}{\theta}$$

$$d = \frac{1.22 \times 5500 \times 10^{-10}}{\frac{1}{60 \times 60} \times \frac{\pi}{180}}$$

$$= \frac{2.41 \text{ mm} \times \frac{180}{\pi}}{\pi} = 13.8 \text{ cm}$$

do not
forget to
convert
degree to
radians.

Q5)

For the grating, the grating equation is

$$(e+d) \sin \theta = n\lambda$$

$$\Rightarrow \theta_1 - \theta_2 = n \frac{(\lambda_1 - \lambda_2)}{e+d}$$

$$= \frac{2 \times 6 \times 10^{-10} \times 6000}{1 \times 10^{-2}}$$

Not
Eqd: 2.5 cm

$$= 0.04127^\circ = 2.476 \text{ minutes}$$

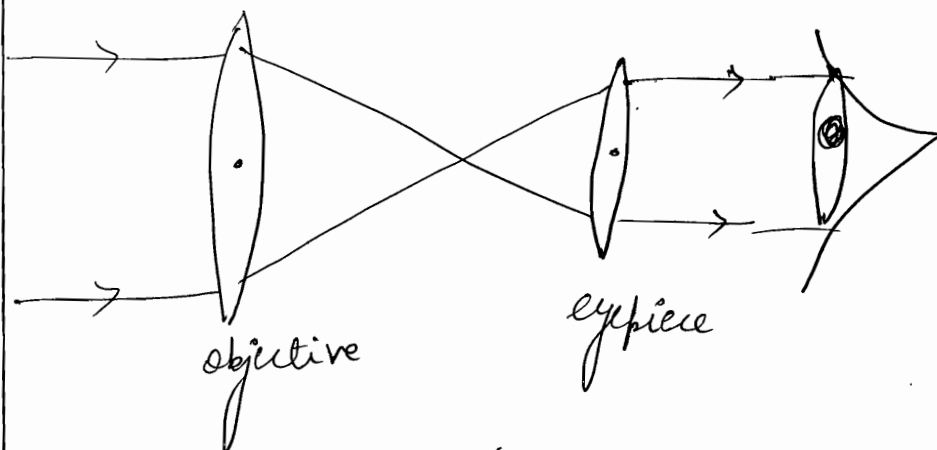
Q6)

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Q7)

Lecture 12

Magnifying Power of Telescope



$$\text{Magnification} = \left(\frac{\text{focal length of Objective}}{\text{focal length of eyepiece}} \right)$$

From the above figure,

$$\text{Magnification} = \frac{f_1}{f_2} = \left(\frac{d_1}{d_2} \right)$$

$$= \left(\frac{\text{Diameter of Objective}}{\text{Diameter of Eyepiece}} \right) = \left(\frac{D}{d} \right)$$

Now diameter of eyepiece = diameter of eye
 $d = d_e$

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Also, limit of resolution of telescope

$$d\theta_T = \frac{1.22 \lambda}{D}$$

Also, limit of resolution of eye

$$d\theta_e = \frac{1.22 \lambda}{d_e}$$

$$\Rightarrow \frac{d\theta_e}{d\theta_T} = \frac{1/d_e}{1/D} = \left(\frac{D}{d_e}\right) = \underline{M}$$

$$\Rightarrow \underline{\text{Optimum Magnification}} = (d\theta_e) \left(\frac{1}{d\theta_T}\right)$$

= Product of resolving power of telescope
and limit of resolution of eye.

Q8)

For grating

$$(e+d) \sin \theta = n\lambda$$

$$\Rightarrow \theta_1 - \theta_2 = \frac{n(\lambda_1 - \lambda_2)}{e+d}$$

$$= \frac{1 \times 6 \text{ \AA} \times 5000}{1 \text{ cm}}$$

$$= \underline{\underline{0.017^\circ}}$$

use a telescope of angular
magnification 10 or increase the no. of
grating per cm to 50000.

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(Q9)

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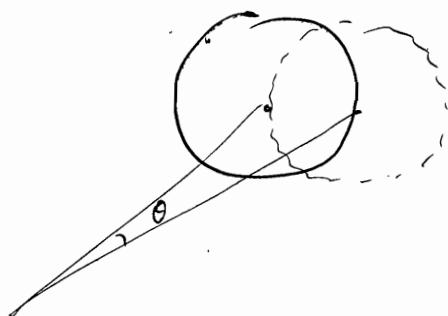
Resolving Power of Telescope = $\left(\frac{D}{1.22\lambda}\right)$

more the D, better the telescope.

Let the distance on the moon that
could be resolved be d

$$\Rightarrow \theta = \frac{d}{D_{EM}} \text{ where } D_{EM} = 384000 \text{ km}$$

Now for the 2 images to be just
resolved, using Rayleigh criteria



$$\alpha = 1.22\pi$$

$$\Rightarrow \frac{\pi d \sin \theta}{\lambda} = 1.22\pi$$

$$\Rightarrow \frac{d \cdot d}{\lambda D_{EM}} = 1.22$$

$$\Rightarrow d = \frac{1.22 \times \lambda \times D_{EM}}{d_e}$$

$$= \frac{1.22 \times 5890 \times 10^{-10} \times 3.84 \times 10^8}{5.08}$$

$$= \underline{\underline{54.31 \text{ m}}}$$

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Q10)

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Q11)

Resolving Power of an optical instrument represents its ability to produce distinctly separate spectral lines of light having two or more close wavelengths.

Rayleigh's Criterion of resolution states that two spectral lines of equal intensity are just resolved by an optical instrument when the principal maxima of the diffraction pattern due to one falls on the first minimum of the diffraction pattern of the other.

$\left(\frac{d\theta}{d\lambda}\right)$

Dispersive power of a grating is defined as the ratio of the difference in the angle of diffraction of any two neighbouring spectral lines to the difference in wavelength between the two spectral lines.

For grating,

$$(e+d) \sin \theta = n\lambda$$

$$\Rightarrow \frac{(e+d) \sin \theta_1}{(e+d) \sin \theta_2} = \frac{n\lambda_1}{n\lambda_2}$$

$$\Rightarrow \frac{(e+d) \sin \theta_1}{(e+d) \sin \theta_2} = \frac{\lambda_1}{\lambda_2}$$

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$$\theta_{\text{Maxima}_{(\lambda+\Delta\lambda)}} = \theta_{\text{Minima}_{\lambda}}$$

$$\frac{n(\lambda+\Delta\lambda)}{e+d} = \frac{n\lambda}{e+d} + \frac{\lambda}{N(e+d)}$$

$$\Rightarrow \frac{n\Delta\lambda}{e+d} = \frac{\lambda}{N(e+d)}$$

$$\Rightarrow \boxed{\frac{\Delta\lambda}{\lambda} = Nn}$$

$$\Rightarrow \frac{500}{0.5} = N \cdot 2$$

$$\Rightarrow N = 500 \text{ lines}$$

$$\Rightarrow 500 = 250 \text{ lines per cm} \times x \text{ cm}$$

$$\Rightarrow \boxed{x = 2 \text{ cm}}$$

Q12) @ lecture 12

Q13) (b) Now for grating,

$$(e+d) \sin \theta = m\lambda$$

$$\Rightarrow (e+d) \sin 30^\circ = 2\lambda$$

$$\Rightarrow \frac{1 \text{ cm}}{4250} \times \frac{1}{2} = 2\lambda \Rightarrow$$

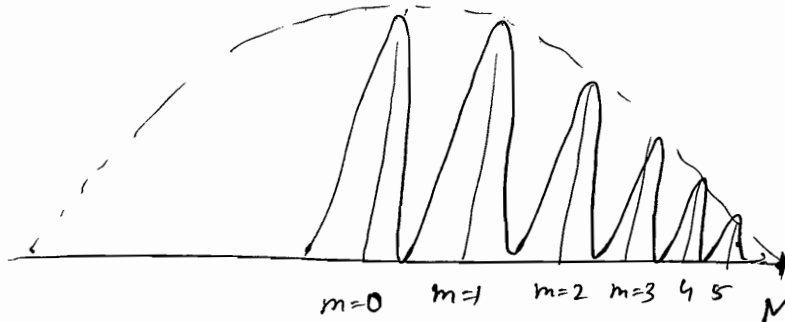
$$\boxed{\lambda = 5882 \text{ \AA}}$$

Q13)

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$$\frac{e+d}{d} = \frac{0.96}{0.16} = 6$$

$m=6$: missing
Order



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Missing
Order

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Tut Sheet 12 : Polarization

Q1)

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From linearly polarized light changing to
Circularly polarized light,

$$\Delta\phi = 90^\circ = \frac{\pi}{2}$$

$$\Delta = \frac{\lambda}{2\pi} \times \frac{\pi}{2} = \left(\frac{\lambda}{4}\right)$$

We know $\Delta = (\mu_e - \mu_o) t$

$$\begin{aligned} \Rightarrow (\mu_e - \mu_o) &= \frac{\lambda}{4t} = \frac{6 \times 10^{-7}}{4 \times 3 \times 10^{-5}} \\ &= \frac{1}{2} \times 10^{-2} \\ &= \underline{\underline{5 \times 10^{-3}}} \end{aligned}$$

Q2)

Let us consider 2 linearly polarized
light waves:

$$\begin{aligned} E_1 &= E_0 \cos \phi_1 \sin(kz - \omega t) \hat{i} \\ &\quad + E_0 \sin \phi_1 \sin(kz - \omega t) \hat{j} \end{aligned}$$

$$\begin{aligned} E_2 &= E_0 \cos \phi_2 \sin(kz - \omega t + \phi) \hat{i} \\ &\quad + E_0 \sin \phi_2 \sin(kz - \omega t + \phi) \hat{j} \end{aligned}$$

Now note that using two beams
will be very confusing. Rather than
take two linearly polarization beams
in $\perp$ directions. Both are equivalent.

As can be seen from addition of above 2 waves

along x and y axis, we can have

$$E_3 \sin(kz - \omega t + \phi_3) \hat{i}, E_4 \sin(kz - \omega t + \phi_4) \hat{j}$$

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Let the 2 polarized beams be

$$E_x = E_0 \sin(kz - \omega t) \hat{i}$$

$$E_y = E_0 \sin(kz - \omega t + \phi) \hat{j}$$

NOW

$$E_y = E_0 \sin(kz - \omega t) \cos\phi + E_0 \cos(kz - \omega t) \sin\phi$$

$$\Rightarrow E_y = E_x \cos\phi + E_0 \sqrt{1 - \frac{E_x^2}{E_0^2}} \sin\phi$$

$$\Rightarrow (E_y - E_x \cos\phi)^2 = (E_0^2 - E_x^2) \sin^2\phi$$

$$\Rightarrow E_y^2 + E_x^2 \cos^2\phi - 2E_x E_y \cos\phi = E_0^2 \sin^2\phi - E_x^2 \sin^2\phi$$

$$\Rightarrow \boxed{E_y^2 + E_x^2 - 2E_x E_y \cos\phi = E_0^2 \sin^2\phi}$$

For $\phi = (2n-1) \frac{\pi}{2}$, we have

$$E_y^2 + E_x^2 - 0 = E_0^2$$

$$\Rightarrow \boxed{E_y^2 + E_x^2 = E_0^2}$$

$$E_x = E_0 \sin(kz - \omega t) \hat{i}$$

$$E_y = E_0 \cos(kz - \omega t) \hat{j}$$

this
defines
a circularly
polarized
wave

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For even n , i.e. $\phi = -\frac{\pi}{2}, \frac{3\pi}{2}, \frac{7\pi}{2}$

We have

$$E_x = E_0 \sin(kz - \omega t) \hat{i}$$

$$E_y = -E_0 \cos(kz - \omega t)$$

LCP

For odd n , i.e. $\phi = \frac{\pi}{2}, \frac{5\pi}{2}, \frac{9\pi}{2}$

$$E_x = E_0 \sin(kz - \omega t)$$

$$E_y = E_0 \cos(kz - \omega t)$$

RCP

Step 1

Pass the beam through a rotating polaroid. The intensity for both will remain same. $\Rightarrow$ circular or unpolarized

Step 2

Take a QWP and pass the beam through it. Now perform step 1.

If intensity drops to 0 twice,
 $\Rightarrow$ initially it was CPW

If intensity remains same, then initially it was Unpolarized.

Q3) For QWP, $\phi = \frac{\pi}{2}$ i.e. $\Delta = \frac{\lambda}{2\pi} \cdot \frac{\pi}{2} = \frac{\lambda}{4}$

$$\Rightarrow (1.658 - 1.486) \cdot t = \frac{5890 \text{ \AA}}{4}$$

$$\Rightarrow t = 0.8561 \mu\text{m}$$

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Q4)

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Optical Rotation

When a plane polarized passes through certain substances, the plane of polarization of light is rotated about the direction of propagation of light through a certain angle. This phenomenon is called OPTICAL ROTATION or ROTATORY POLARIZATION.

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The substances which rotate the plane of polarization are called "Optically active" and the property is called "Optical Activity".

Fresnel Theory of Optical Rotation

It's based upon the principle that dynamics of a linear vibration may be described as the resultant of two opposite circular motions of the same frequency.

The incident plane-polarized light on entering a substance is broken up into two circularly-polarized waves, one clockwise and the other anti-clockwise.

In an optically active substance, the two waves travel with different velocities,
(In levorotatory, $v_{\text{anticlockwise}}$ is more,
in dextrorotatory, $v_{\text{clockwise}}$ is more)

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Hence a phase difference is developed between them as they transverse the substance.

On emergence, the two circular components, recombine to form plane polarized light by an angle depending upon the phase difference between them.

Let the vibration in incident light be represented by $y = A \cos \omega t$
These vibrations, upon entering the crystal, are broken down into two equal & opposite circular motions.

$$\left. \begin{aligned} x_1 &= \frac{a}{2} \sin(\omega t) \\ y_1 &= \frac{a}{2} \cos(\omega t) \end{aligned} \right\} \text{RCP (clockwise)}$$

$$\left. \begin{aligned} x_2 &= -\frac{a}{2} \sin(\omega t) \\ y_2 &= \frac{a}{2} \cos(\omega t) \end{aligned} \right\} \text{RCP (anticlockwise)}$$

These circular components are propagated through the quartz crystal with different velocities. Therefore, when they emerge, they have a phase difference δ between them.

$$\left. \begin{aligned} x_1 &= \frac{a}{2} \sin(\omega t + \delta) \\ y_1 &= \frac{a}{2} \cos(\omega t + \delta) \end{aligned} \right\} \quad \left. \begin{aligned} x_2 &= -\frac{a}{2} \sin(\omega t) \\ y_2 &= \frac{a}{2} \cos(\omega t) \end{aligned} \right\}$$

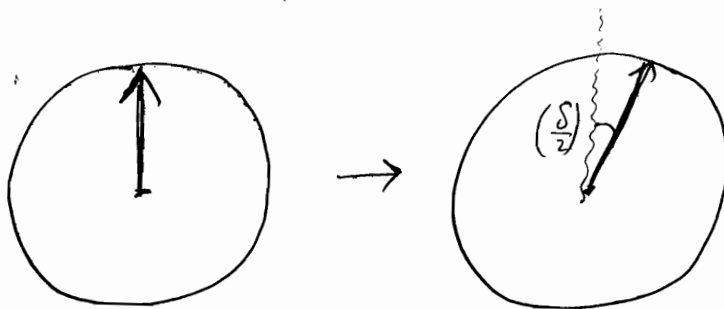
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The resultant displacement along the two
axis are :

$$\begin{aligned} x &= \frac{a}{2} [\sin(\omega t + \delta) - \sin(\omega t)] \\ &= \frac{a}{2} [2 \sin\left(\frac{\delta}{2}\right) \cos\left(\omega t + \frac{\delta}{2}\right)] \\ &= \underline{\underline{a \sin\left(\frac{\delta}{2}\right) \cos\left(\omega t + \frac{\delta}{2}\right)}} \end{aligned}$$

$$\begin{aligned} y &= \frac{a}{2} [\cos(\omega t + \delta) + \cos \omega t] \\ &= \frac{a}{2} [2 \cos\left(\omega t + \frac{\delta}{2}\right) \cos\left(\frac{\delta}{2}\right)] \\ &= \underline{\underline{a \cos\left(\frac{\delta}{2}\right) \cos\left(\omega t + \frac{\delta}{2}\right)}} \end{aligned}$$



$$\Rightarrow \text{Rotation } \theta = \left(\frac{\delta}{2}\right)$$

$$\theta = \frac{\delta}{2} = \frac{(\mu_A - \mu_B) t}{2\lambda} \cdot 2\pi = \underline{\underline{\frac{\pi t}{\lambda} (\mu_A - \mu_B)}}$$

$$\boxed{\theta \propto \frac{1}{\lambda}}$$

$$\theta = \frac{\pi * 10^{-3} \times 12 \times 10^{-5}}{4500 \times 10^{-10}} = \underline{\underline{48^\circ}}$$

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Q4

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different velocities of $\Rightarrow$ dextrorotatory & levorotatory motion are produced.

$$\frac{(\mu_o - \mu_e) t}{\lambda} \cdot 2\pi = \frac{\pi}{2}$$

$$\Rightarrow (\mu_o - \mu_e) = \frac{\lambda}{4t} = \frac{589 \times 10^{-9}}{4 \times 2.945 \times 10^{-5}} = \underline{\underline{0.005}}$$

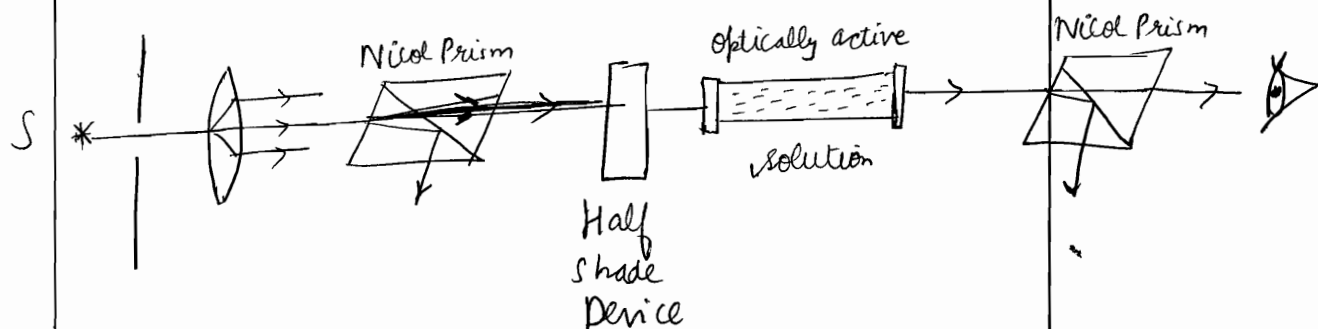
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Q5)

a) Q4

b) Half Shade Device

It's an instrument to measure the optical rotation of certain substances.

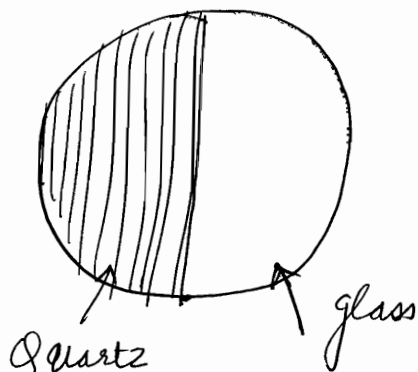


In the absence of Half Shade Device, in order to measure the angle θ of optical rotation, the precise measurement of complete darkness has to be observed. But eye cannot judge position of complete darkness accurately and field appears dark over a considerable range of rotation.

Half shade device overcomes this

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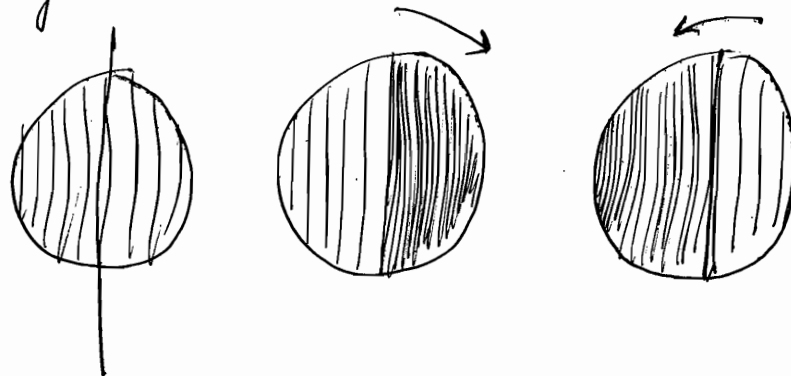
problem



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~~Half wave~~
Half shade device is combination of 2 semi circular plates, One of glass and other of quartz cut along the optic axis. Thickness of plate is such that it introduces a path difference of $\frac{\lambda}{2}$. (Note that Half shade device is constructed for a particular λ)

Due to such a structure, slight rotation left or right of the total darkness angle will lead to half illumination, and therefore the rotation angle could be accurately measured.



Q7)

$$\Delta\phi = \frac{2\pi}{\lambda} \cdot \Delta = \frac{2\pi}{\lambda} \cdot t(\mu_o \sim \mu_e)$$

$$= \frac{2\pi}{6000 \times 10^{-10}} \times 2.3 \times 10^{-5} \times 9.4 \times 10^{-3}$$

$$= \underline{\underline{0.72\pi}}$$

Elliptically Polarized Beam will come out with $\Delta\phi = \frac{\pi}{2} + 0.72\pi$

$$= \underline{\underline{1.2\pi}}$$

Q8)

Numerical

For ϕ WP : $\Delta\phi = (2n-1) \frac{\pi}{2}$

$$\Rightarrow \Delta = \frac{\lambda}{2\pi} \cdot (2n-1) \frac{\pi}{2} = (2n-1) \left(\frac{\lambda}{4} \right)$$

Also $(\mu_o \sim \mu_e) t = (2n-1) \cdot \frac{\lambda}{4}$

$$\Rightarrow t = \frac{(2n-1) \cdot \lambda}{(\mu_o \sim \mu_e) 4}$$

$$= \frac{5472 \times 10^{-10} (2n-1)}{4 \times 0.171}$$

$$= 0.8 \mu\text{m} (2n-1)$$

$$\boxed{n=20}$$

$$t = 0.8(39) \mu\text{m} \\ = \underline{\underline{31.2 \mu\text{m}}}$$

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Q9)

QWP

$$(1.5533 - 1.5443) 0.1436 \times 10^{-3} = \frac{\lambda}{4} (2n-1)$$

$$\frac{51696 \text{ \AA}}{2n-1} = \lambda$$

For λ should lie between $4000 \text{ \AA} \Delta 7500 \text{ \AA}$

$$\lambda = 7385.14 \text{ \AA} \quad n=4$$

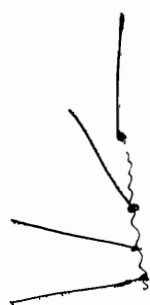
$$\lambda = 5744 \text{ \AA} \quad n=5$$

$$\lambda = 4699 \text{ \AA} \quad n=6$$

$$\lambda = 3976.6 \text{ \AA} \quad n=7$$

odd multiples
of $\left(\frac{\lambda}{4}\right)$

Q10)



$$\begin{aligned} I &= I_0 (\cos^2 30^\circ)^3 \\ &= I_0 \left(\frac{\sqrt{3}}{2}\right)^6 \\ &= I_0 \left(\frac{3\sqrt{3}}{8}\right)^2 = \underline{\underline{0.42 I_0}} \end{aligned}$$

Now $I_1 = I_0 \langle \cos^2 \theta \rangle$ ($\because$ initial light is unpolarized)

$$\Rightarrow I = I_0 * 0.4$$

around 21% of intensity is transmitted

Crossed
 $\Rightarrow \perp$

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Q11) After Polarizer, intensity = $\left(\frac{I_0}{2}\right)$
 After analyzer, intensity = $\frac{I_0}{2} \cos^2 \theta$
 Now we want,

$$\frac{I_0}{2} \cos^2 \theta = \frac{I_0}{4}$$

$$\Rightarrow \boxed{\theta = 45^\circ}$$

Q12) If we take 2 polarizers and put them in series with their optical axis $\perp$ to each other, we get 0 output. This proves that light is transverse.

Q13) अरे मूर्ख μ ती दे हो

Q14) Calcite : $\ominus$ ve Quartz : $\oplus$ ve

Q15) Chiral Molecule has non-superimposable mirror image. Optical Activity is rotation of linearly polarized light as it travels through a solution of chiral molecules. If the solution is composed of molecules of opposite chirality (called enantiomers), the rotation in opposite direction takes place.

Also rotation is proportional to concentration of solution i.e. more the molecules more the rotation !!

$$\boxed{\theta = \theta_c \cdot c \cdot l}$$

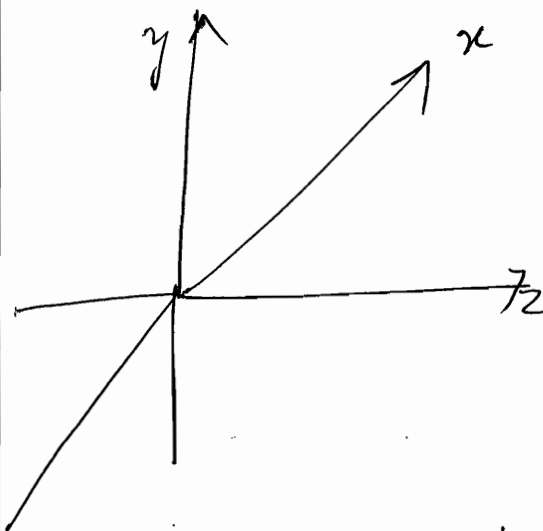
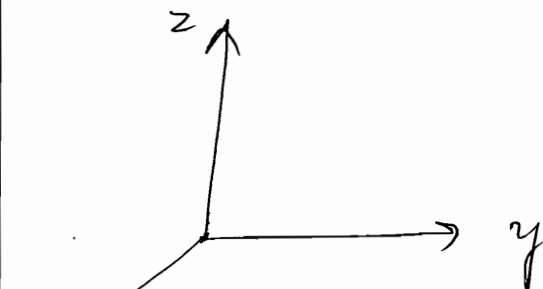
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Q16) $\theta = \tan^{-1} \left(\frac{n_z}{n_i} \right) = \tan^{-1}(1.33) = \underline{53^\circ}$ with vertical

<Q27 With horizon, $\angle = 36.9^\circ$

Q17)



For circular polarization,
1st condition $a = b = A$ (say) < derive the

$\Rightarrow E_y = A \cos(kx - \omega t)$ $\hat{y}$ ^{trig identity}
 $E_z = A \cos(kx - \omega t + \phi)$ $\hat{z}$ ^{equation} & then conclude!!

Take $\phi = \left(\frac{\pi}{2} \right)$

$\Rightarrow E_y = A \cos(kx - \omega t)$
 $E_z = -A \sin(kx - \omega t)$

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This is RLP i.e. clockwise from Point of View of Observer travelling with the wave.

$$\text{Take } \phi = -\frac{\pi}{2}$$

$$\Rightarrow E_y = A \cos(kx - \omega t) \hat{y}$$

$$E_z = A \sin(kx - \omega t) \hat{z}$$

This is LCP i.e. anticlockwise from Point of view of Observer goes along the wave.

Q18)

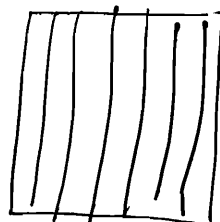
(i) Lecture 16

Quartz
⊕ve
Crystal

(ii) For Quartz

$$\mu_e = 1.554$$

$$\mu_o = 1.544$$



Remember this!!

$$\text{Now } \Delta = (1.554 - 1.544) t$$

$$\phi = \frac{2\pi}{\lambda} \Delta = \frac{\pi}{3}$$

$$\Rightarrow \Delta = \left(\frac{\lambda}{6}\right)$$

$$\Rightarrow 0.010 t = \frac{\lambda}{6}$$

$$\Rightarrow t = \frac{100\lambda}{6}$$

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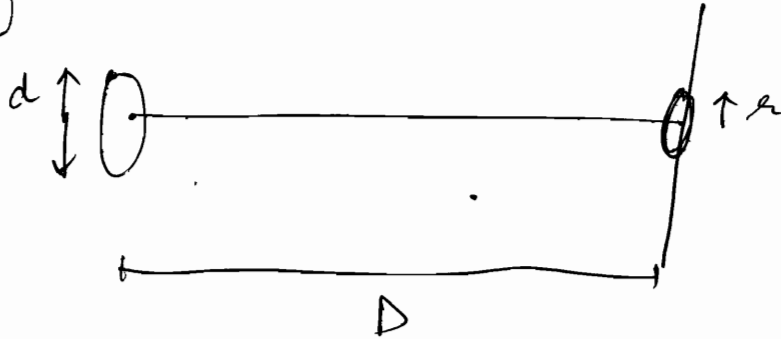
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Tut 13 LASERS

Q1)

Laser is diffraction limited, due to diffraction on airy pattern is formed.



For large distance

Put $\alpha = 2$

$$\frac{\pi d \sin \theta}{\lambda} = 2$$

$$\theta = \left(\frac{\lambda}{\pi a_0} \right)$$

$$r_f = \left(\frac{\lambda}{\pi a_0} \right)$$

$$r_f = \left(\frac{d \lambda}{\pi a} \right)$$

For the first ring of airy pattern

$$\alpha = 1.22 \pi \quad \text{where } \alpha = \frac{\pi d \sin \theta}{\lambda}$$

$$\Rightarrow \frac{d \sin \theta}{\lambda} = 1.22$$

$$\Rightarrow \frac{d}{\lambda} \cdot \frac{r}{D} = 1.22$$

$$\Rightarrow r = \frac{1.22 \lambda D}{d}$$

$$\Rightarrow \text{diameter} = 2.44 \times \frac{\lambda D}{d}$$

$$44.2 \text{ m}$$

$$= \frac{2.44 \times 6943 \times 10^{-9} \times 1000 \times 10^3}{10^{-2}}$$

$$= 169.4 \text{ m}$$

एक intensity
आएगी इसे बड़े
area में !!

Q2)

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Lecture 17



Let a consider a time of 1 s

$$\text{Energy emitted} = 10 \text{ mJ} = 10^{-2} \text{ J}$$

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For focusing of laser via a lens: we have [NOT CIRCULAR APERTURE]

radius of focussed spot = $\left(\frac{\lambda_0 f}{D_{\text{lens}}}\right) = r_0$

Intensity = $P / \pi r_0^2$

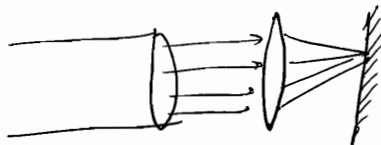
No. of pulses emitted = 500

Energy of 1 pulse = $\frac{10^{-2}}{500} \text{ J}$

Power of 1 pulse = $\frac{10^{-2} \times 10^9}{500 \times 3} = \underline{\underline{6666.6 \text{ W}}}$

$\lambda = 6000 \text{ \AA}$
 $d_{\text{laser}} = 1 \text{ cm}$
 $P_{\text{laser}} = 0.2 \text{ W}$

$r_0 = \frac{\lambda_0 f}{D} = \frac{6000 \times 10^{-10} \times 1}{2 \times 100} \times 10$



$f = 10 \text{ cm}$

For diffraction

$\alpha = 1.22\pi$

$\pi \frac{d \sin \theta}{\lambda} = 1.22\pi$

$\sin \theta = \frac{1.22 \lambda}{d}$

$\frac{r_0}{f} = \frac{1.22 \lambda}{d}$

$r_0 = \frac{1.22 \lambda f}{d} = \underline{\underline{7.32 \times 10^{-6} \text{ m}}}$

$A = \pi r_0^2 = 1.68 \times 10^{-10} \text{ m}^2$

$I = \frac{0.84 P}{A} = \underline{\underline{9.98 \times 10^8 \text{ W/m}^2}}$

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~~Note that the diameter of the airy pattern will be less or more depending upon where the image is formed.~~

It is not Airy disk

4.42×10^8

Q4)

Lecture 17

$$P = \frac{3J}{1 \times 10^{-5}} = \underline{\underline{3 \times 10^9 W}}$$

Q5)

Ruby laser

Q6)

Now we are talking !!

$$P(\text{stimulated emission}) = B_{21} u(\nu)$$

$$\text{Rate of stimulated emission} = N_2 B_{21} u(\nu)$$

$$P(\text{spontaneous emission}) = A_{21}$$

$$\text{Rate of spontaneous emission} = A_{21} N_2$$

$$\text{Now } \frac{P(\text{S.E.})}{P(\text{Sp.E.})} = \left[\frac{B_{21} u(\nu)}{A_{21}} \right]$$

Prob. vs Rate

Now let Probability of stimulate
absorption $P(\text{S.E.A.}) = B_{12} u(\nu)$

$$\text{Rate of stimulated absorption} = B_{12} u(\nu) N_1$$

Now at equilibrium

$$\text{rate of absorption} = \text{rate of emission}$$

$$\Rightarrow B_{12} N_1 u(\nu) = B_{21} N_2 u(\nu) + A_{21} N_2$$

$$\Rightarrow u(\nu) = \frac{A_{21} N_2}{B_{12} N_1 - B_{21} N_2}$$

$$= \left(\frac{A_{21}}{B_{21}} \right) \quad \text{--- (1)}$$

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$$\left(\frac{B_{12}}{B_{21}} \right) e^{\frac{h\nu}{kT}} - 1$$

Now we also know,

$$u(\nu) = \frac{8\pi h}{c^3} \nu^3 \cdot \frac{1}{(e^{\frac{h\nu}{kT}} - 1)} \quad \text{--- (2)}$$

$$\frac{8\pi p^2 dp}{h^3}$$

$$p = \frac{h\nu}{c}$$

$$\frac{8\pi \cdot \frac{h^2 \nu^2}{c^2} \cdot \frac{h \nu}{c} d\nu}{h^3}$$

$$\frac{8\pi h \nu^3 d\nu}{c^3 (e^{\frac{h\nu}{kT}} - 1)}$$

Comparing (1) and (2), we get

$$\frac{B_{12}}{B_{21}} = 1 \Rightarrow$$

$$B_{12} = B_{21}$$

$$\frac{A_{21}}{B_{21}} = \frac{8\pi h \nu^3}{c^3}$$

$$\frac{P(\text{st.e})}{P(\text{sp.e})} = \frac{B_{21} u(\nu)}{A_{21}}$$

$$= \frac{u(\nu)}{\frac{8\pi h \nu^3}{c^3}} = \left[\frac{1}{e^{\frac{h\nu}{kT}} - 1} \right]$$

For $\lambda = 6000 \text{ \AA}$
 $T = 10^3 \text{ K}$

$$\frac{h\nu}{kT} = \left(\frac{hc}{\lambda kT} \right) = 25$$

$$\Rightarrow \frac{P(\text{st.e})}{P(\text{sp.e})} = \frac{1}{e^{25}} \approx 10^{-10}$$

Hence Probability of stimulated emission is negligible.

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Q7)

E : Bohr

→ W : ✓

→ ~~W~~ : ✓

→ $U(W)$: ✓

Prob : $B_1 U$

Rate : $B_1 U N_2$

A_2 : given

$\frac{B_1}{A_2}$: known

A_2

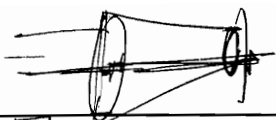
→ B_1 : ✓

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$$\text{Area} = \pi \left(\frac{\lambda f}{x_{\text{lens}}} \right)^2$$

$$\text{Intensity} = \left(\frac{P}{\text{Area}} \right)$$

(Q8)

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We know that laser is diffraction limit
 $\alpha = 1.22\pi$ (for central zone of airy pattern or Airy disk)

$$\frac{\lambda d \sin \theta}{\lambda} = 1.22\lambda$$

$$\sin \theta = \left(\frac{1.22 \lambda}{d} \right)$$

$$\frac{\lambda}{f} = \frac{1.22 \lambda}{d}$$

$$r = \frac{f 1.22 \lambda}{d}$$

$$\Rightarrow \text{Area} = \pi \left(\frac{1.22 f \lambda}{d} \right)^2 = 1.168 \times 10^{-10} \text{ m}^2$$

$$I = \frac{0.84 P}{\text{Area}}$$

$$= \frac{0.84 \times 3 \times 10^6}{\pi \left(\frac{1.22 f \lambda}{d} \right)^2}$$

$$= \frac{0.84 \times 3 \times 10^6 \times 10^{-4}}{\pi (1.22)^2 \times (5 \times 10^{-2})^2 \times (10000 \times 10^{-10})^2}$$

$$= \underline{\underline{2.156 \times 10^{16} \text{ W/m}^2}}$$

(Q9)

$$\checkmark N_1 = N_0 e^{-\frac{\Delta E}{kT}} \quad N_1 > N_0 \Rightarrow \frac{\Delta E}{kT} < 0 \Rightarrow T < 0$$

(Q10)

$$d \approx \lambda$$

$$\alpha = 1.22\pi \Rightarrow d \sin \theta = 1.22 \lambda$$

$$\Rightarrow \sin \theta = \frac{1.22 \lambda}{d}$$

$$\Rightarrow \frac{r}{f} = \frac{1.22 \lambda}{d} \Rightarrow r = 1.22 f \frac{\lambda}{d} \approx \underline{\underline{1.22 \lambda}}$$

$$\text{Area} = \pi (1.22 \times 6328 \times 10^{-10})^2 = 1.87 \times 10^{-12} \text{ m}^2$$

★ Focussing of lens is not Airy disk,
 so do not use 84%. To calculate x_0 , use

$$\alpha = 1.22\pi$$

$$r_0 = \frac{\lambda f}{d}$$

Q11)

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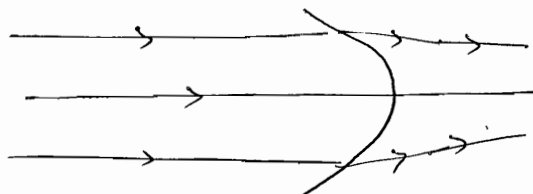
Q12)

Lecture 17

Self Focussing of Laser Beam

A non-linear optical phenomenon due to the variation of refractive index due to high intensity of laser as a result of which medium acts a converging lens. The beam is said to undergo self focussing.

$$n^2 = n_0^2 \left[1 - \left(\frac{I}{I_0} \right)^2 \right]$$



The medium acts as converging lens
of focal length = $\frac{\pi}{2} \alpha$

Therefore, the spreading of the beam due to diffraction will be balanced by self focussing, due to high directionality of LASERS can be explained.

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Q13)
Q14)
Q15)
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lecture 17

For rate of stimulated emission to be equal to rate of spontaneous emission

$$B_{21} u(\nu) N_2 = A_{21} N_2$$

$$\Rightarrow u(\nu) = \frac{A_{21}}{B_{21}}$$

$$\Rightarrow e^{\frac{h\nu}{kT}} = 2$$

$$\Rightarrow \frac{hc}{\lambda kT} = \ln 2$$

$$\Rightarrow T = \frac{hc}{\ln 2 \lambda k} = \underline{\underline{41.28 \times 10^3 \text{ K}}}$$

(ii) copy

(iii) Fiber attenuation = $10 \log_{10} \left(\frac{P_{in}}{P_{out}} \right) \text{ dB}$

$$= \frac{10}{100} \log_{10} \left(\frac{1}{0.3} \right) \text{ per unit length}$$

$$= \underline{\underline{0.052 \text{ dB/m}}}$$

Q17) $\frac{1}{2} \epsilon_0 E_0^2 c = \text{Intensity} = \frac{\text{Power}}{\pi r_0^2}$

$$\Rightarrow E_0 = \sqrt{\frac{2 \text{Power}}{\epsilon_0 c \pi r_0^2}} \quad \checkmark \quad r_0 = \left(\frac{\lambda f}{a} \right)$$

$$k = \left(\frac{2\pi}{\lambda} \right) \checkmark \quad \omega = \left(\frac{2\pi c}{\lambda} \right) \checkmark$$

$$E_0 = E_0 \sin(kz - \omega t)$$

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(Q1)

lecture 18

(Q5)

(Q7)

(Q8)

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(Q6) Not in Syllabus

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(Q2)

Coherence Length

A wave which appears to be a pure sine wave for an infinitely large period of time, or in an infinitely extended space, is said to be a perfectly coherent wave. In such a wave, there is a definite relationship between the phase of the wave at a give time and at a certain time later, or at a given point and at a certain distance away. No actual light source, however, emits a perfectly coherent wave. Light waves that are coherent only for a limited of time or for a limited space are called "partially coherent waves".

✓ The average time-interval for which the field remains sinusoidal is known as "coherence time" or "temporal coherence". Its denoted by T .

✓ The distance L for which field is sinusoidal is called coherence length

$$L = TC$$

★ For Young slit : Spatial Coherence is reqd.
For Michelson : temporal coherence is reqd.

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In interference experiments, distance
after which pattern disappear is
called coherence length.

$$L = \frac{\lambda^2}{\Delta\lambda}$$

$$= \frac{623.8 \times 623.8 \times 10^{-8}}{2 \times 10^{-12}}$$

$$= \underline{\underline{0.194 \text{ m}}}$$

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Proof For disappearance of pattern,
we have : $\frac{2d}{\lambda} - \frac{2d}{\lambda + \frac{\Delta\lambda}{2}} = \frac{1}{2}$

$$\frac{2d\Delta\lambda}{2(\lambda)(\lambda + \frac{\Delta\lambda}{2})} = \frac{1}{2}$$

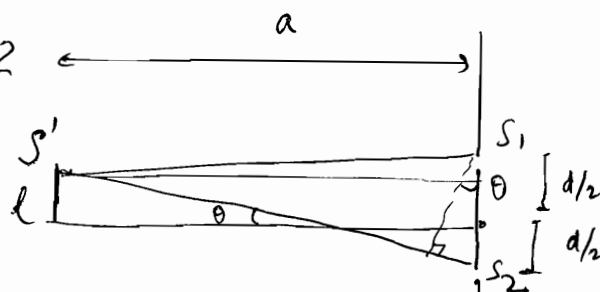
$$\frac{2d\Delta\lambda}{\lambda^2} = 1 \Rightarrow \boxed{2d = \frac{\lambda^2}{\Delta\lambda}}$$

Q3) Q2

Q4)

(a) Q2

(b)



$\Delta \ll 1$ i.e. $d \cdot \sin \theta \ll 1$ for good coherence

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Therefore for an extended incoherent source,
interference fringes of good contrast will
be formed when

$$l \ll \frac{\lambda a}{d}$$

$$s_1 s_2 = \left[a^2 + \left(\frac{d}{2} + l \right)^2 \right]$$

$$s_1 s_1 = \left[a^2 + \left(\frac{d}{2} - l \right)^2 \right]$$

$$\Rightarrow s_1 s_2 - s_1 s_1 \approx \frac{1}{2a} \cdot 2l \cdot d = \left(\frac{dl}{a} \right)$$

~~For fringes to disappear, $\left(\frac{dl}{a} \right) = \left(\frac{\lambda}{2} \right)$ for good contrast.~~

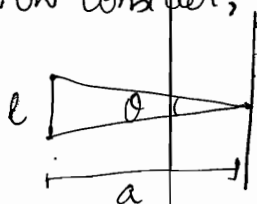
For fringes to disappear,

$$s_1 s_2 - s_1 s_1 = \left(\frac{\lambda}{2} \right) \Rightarrow l = \left(\frac{\lambda a}{2d} \right)$$

$\therefore$ If we have a linear source of dimension $\frac{\lambda a}{d}$, for every point on the source, there exists a point at distance $\left(\frac{\lambda a}{2d} \right)$ which produces fringes shifted by λ .

$$\therefore \text{for good pattern } l \ll \frac{\lambda a}{d}$$

Now consider,



$$\text{Now } \theta \approx \left(\frac{l}{a} \right)$$

$$\Rightarrow d \ll \left(\frac{\lambda}{\theta} \right)$$

P.T.O.

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P-17.5

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छोड़ी आसान भाषा में करें तो,

The path difference between the rays
reaching the slits, should be much less
than λ

$$\Delta x \ll \lambda$$

Similar to Δx being $d \sin \theta$ in Young
double slit, here Δx is $l \sin \theta$

$$\Rightarrow l \sin \theta \ll \lambda$$

$$\Rightarrow \frac{dl}{a} \ll \frac{\lambda}{d}$$

$$\Rightarrow l \ll \frac{\lambda a}{d}$$

$$\theta = \left(\frac{l}{a} \right)$$

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